

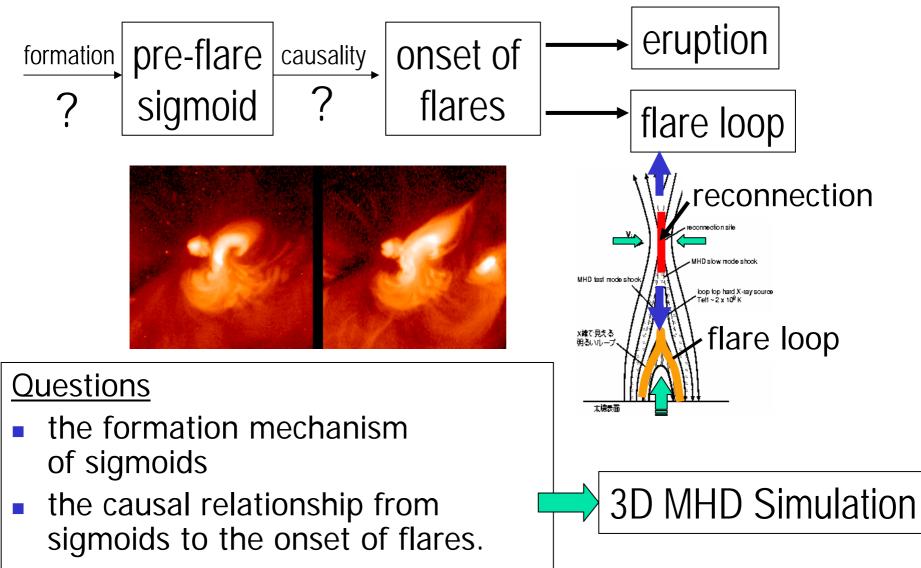


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Simulation Study on the Self-Organization of Sigmoidal Structure and the Onset of Solar Flares

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Introduction



Mechanism of sigmoid formation

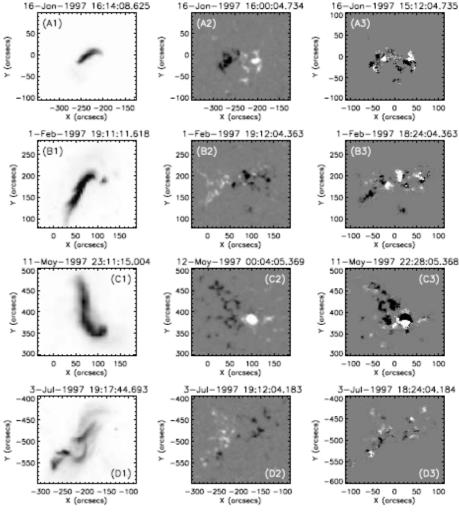
- kink mode instability
 - proposed by Rust & Kumar 1996
 - simulations
 - Fan & Gibson 2004, Fan 2005
 - Kliem, Titov & Török 2004, Török & Kliem 2005

However, the several observations suggested that the field twist in sigmoids is insufficient for the kink instability. (Leamon et al. 2003, Yamamoto et al. 2005)

■ resistive tearing instability (Kusano 2005 ApJ in press)
our thin current sheet (reversed-shear layer)
model → self-organization → sigmoid formation
→ double reconnections → flare onset

Helicity Injection Measurement

Yamamoto et al. 2005 ApJ (method developed by Kusano et al. 2002) r (arcse LCT + induction equation $\left[\frac{\partial \mathbf{B}}{\partial t}\right] = \left[\nabla \times (\mathbf{V}_t \times \mathbf{B}_n + \mathbf{V}_n \times \mathbf{B}_t)\right]_n$ $\dot{H} = \int \mathbf{E} \times \mathbf{A}_{p} \cdot d\mathbf{S}, \ \mathbf{E} = \mathbf{V} \times \mathbf{B}$ $\Omega = \frac{H}{\phi^2}$ normalized twist rate $(24h) \cdot \Omega \approx O(0.01) \ll 1$ much less than one turn twist



Simulation Model (3D MHD)

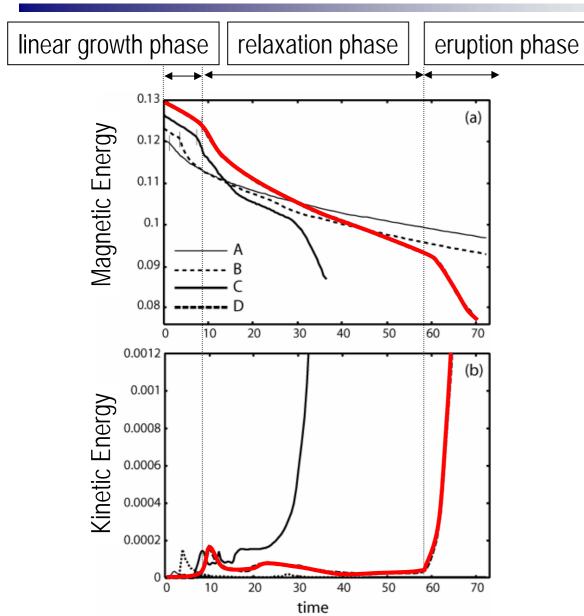
- **Basic equations** $\frac{\partial \mathbf{V}}{\partial t} + \mathbf{V} \cdot \nabla \mathbf{V} = \mathbf{J} \times \mathbf{B} + v \nabla^2 \mathbf{V}, \quad \frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{V} \times \mathbf{B} \eta \mathbf{J})$
 - zero- β version & finite β version
- Finite difference: 512 X 512 X 1024 (Δ~10⁻³)
 - parallelization in terms of MPI library (domain decomposition)
- Boundary Condition
 - line-tide on the bottom, periodic for the axial

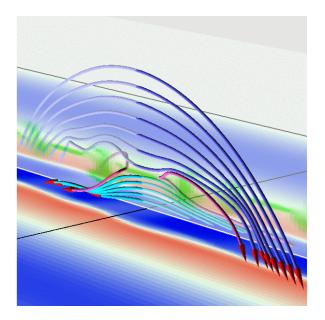
(current sheet)

- Initial state
 - force-free field with <u>reversed-shear layer</u>
- Anomalous resistivity

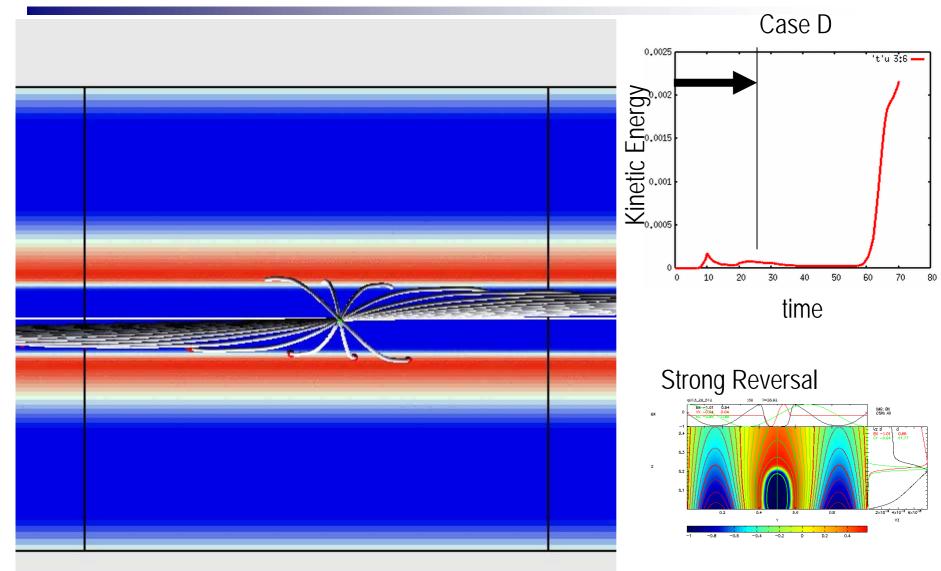
$$\eta = \begin{cases} \eta_0 & (J < J_0) \\ \eta_0 + \eta_1 \frac{J - J_0}{J_0} & (J > J_0) \end{cases}$$
$$\eta_0 = 10^{-5}, \quad \eta_1 = 5 \times 10^{-4}$$

Time Evolution of Energy

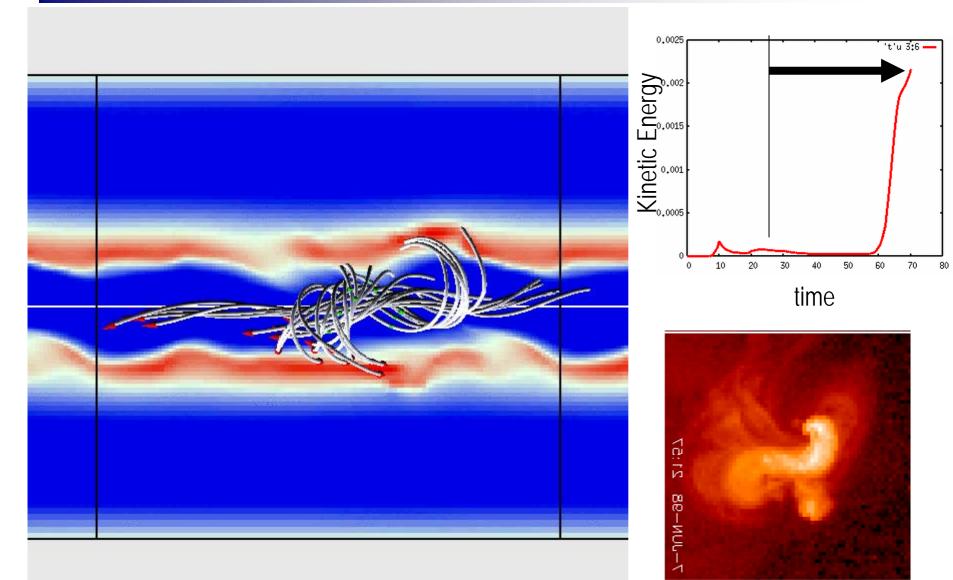




Formation of Sigmoidal Structure



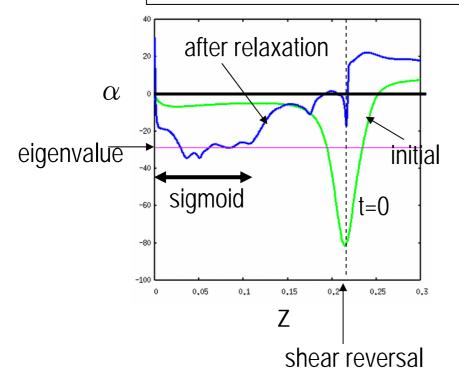
Formation of Sigmoidal Structure

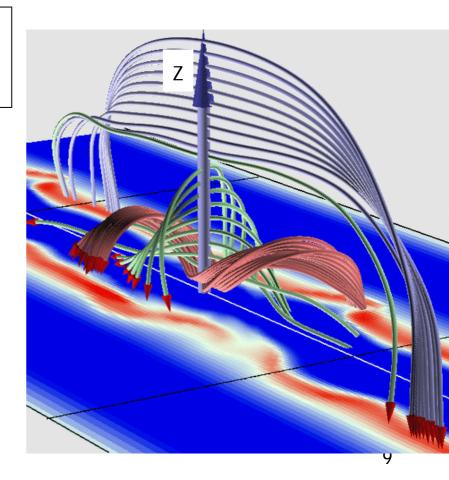


distribution of $\boldsymbol{\alpha}$ in sigmoid

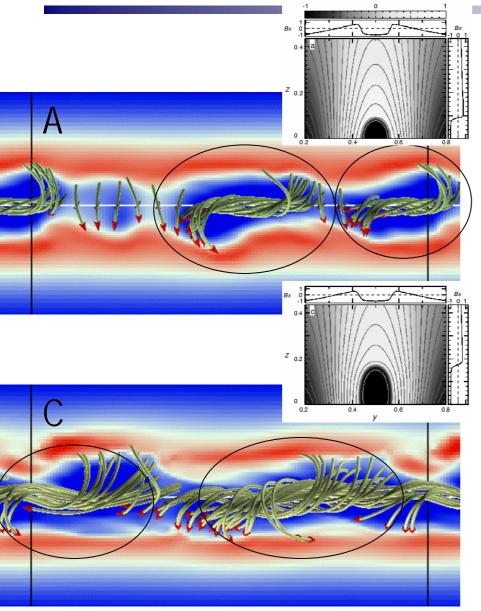
- α is flattened inside the sigmoid.
- α is limited by eigenvalue.

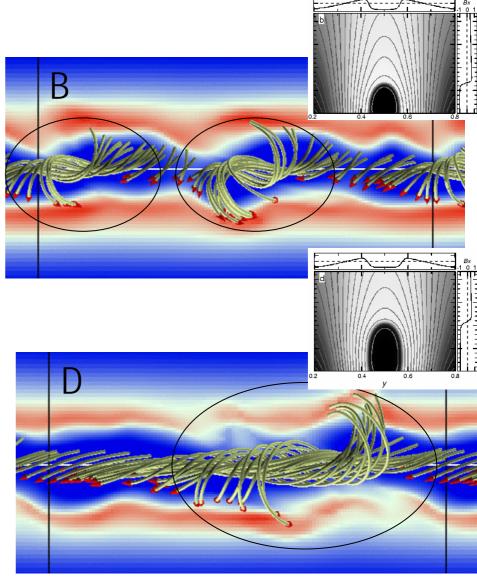
consistent with Taylor's relaxation theory





dependency on the initial state

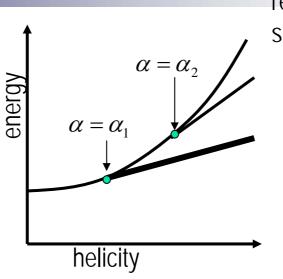


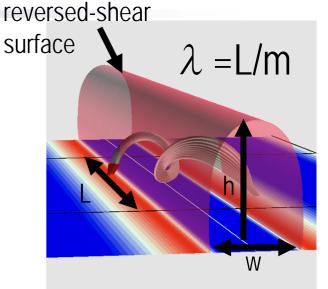


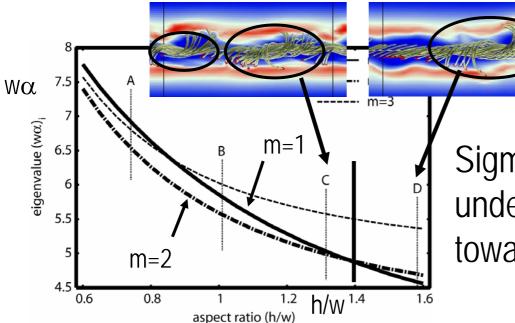
Comparison with Taylor relaxation

LFFF bifurcates into multiple solution, when α is larger than the eigenvalue α_1 . (Taylor, 1986)

 $\nabla \times \mathbf{B} = \alpha_i \mathbf{B}$



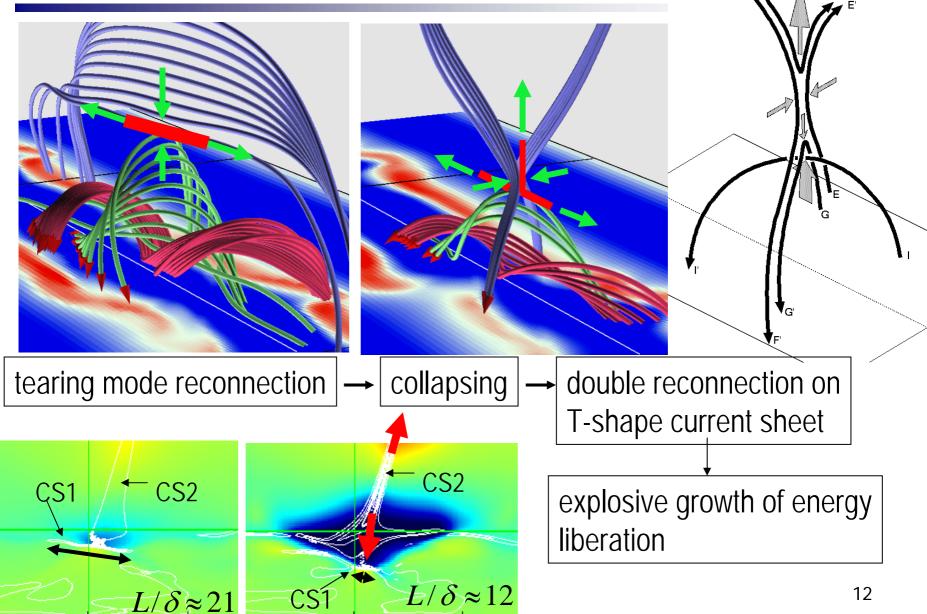




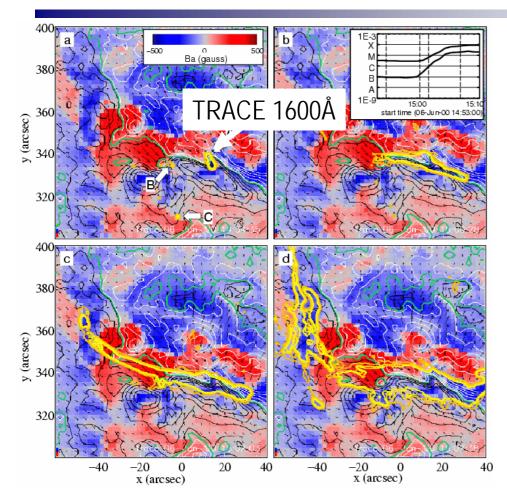
Sigmoidal wavelength is consistent with the Taylor's minimum energy principle.

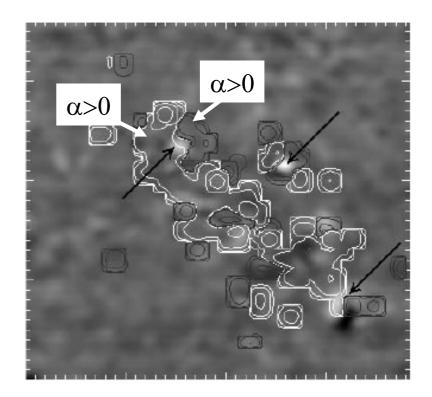
Sigmoidal formation can be understood as the self-organization toward the Taylor type state.

Sigmoid to Eruption



flares from reversed-shear

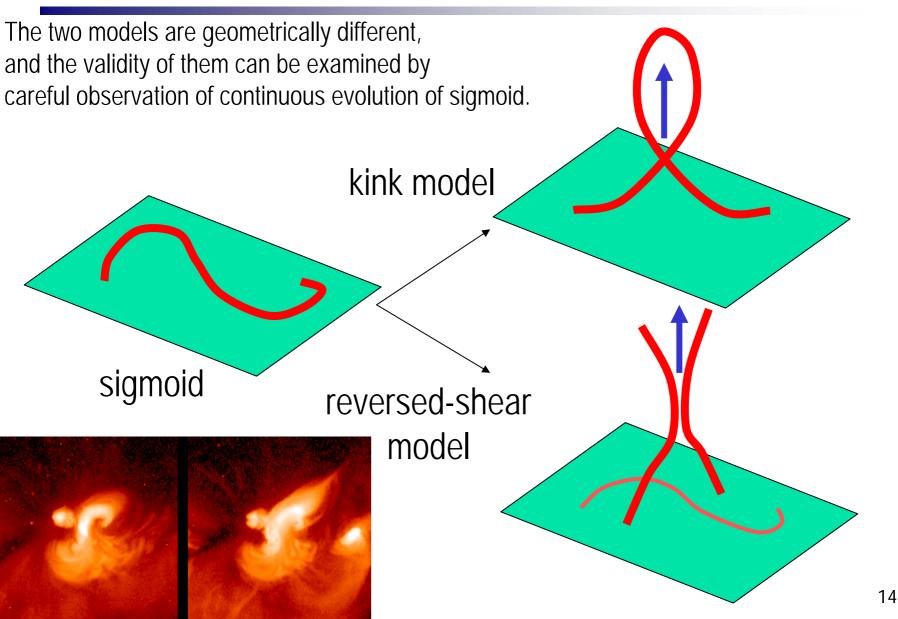




Hahn et al. 2005

Maeshiro et al. 2005

Kink Model vs. Reversed-Shear Model



Conclusions

- The resistive tearing mode instability on the reversed-shear layer may cause both the formation of sigmoids and the onset of eruption.
- The sigmoidal formation is consistent with the self-organization toward Taylor state.
- The transition from the quasi-steady sigmoid to the sudden onset of eruption can be explained by the arcade collapsing and the feed-back of double reconnections.
- Reversed-shear flare model.
 - Kusano et al. 2004 ApJ, Kusano 2005 ApJ (in press)