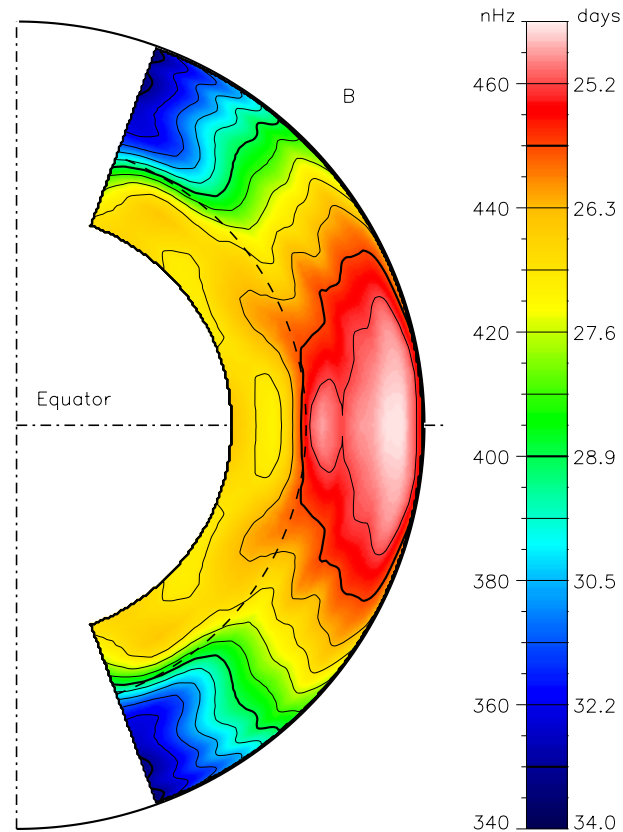
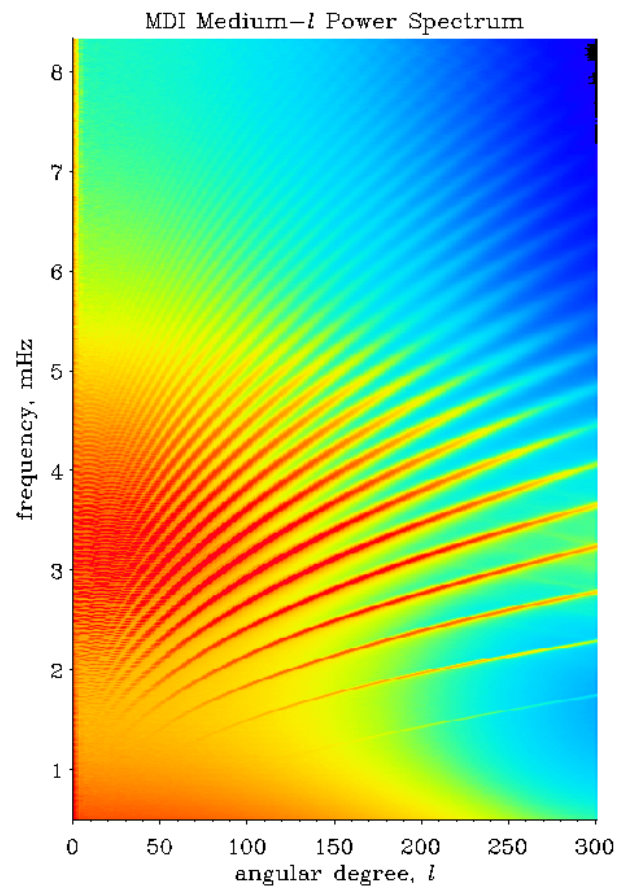


IMPRS, February 2002

# Helioseismology and Internal Rotation

Dieter Schmitt (Katlenburg-Lindau)



IMPRS, 2/2002

Helioseismology and  
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D. Schmitt

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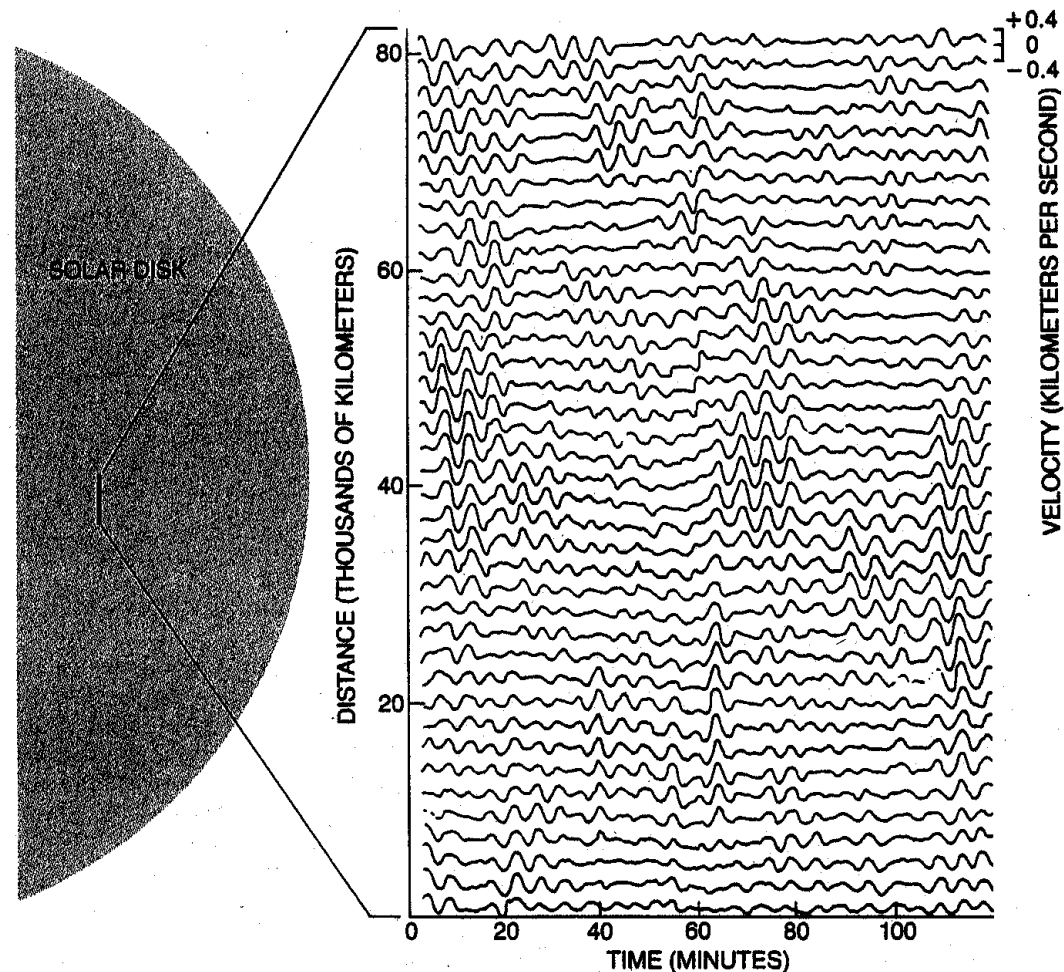
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# 1. Introduction

- Earth quakes  $\leadsto$  compressional and shear waves  $\leadsto$  propagation through interior  $\leadsto$  information about Earth  $\leadsto$  **seismology**
- free oscillations of Earth  $\leadsto$  analogy to Sun
- Leighton, Noyes & Simon (1960): 5-min oscillation  
 $P \approx 5 \text{ min}$ ,  $v \approx 1/2 \text{ km s}^{-1}$   
spatial coherence  $\approx 30\,000 \text{ km}$ , temporal coherence  $\approx 1/2 \text{ h}$



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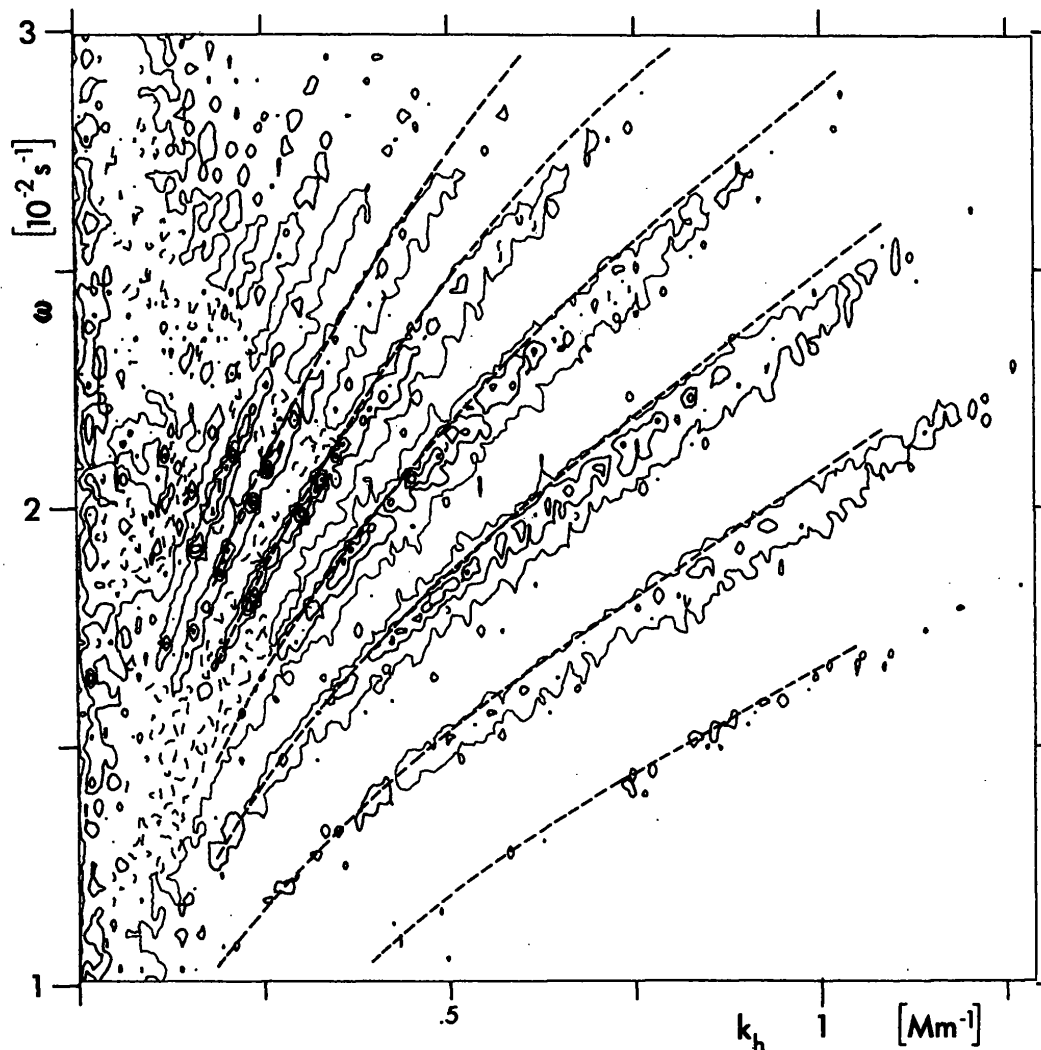


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- Ulrich (1970), Noyes & Simon (1971): superposition of many ( $\sim 10$  Mio) discrete global acoustic oscillations of Sun individual amplitudes  $< 20 \text{ cm s}^{-1}$
- Deubner (1975): confirmation by observation spatial and temporal Fourier analysis, power  $k_h \omega$ -diagram, ridges

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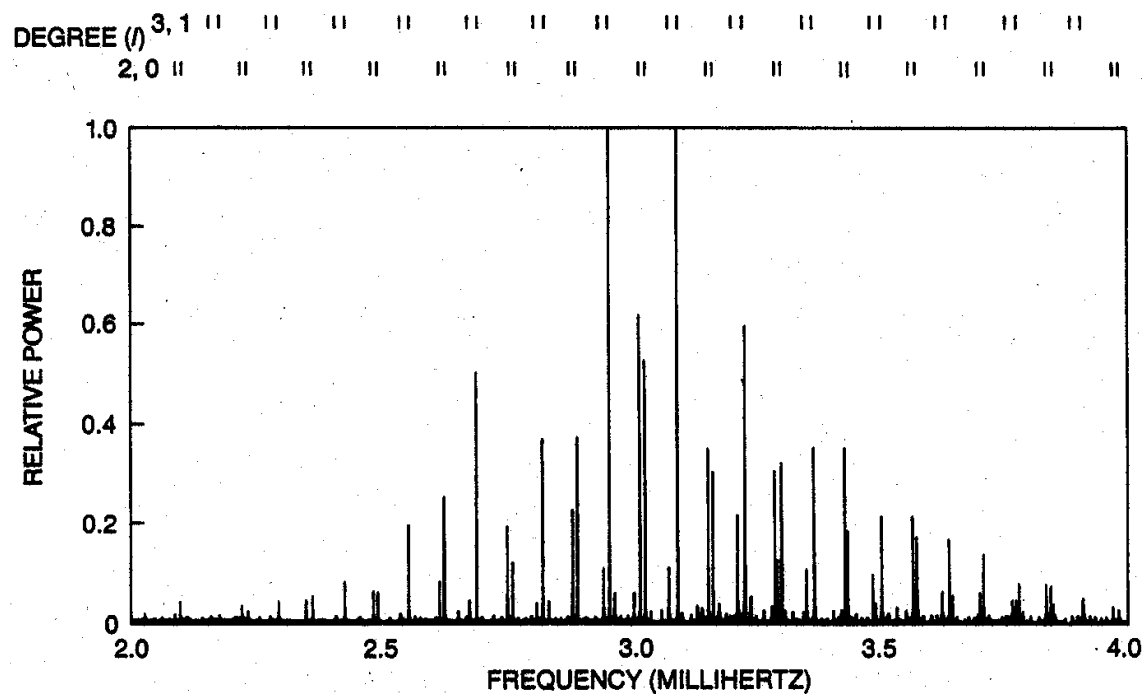
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- Isaak et al.: discrete oscillations of spatially unresolved Sun  $\leadsto$  basis for **asteroseismology**



- Each wave contains information about the solar interior, although largely smoothed  
appropriate combination of certain modes narrows contribution range  $\leadsto$  spatial information  $\leadsto$  **helioseismology**
- Physics of waves, approximate theory

## 2. Observational Constraints

- Spectral lines, Doppler shift  $\leadsto v(r, t)$
- Intensity oscillations, solar irradiance variation
- Fourier transformation  $f(\mathbf{k}, \omega) = \int v(\mathbf{r}, t) \exp(-i\mathbf{k}\cdot\mathbf{r} - i\omega t) d\mathbf{r} dt$
- power  $p(k, \omega) = ff^*$ ,  $k = |\mathbf{k}|$
- Nyquist theorem, time resolution,  $\geq 2$  measurements per period, e.g. every 90 s
- Frequency resolution,  $T \geq 2\pi/\Delta\omega$   
 $\Delta\omega/\omega \approx \Omega/\omega \approx 10^{-4}$ ,  $T \geq 30$  days
- Night gaps  $\leadsto$  side peaks  
 $\leadsto$  south pole, GONG, SOHO
- Spatial resolution, small wave length, large wave number
- Wave number resolution, whole Sun

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### 3. Properties of p-Modes

- Sound and gravity waves or p- and g-modes
- Internal gravity waves  $\neq$  gravitational waves  
gravitation / buoyancy restoring force  
levels inhomogenities on horizontal surfaces  
g-modes propagate in convectively stable layers  
i.e. in solar interior, evanescent in convection zone  
frequencies small  $< N$ ,  $N$  buoyancy or Brunt-Väisälä frequency  
periods  $\geq 30$  min, not yet observed,  
thus not used in helioseismology
- Sound waves  
pressure gradient restoring force  
compression and expansion, longitudinal waves  
periods small 3 . . . 12 min, maximum power at 5 min  
individual amplitudes small, linear treatment

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- Infinite homogenous fluid

- wave equation for pressure perturbation  $\ddot{p}_1 - c^2 \Delta p_1 = 0$   
with  $c^2 = \gamma p_0 / \rho_0 = \gamma R T_0$ ,  $c$  adiabatic sound speed
- plane wave solution  $p_1(\mathbf{r}, t) = \hat{p}_1 \exp i(\mathbf{k} \cdot \mathbf{r} - \omega t)$
- angular frequency  $\omega$ , wave vector  $\mathbf{k}$
- dispersion relation  $\omega^2 = k^2 c^2$  with  $k^2 = |\mathbf{k}|^2$

- Sun: stratification  $c(r)$  and spherical geometry

$$p_1(r, \theta, \phi, t) = \hat{p}_1(r) Y_l^m(\theta, \phi) \exp(-i\omega t)$$

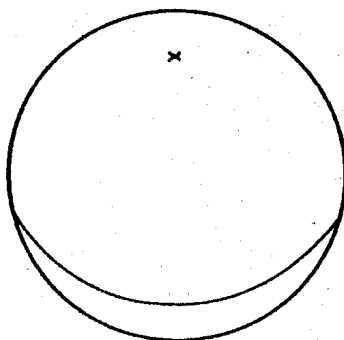
- Spherical harmonics  $Y_l^m(\theta, \phi) = P_l^m(\cos \theta) \exp(im\phi)$ ,  
Legendre functions  $P_l^m$

degree  $l = 0, 1, 2, \dots$ , order  $m = -l, \dots, 0, \dots, +l$

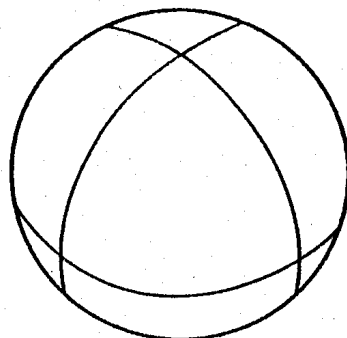
$l$  nodal lines,  $l - |m|$  meridional nodal lines,

$|m|$  azimuthal nodal lines

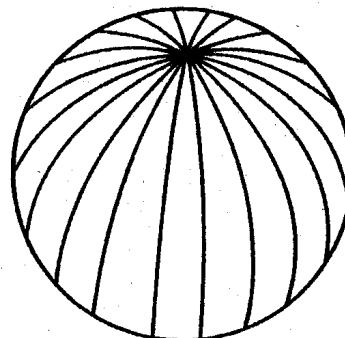
$l=1 \quad m=0$



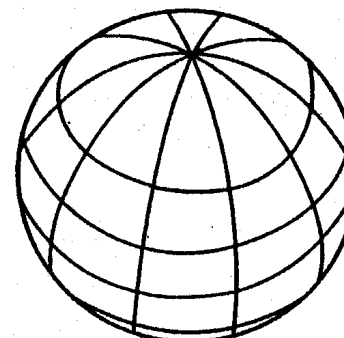
$l=3 \quad m=2$



$l=10 \quad m=10$



$l=10 \quad m=5$


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$l, m$  determine horizontal wave numbers  $k_\theta, k_\phi$   
discrete numbers: wavelengths must fit on spherical surfaces  
 $Y_l^m$  eigenfunctions of horizontal Laplace operator

$$\Delta_\theta^\phi Y_l^m = -\frac{1}{r^2} \left[ \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right] Y_l^m = \frac{l(l+1)}{r^2} Y_l^m$$

$\leadsto$  eigenvalue  $k_h^2 = k_\theta^2 + k_\phi^2 = \frac{l(l+1)}{r^2} \geq 0$  not dependent on  $m$

$$k_\phi^2 = \frac{m^2}{r^2 \sin^2 \theta} \geq 0, \quad k_\theta^2 = \frac{l(l+1) - m^2 / \sin^2 \theta}{r^2} \begin{cases} > 0 \text{ running} \\ < 0 \text{ evanescent} \end{cases}$$

running for  $\sin^2 \theta \geq \frac{m^2}{l(l+1)}$  latitudinal info for helioseismology

- No rotation  $\leadsto$  no prescribed pole  
 $m$  does not occur in dispersion relation, degeneration  
 $2l(l+1)$  eigenfunctions have same eigenfrequency  $\omega$

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- Information on depth

stratification, coefficients of wave equation depend on  $r$

local analysis,  $\hat{p}_1(r) = \rho_0^{1/2} \exp(ik_r r)$

approximate dispersion relation for  $\omega > N$ :

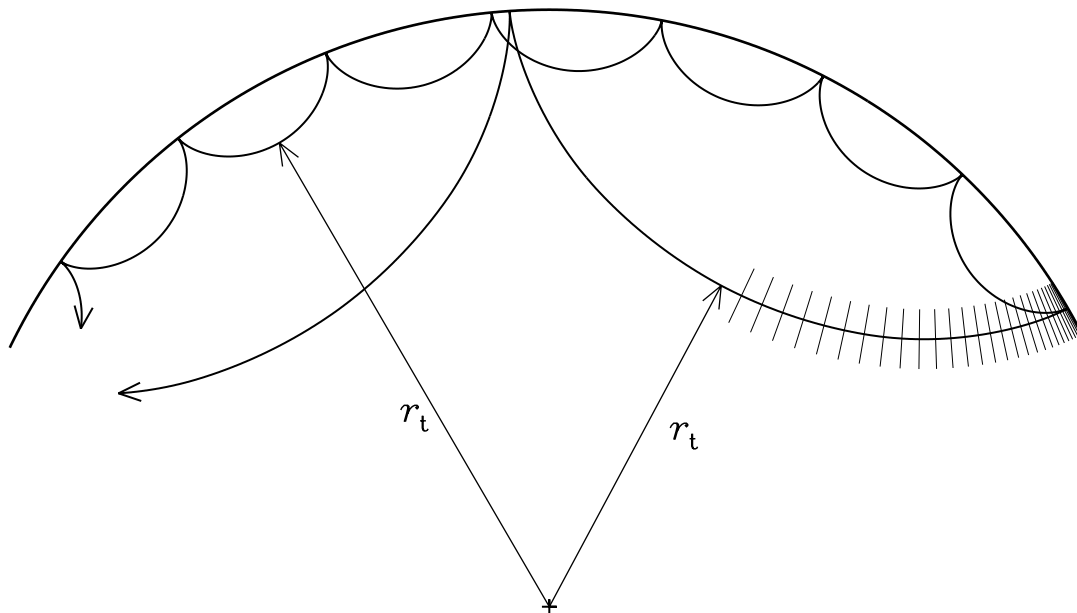
$$\omega^2 = (k_r^2 + k_h^2) c^2 + \omega_{ac}^2$$

- Acoustic cut-off:  $\omega > \omega_{ac}$ ,  $\omega_{ac} = c/2H$

$c$ ,  $\omega_{ac}$ ,  $k_h$  functions of  $r$   $\curvearrowright$   $k_r(r)$  local wavenumber

$$k_r^2 = \frac{\omega^2 - \omega_{ac}^2}{c^2} - \frac{l(l+1)}{r^2} \begin{cases} > 0 \text{ running} \\ < 0 \text{ evanescent} \end{cases}$$

$r \searrow$   $c \nearrow$   $k_r \searrow$


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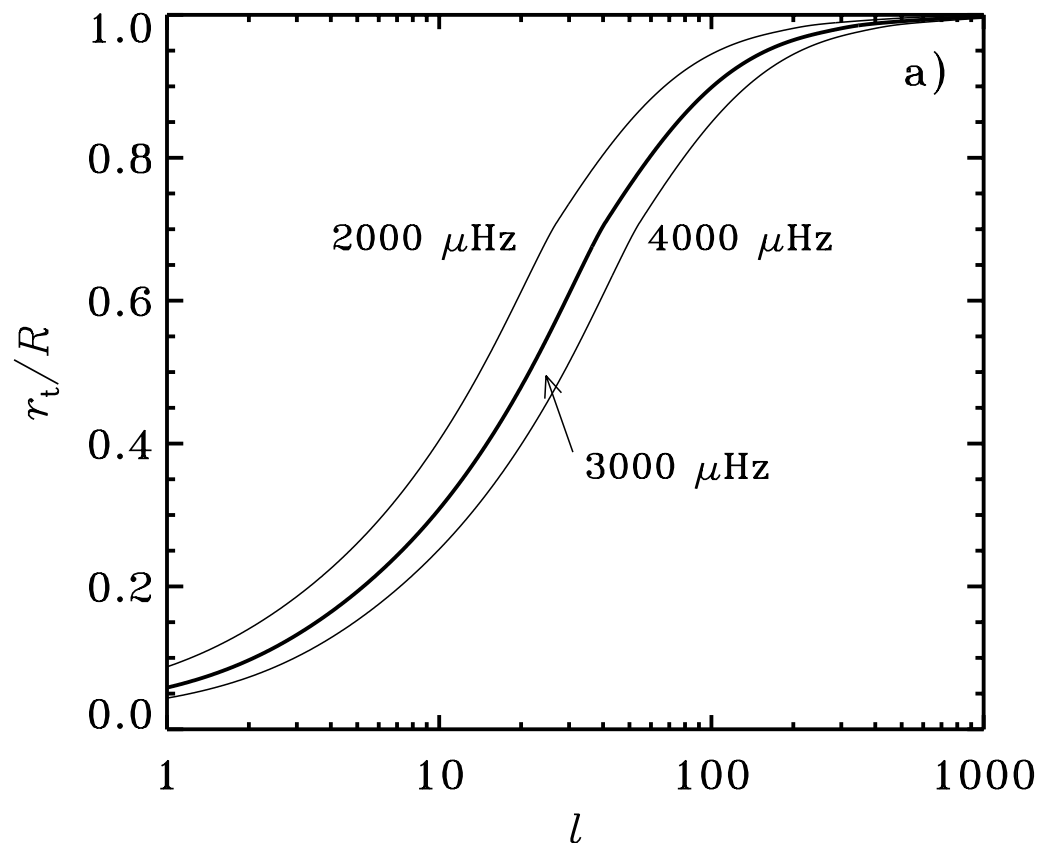
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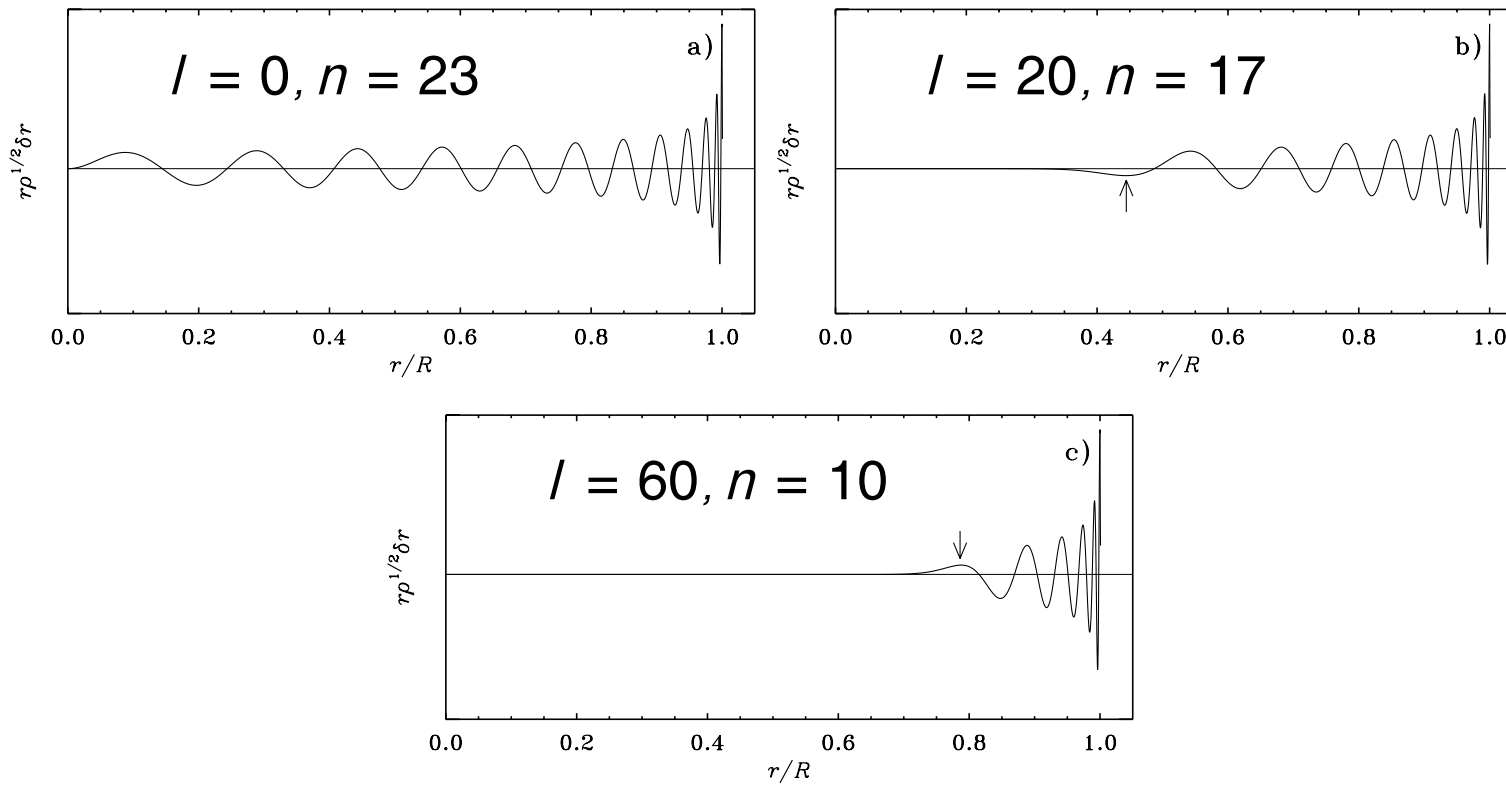
- Inner turning point  $r_t$  where  $k_r = 0$ , in solar interior  $\omega_{ac} \ll \omega$

$$\frac{\omega^2}{c^2(r_t)} = \frac{l(l+1)}{r_t^2} \quad \text{or} \quad r_t = \frac{\sqrt{l(l+1)}}{\omega} c(r_t) \quad \text{i.e. function of } \frac{l}{\omega}$$



- Upper reflection point near surface,  $\omega_{ac}$  increases ( $H$  decreases)  
 $R_t$  approx given by  $\omega = \omega_{ac}(R_t)$ ,  $R_t \approx r_{\odot}$
- Cavity, constructive interference, discrete spectrum of standing sound waves

- Overtones, index  $n$



- Individual mode: three indices  $l, m, n$

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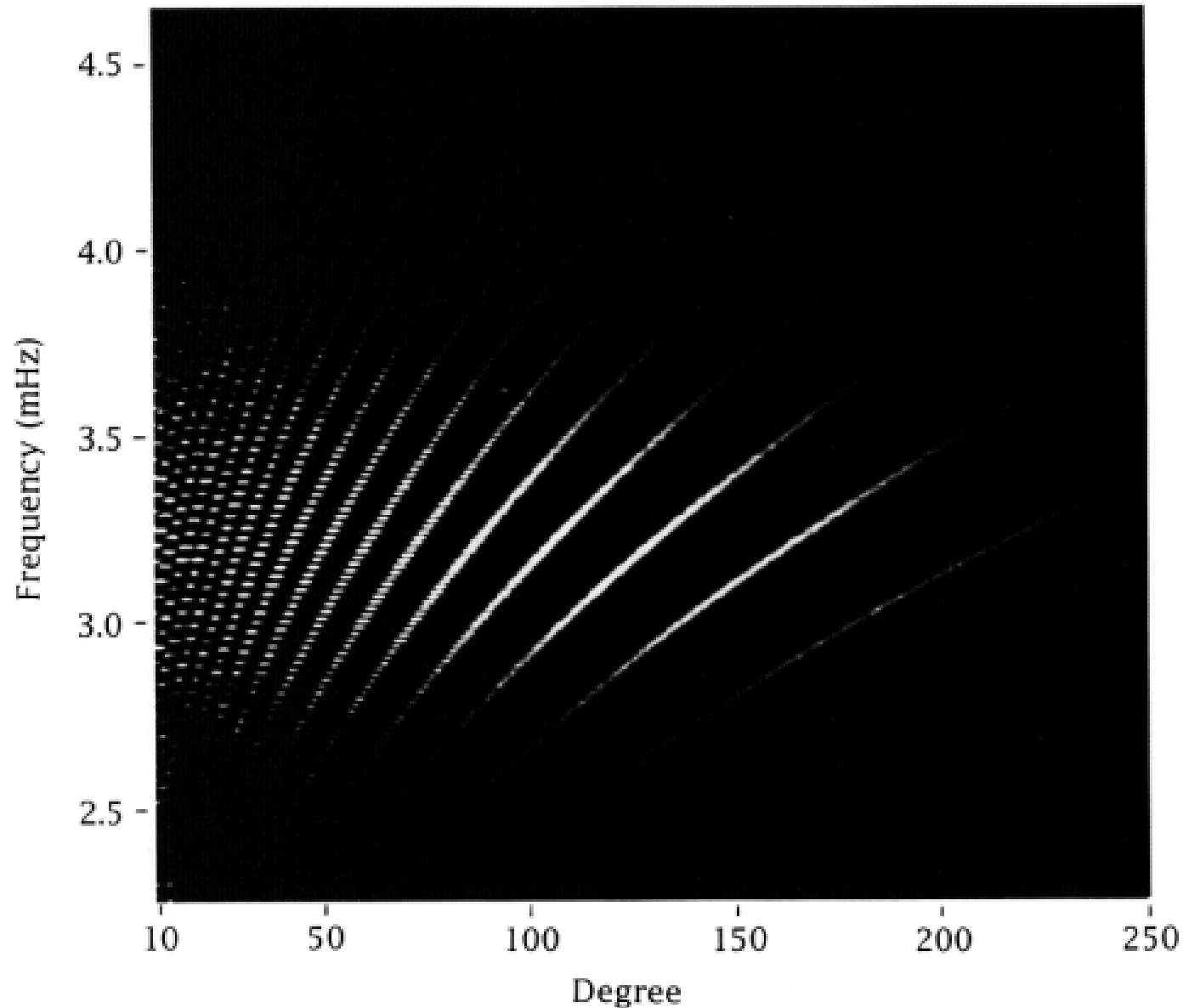


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- Interpretation of Deubner's observation,  $\omega(n, l) \leftrightarrow k_h(l)$   
ridges various  $n$ ,  $\omega(l)$  not resolved



Duvall et al. (1988)

south pole, 50 h, resolution  $\Delta l \approx 3 \dots 5$

$\Delta l = 1$  needs information from around Sun ( $l \approx k_h r_\odot$ )

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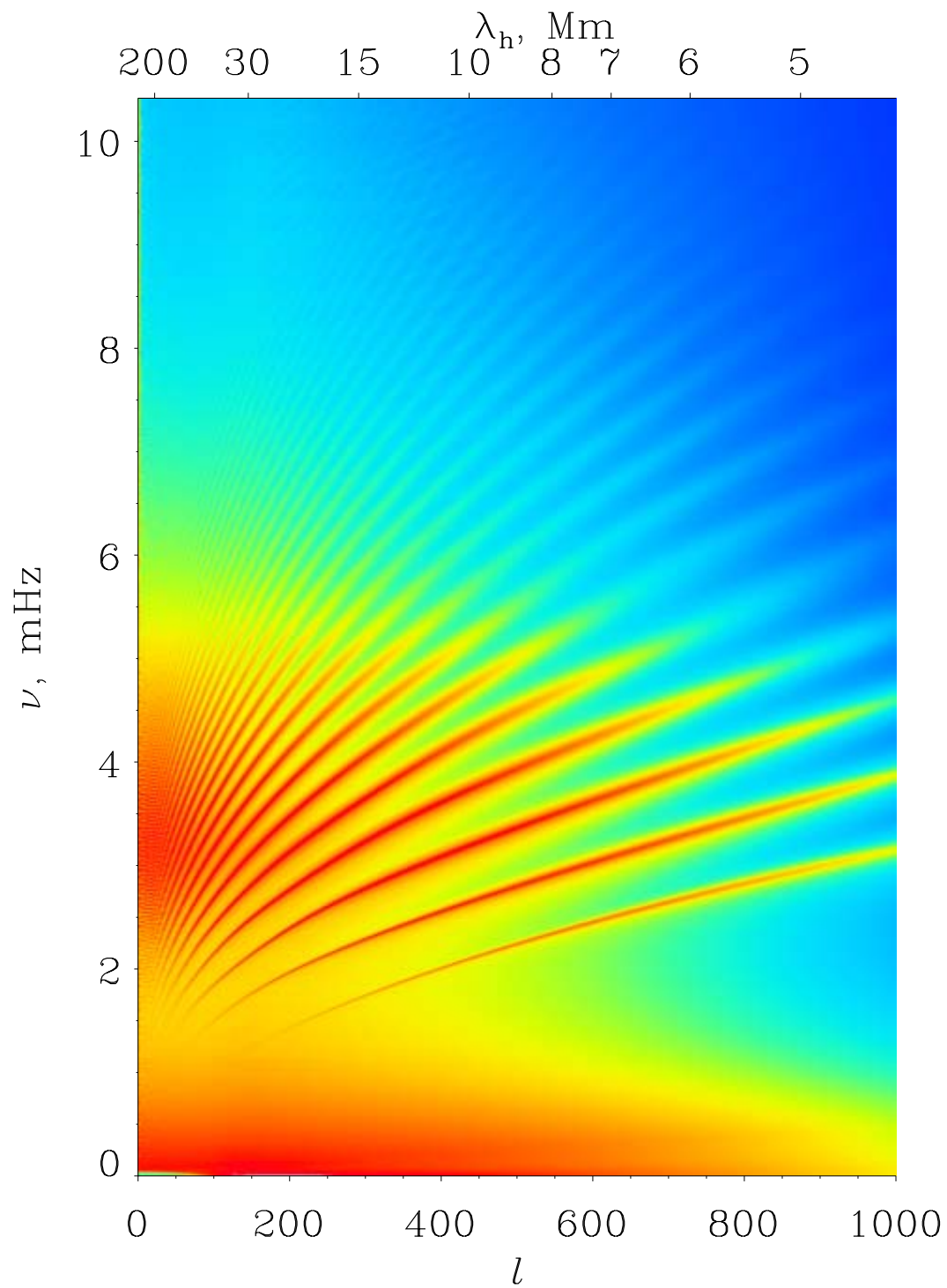


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• SOHO data



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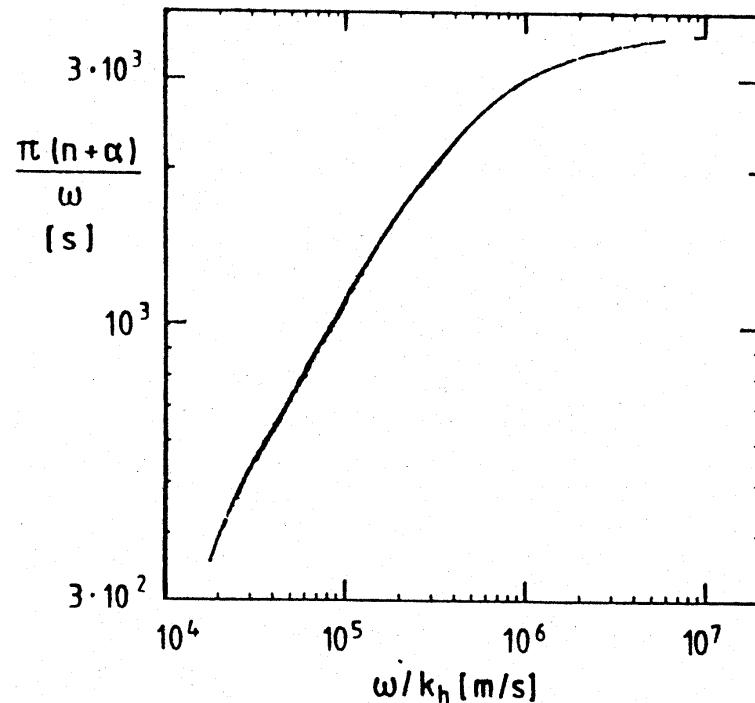


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• Duvall's law:



$$\frac{\pi(n + \alpha)}{\omega} \leftrightarrow \frac{\omega}{k_h} \quad \text{one curve}$$

phase difference  $\Delta\Psi = \int_{r_t}^{r_\odot} k_r dr = \pi(n + \alpha)$

$\alpha$  because of evanescent boundaries,  
string  $\alpha = 0$ , organ pipe  $\alpha = 1/2$

$$r_t = \left( \frac{l(l+1)}{\omega^2} \right)^{1/2} c(r_t), \quad k_r = \frac{\omega}{r} \left( \frac{r^2}{c^2} - \frac{l(l+1)}{\omega^2} \right)^{1/2}$$

$$\frac{\pi(n + \alpha)}{\omega} = \int_{r_t}^{r_\odot} \left( \frac{r^2}{c^2} - \frac{l(l+1)}{\omega^2} \right)^{1/2} \frac{dr}{r} = F \left( \frac{l(l+1)}{\omega^2} \right)^{1/2} = F \left( \frac{k_h}{\omega} \right)$$

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- Ridges for large  $l$ :  $F \sim (ld)^{-1}$ ,  $\omega^2 \sim k_h$
- Frequencies for small  $l$  and large  $n$

$$\int_{r_t}^{r_\odot} k_r dr \approx \frac{\omega}{\bar{c}} r_\odot - \left(l + \frac{1}{2}\right) \frac{\pi}{2} = \pi(n + \alpha)$$

$$\leadsto \omega \approx \left(n + \frac{l}{2} + \alpha + \frac{1}{4}\right) \frac{\pi \bar{c}}{r_\odot}$$

$$\leadsto \omega_{nl} \approx \omega_{n-1, l+2} \quad \text{and} \quad \omega_{nl} \approx \frac{1}{2}(\omega_{n-1, l+1} + \omega_{n+1, l+1})$$

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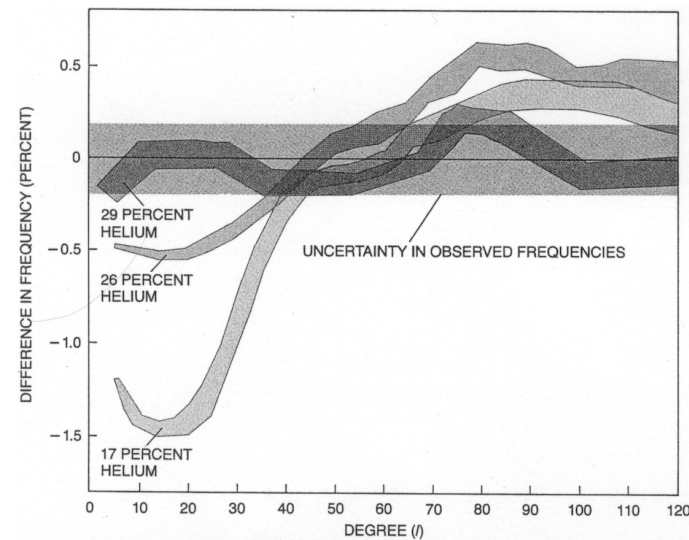
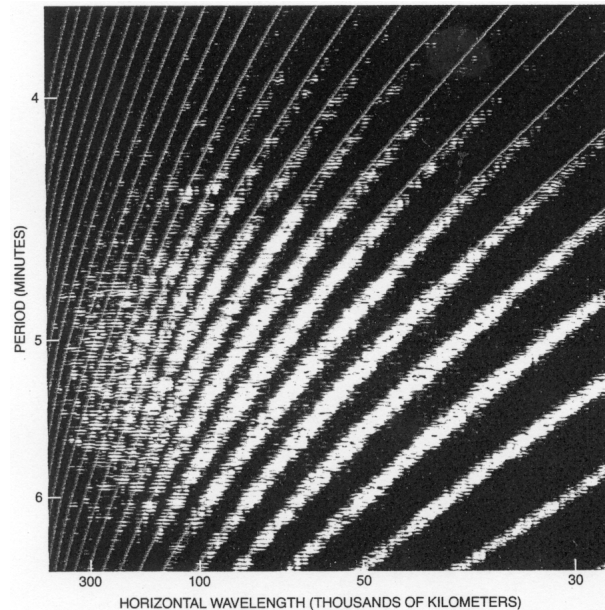
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## 4. Direct Methods

- Theoretical model of Sun  $\leadsto$  theoretical p-modes  
 $\leadsto$  comparison with observation  $\leadsto$  variation of solar model
- Good theory of p-modes necessary
- Often frequency differences  $\delta\omega_{nl} = \omega_{nl} - \omega_{n-1,l+2}$  compared



- Determination:
  - equation of state: electrostatic correction to perfect gas law
  - $Z \leadsto$  opacity  $\leadsto T(0) \leftrightarrow$  neutrino flux
  - $\gamma$
  - mixing length, depth of convection zone 200 000 km
- Confirmation of standard solar model

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## 5. Inversion: Sound Speed

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$$\text{Duvall's law} \quad \int_{r_t}^{r_\odot} \left( \frac{r^2}{c^2} - \frac{l(l+1)}{\omega^2} \right)^{1/2} \frac{dr}{r} = F \left( \frac{l(l+1)}{\omega^2} \right)^{1/2}$$

$$u = \frac{l(l+1)}{\omega^2}, \quad \xi = \left( \frac{r}{c} \right)^2, \quad \xi_t = \left( \frac{r_t}{c(r_t)} \right)^2 = \frac{l(l+1)}{\omega^2} = u, \quad \xi_\odot = \xi(r_\odot)$$

$$F(u) = \int_u^{\xi_\odot} (\xi - u)^{1/2} \underbrace{\frac{1}{r} \frac{dr}{d\xi}}_{d \ln r / d\xi} d\xi \quad \text{known from observation}$$

$$\frac{dF}{du} = -\frac{1}{2} \int_u^{\xi_\odot} \frac{d \ln r / d\xi}{(\xi - u)^{1/2}} d\xi, \quad \text{Abel's integral equation}$$

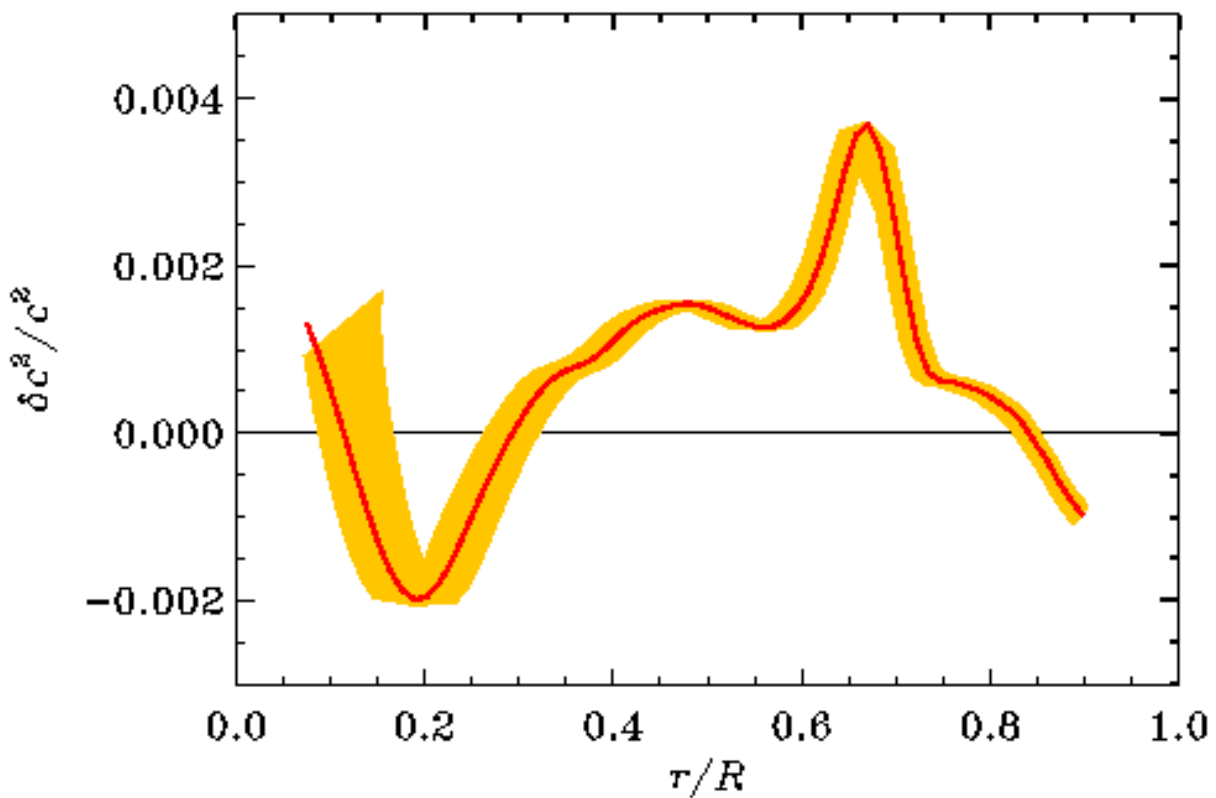
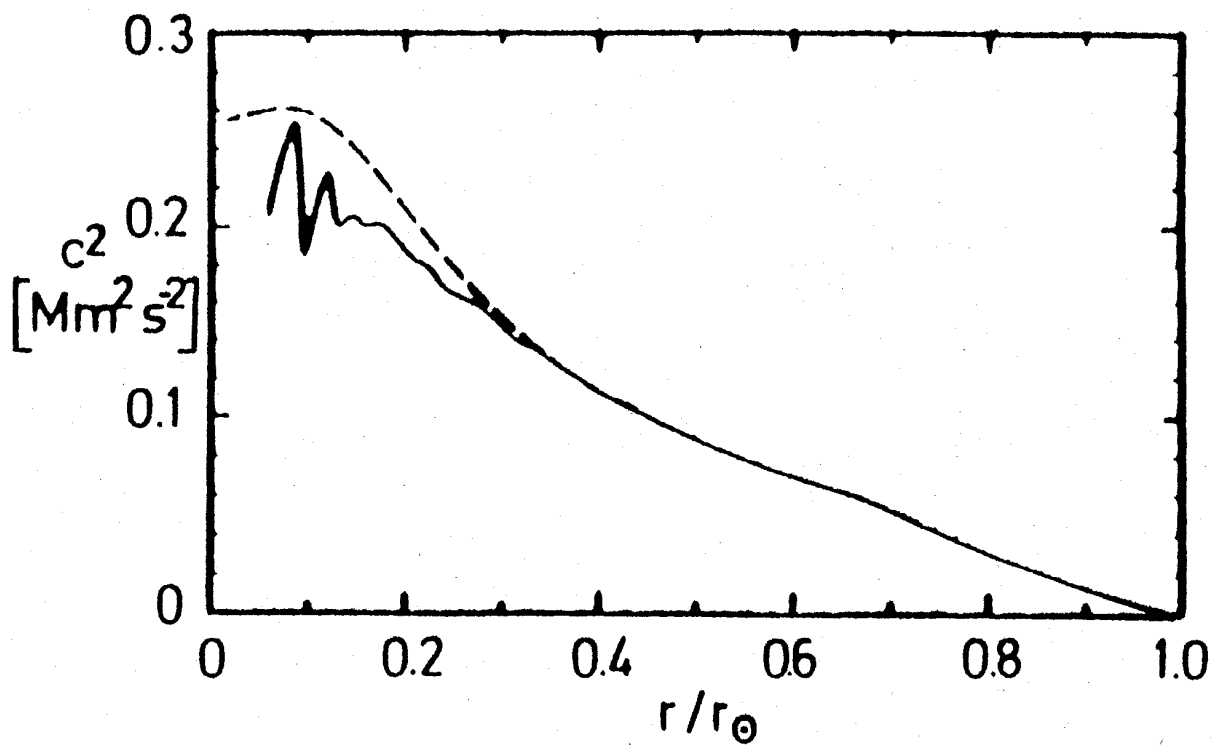
$$\text{solution} \quad \ln r(\xi) - \ln r(\xi_\odot) = -\frac{2}{\pi} \int_{\xi_\odot}^{\xi} \frac{dF/du}{(u - \xi)^{1/2}} du$$

$$r = r_\odot \exp \left( -\frac{2}{\pi} \int_{\xi_\odot}^{\xi} \frac{dF/du}{(u - \xi)^{1/2}} du \right) = r(\xi)$$

$$r(\xi) \rightsquigarrow \xi(r) \rightsquigarrow c(r) \rightsquigarrow T(r)$$

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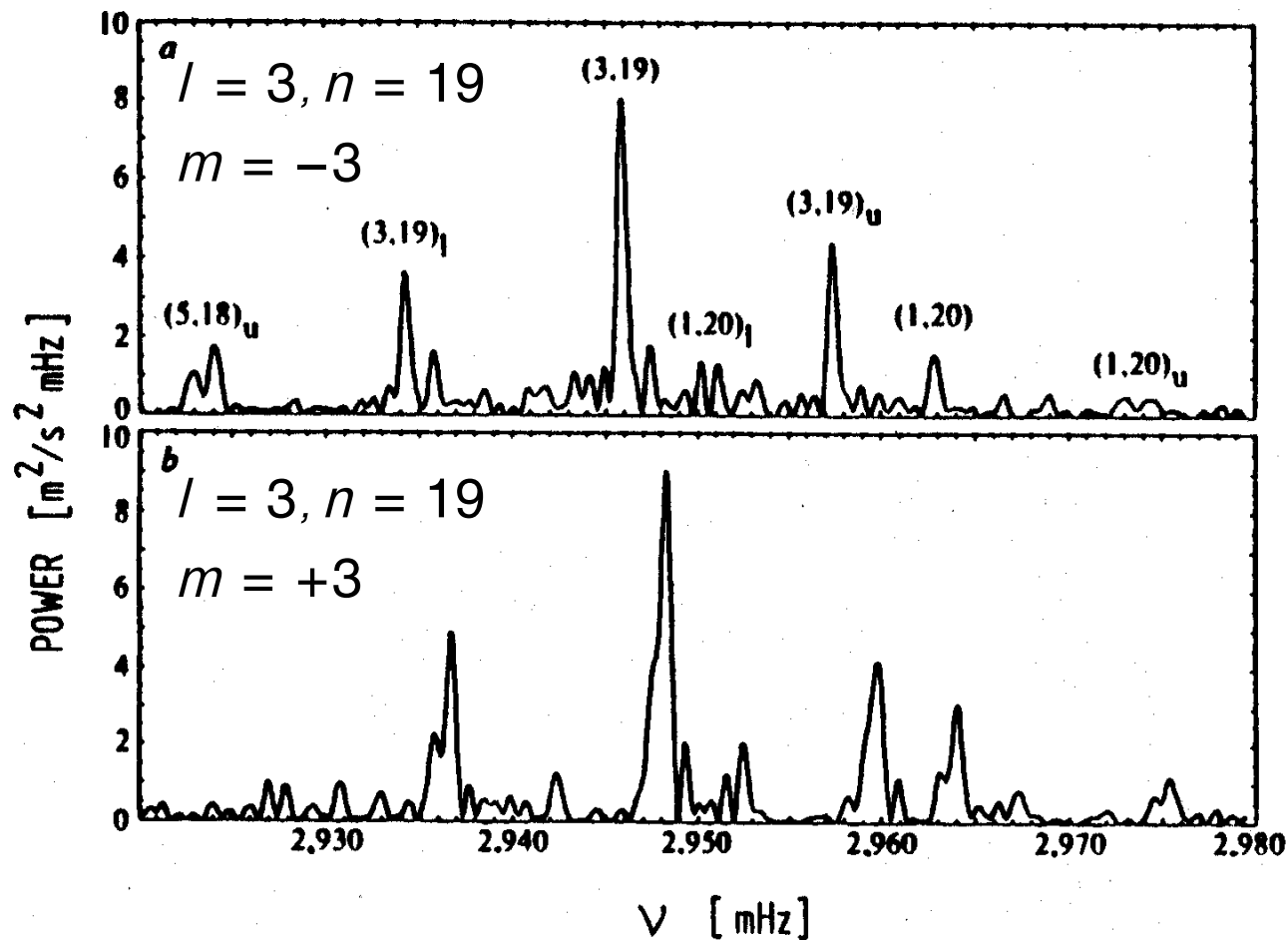
## 6. Internal Rotation

- No rotation:  $\omega_{nlm}$  not dependent on azimuthal wavenumber  $m \leadsto$  degeneration
- With rotation no degeneration
  - assumption for simplicity: rigid rotation
  - $\omega_{nlm} = \omega_{nl0} \pm m\Omega$ ,  $m = -l, \dots, +l$
  - standing wave = superposition of two identical travelling waves in positive and negative azimuthal direction
  - rotation, Doppler effect, frequency shift and splitting
- $\omega_{nlm} - \omega_{nl,m-1} = \Omega$ ,  $\frac{\Omega}{\omega_{nl0}} \approx \frac{5 \text{ min}}{30 \text{ d}} \approx 10^{-4}$
- Observing time  $T \geq 2\pi/\Omega \approx 30 \text{ d}$  to resolve  $\Omega$  with  $\geq 2$  measurements per period (Nyquist)
- Duvall and Harvey (1984)  
 $l = 3, n = 19, m = +3, -3$

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- Side peaks due to night gaps  
 $\curvearrowright$  south pole, GONG, SOHO
- Frequency splitting weighted average of  $\Omega(r, \theta)$

$$\Delta\omega_{n/m} = \int K_{n/m}(r, \theta)\Omega(r, \theta)rdrd\theta$$

- Kernel  $K_{n/m}$  given, i.e.  $|\xi|^2$  of eigenfunction
- $n$  of little influence,  $l$  determines contribution in depth,  $m$  in latitude

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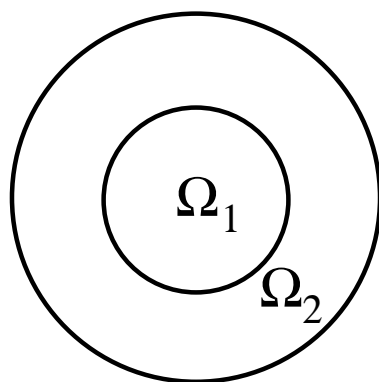
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- Various methods to extract  $\Omega$ , here optimal kernels  
(Backus & Gilbert, 1970):

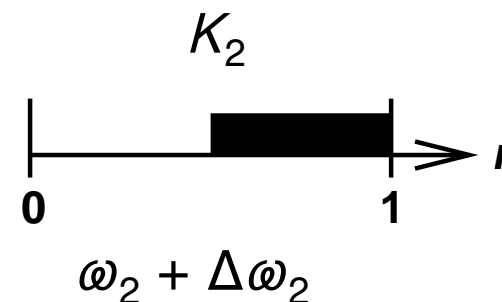
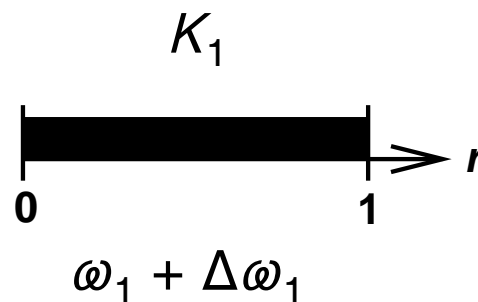
$$\sum_i a_i(r_0) K_i(r) = \delta(r - r_0), \text{ combination of kernels to } \delta \text{ functions}$$

$$\sum_i a_i(r_0) \Delta\omega_i = \int \delta(r - r_0) \Omega(r) d^3r = \Omega(r_0)$$

- Example:

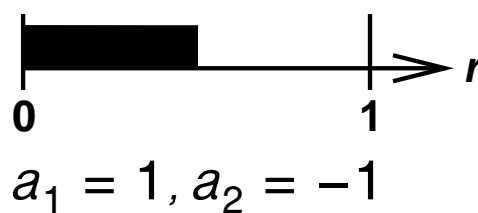


2 waves

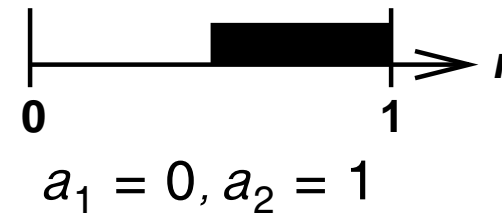


optimal kernels:

$$\tilde{K}_1 = K_1 - K_2$$



$$\tilde{K}_2 = K_2$$



$$\Omega_1 = \Delta\omega_1 - \Delta\omega_2, \quad \Omega_2 = \Delta\omega_2$$

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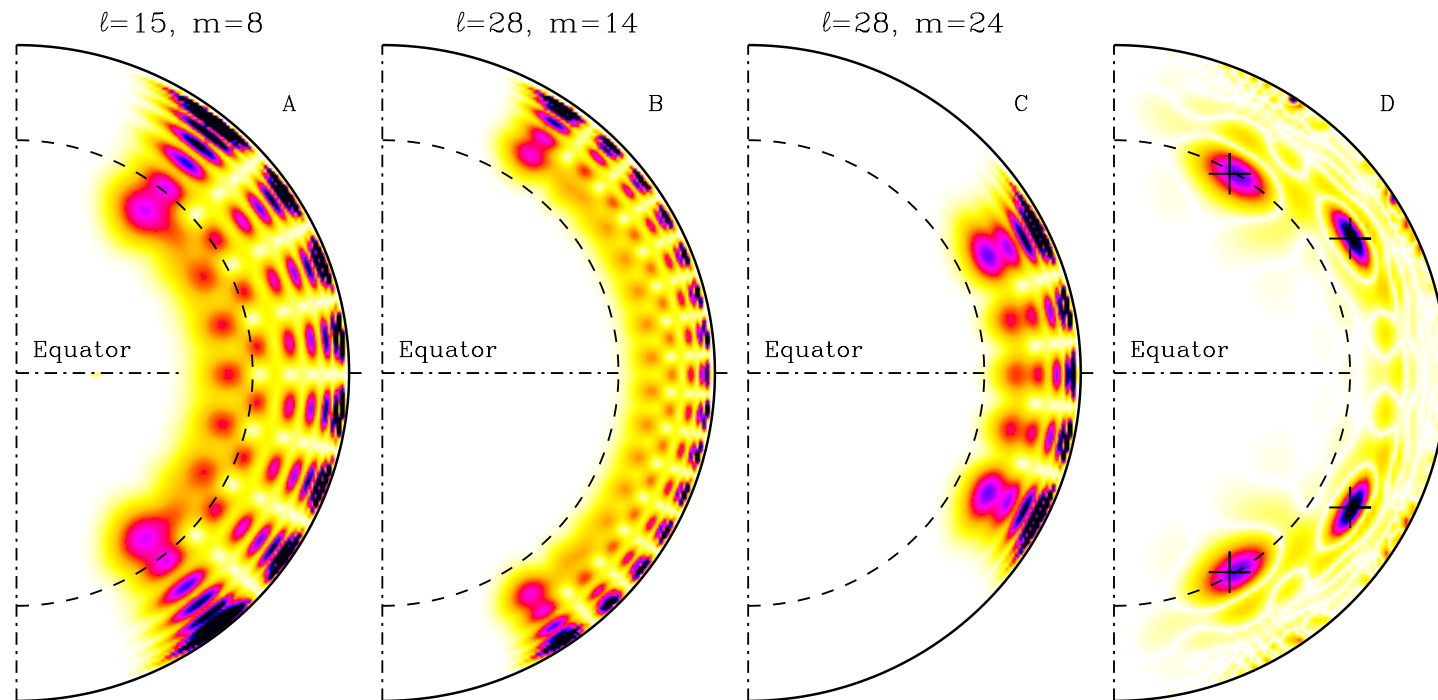
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$$\Delta\omega_1 = \int K_1 \Omega dr = K_1 \left( \Omega_1 \int_0^{1/2} dr + \Omega_2 \int_{1/2}^1 dr \right) = \frac{K_1}{2} (\Omega_1 + \Omega_2)$$

$$\Delta\omega_2 = \int K_2 \Omega dr = \frac{1}{2} K_2 \Omega_2, \quad K_1 = K_2 = 2$$

- Solar p-mode kernels



- Such each location in Sun addressable  $\leadsto \Omega(r, \theta)$

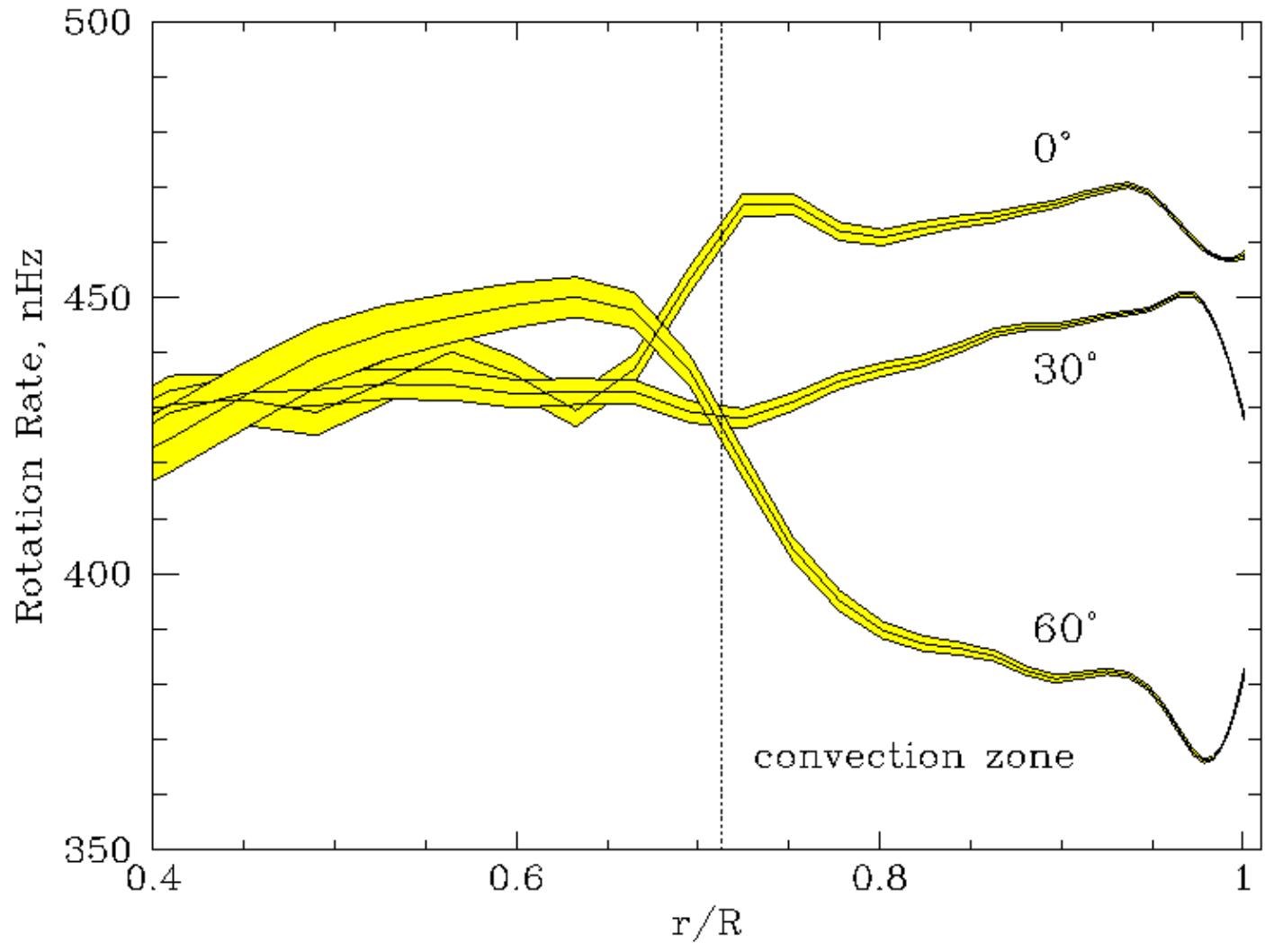
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• Result



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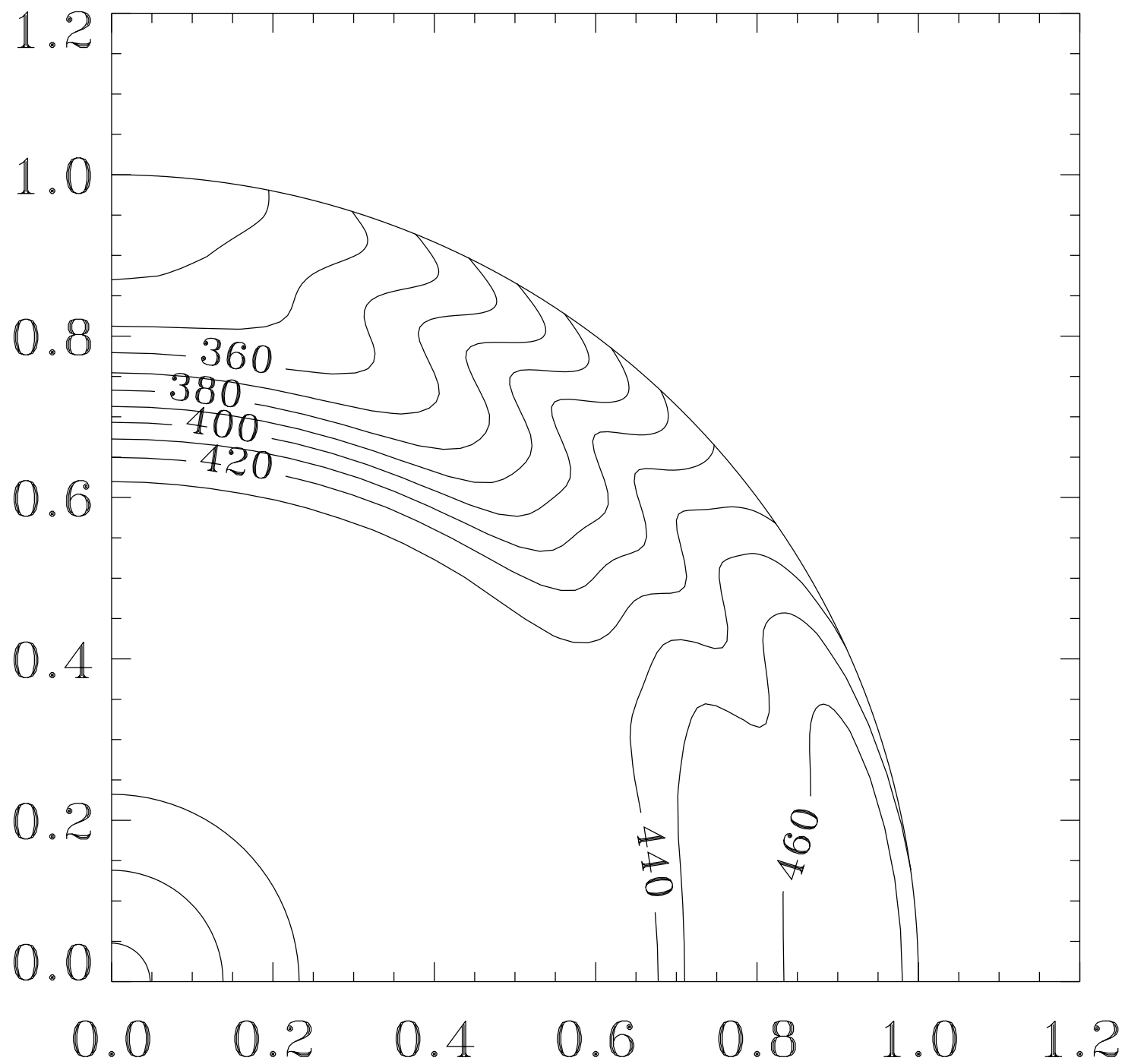
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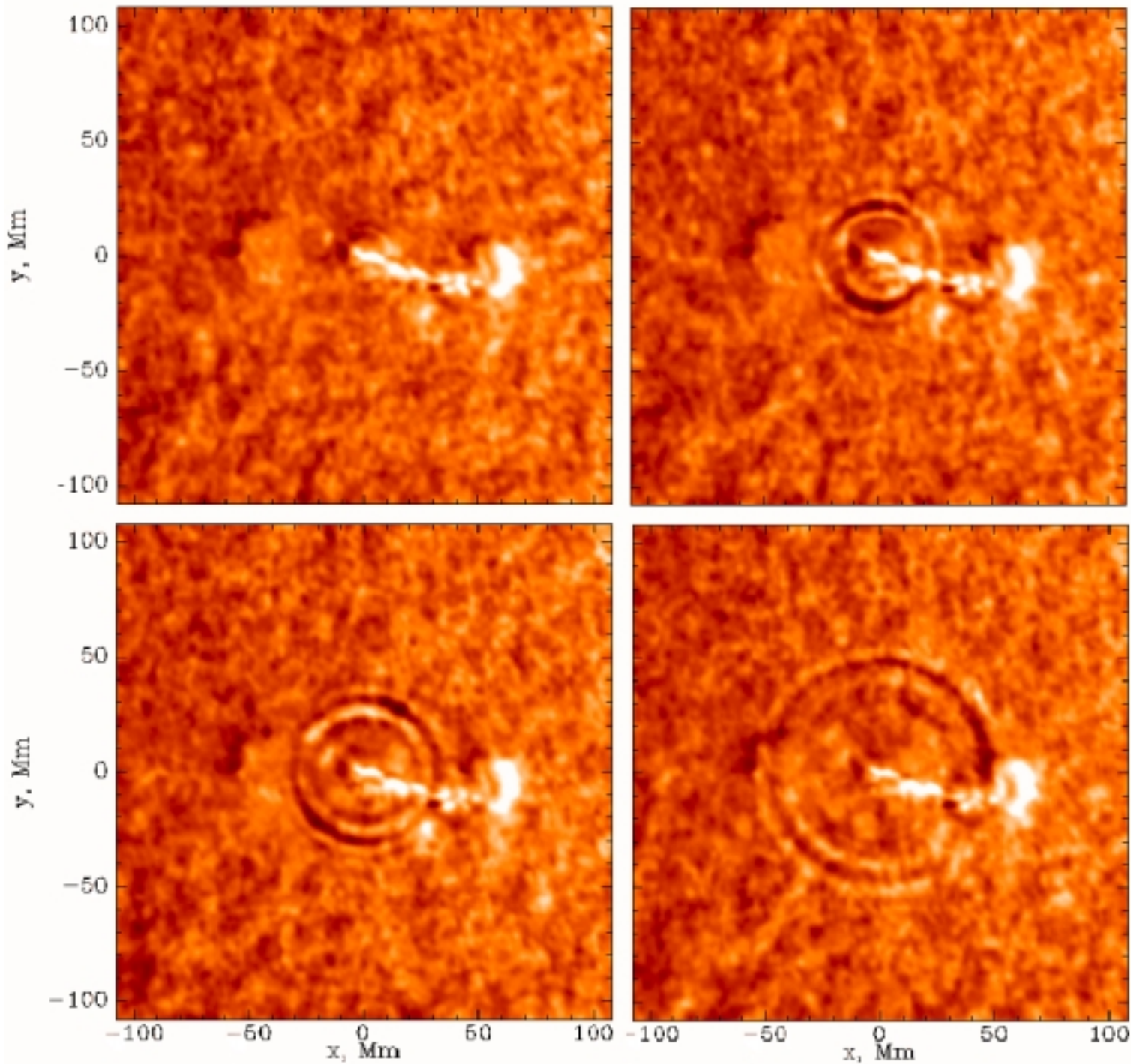


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# 7. Time-Distance Helioseismology



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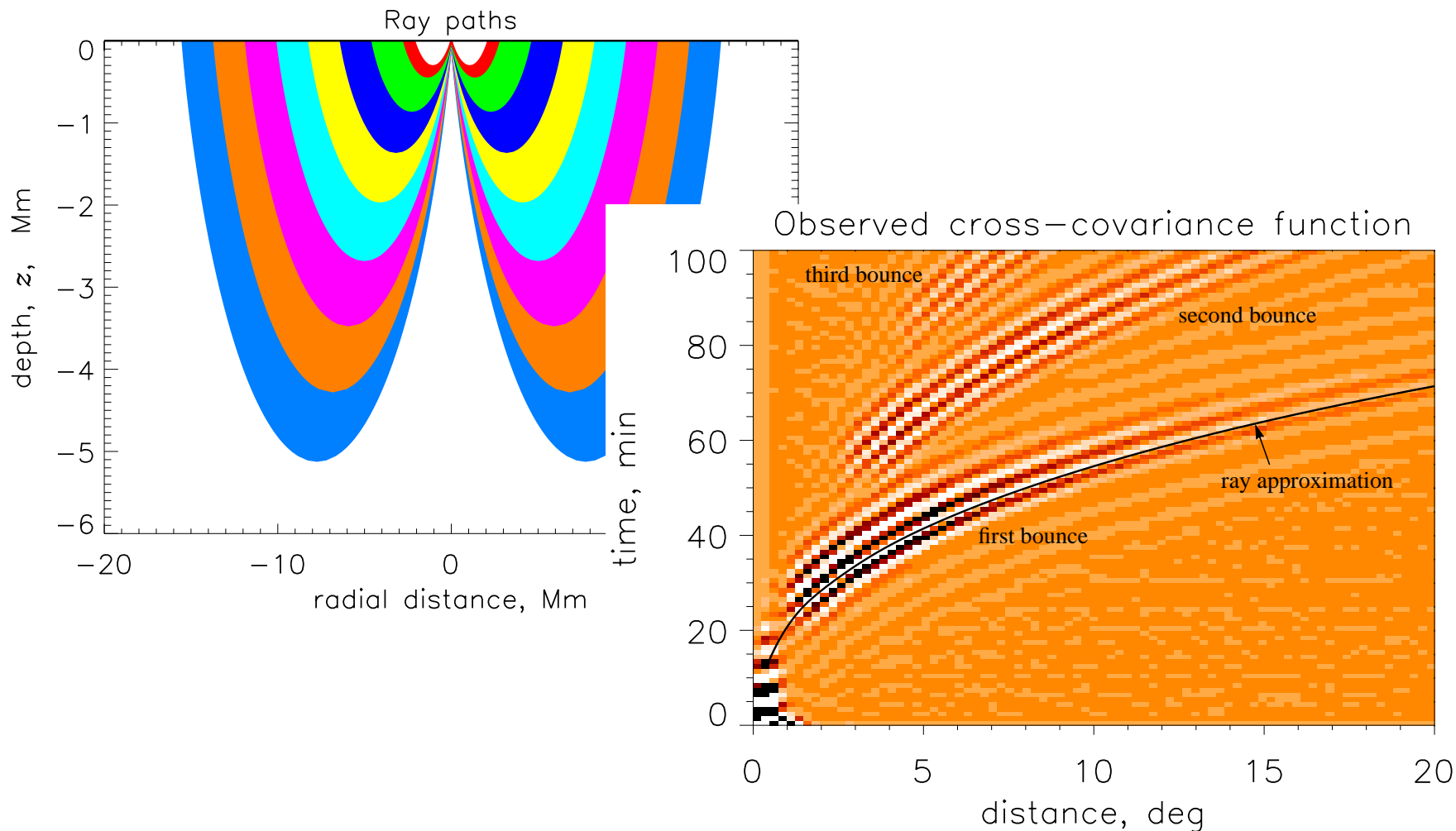
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- Ray paths: horizontal distance  $\approx \pi \cdot \text{depth}$ , length  $\approx 4 \cdot \text{depth}$
- Cross-correlation of time series at two separate points as function of travel time
- Speed up in hotter regions
- Speed up for wave in flow direction, slow down in reciprocal direction
- Magnetic field  $\curvearrowright$  anisotropy in travel time

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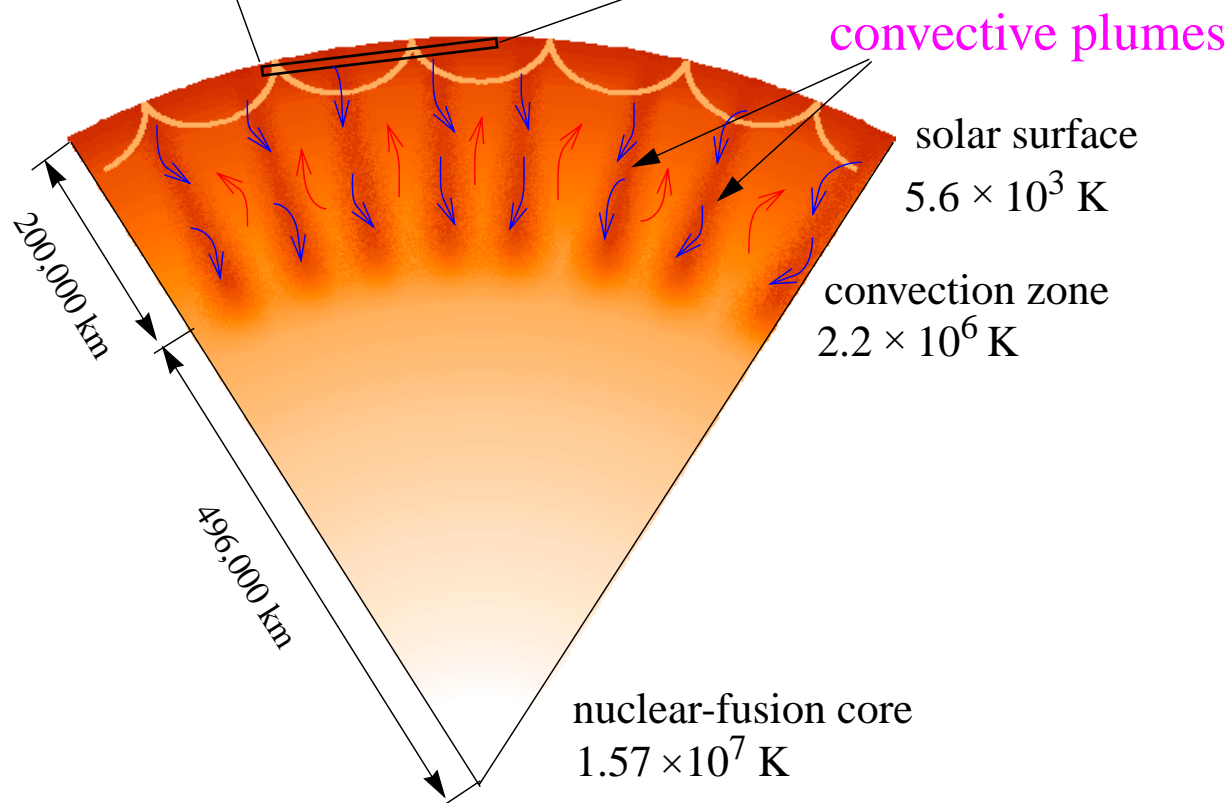
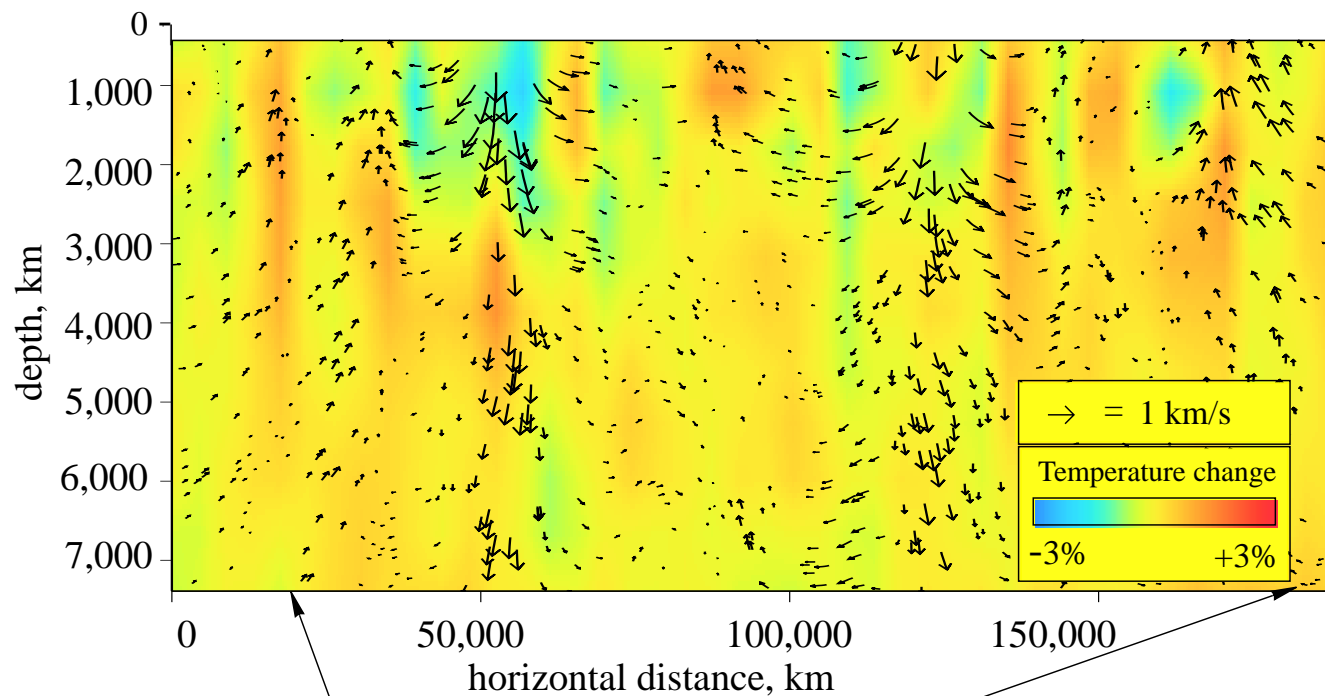
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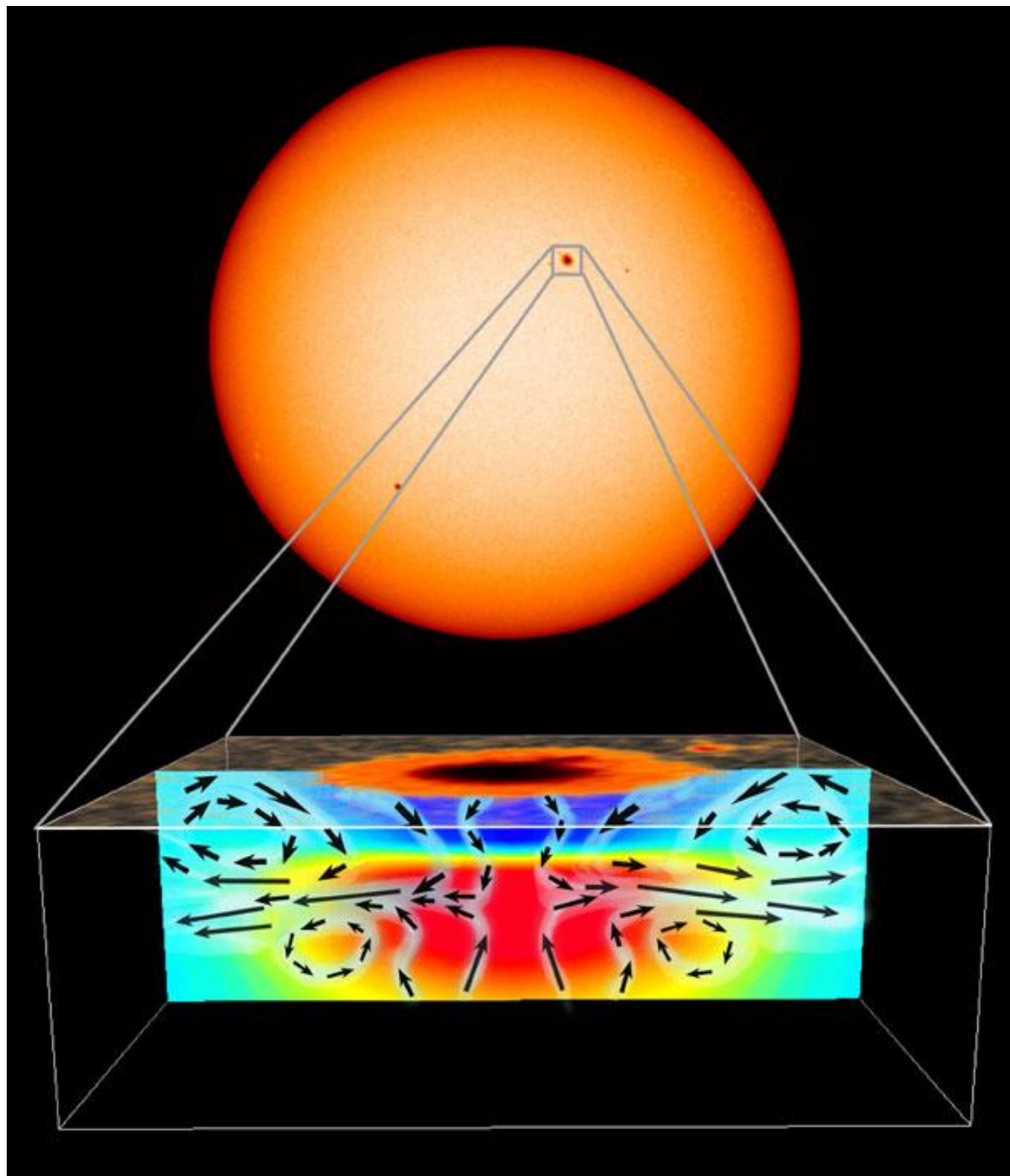


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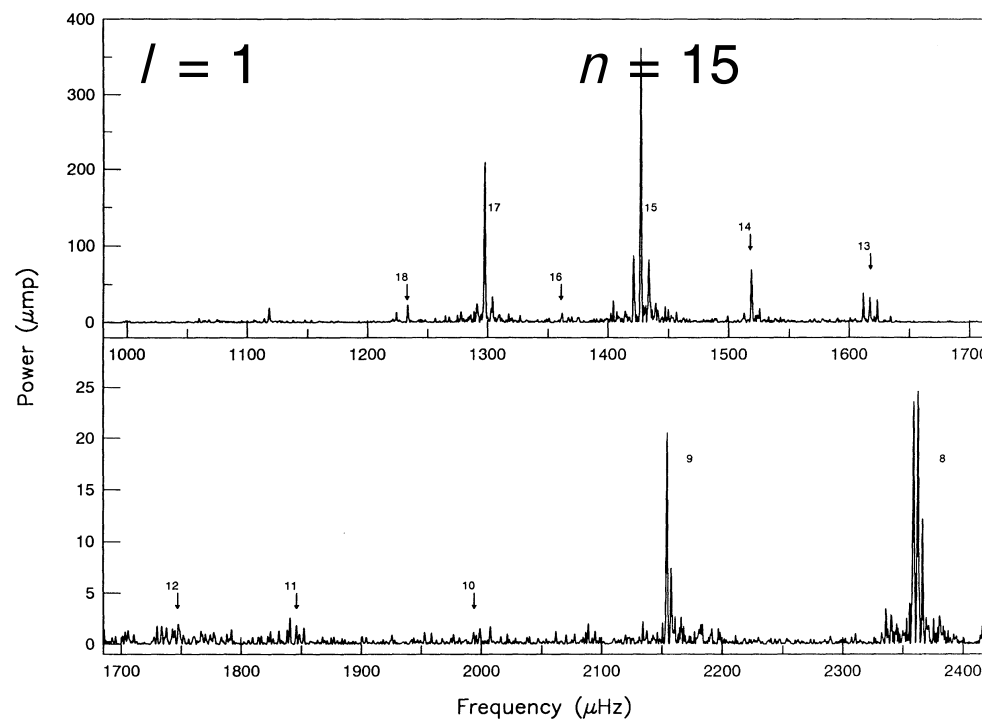
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## 8. Asteroseismology

- Whole solar disk, small  $l$   
spectral line shift – Doppler effect
- Oscillations in intensity  
→ asteroseismology by photometry  
good photometers  $\Delta m \sim 0.01$   
long time series (weeks): Whole Earth Telescope (WET)
- Spectrum of white dwarf GD 358



→ g-modes

→ determination of mass, rotation, magnetic field, ...

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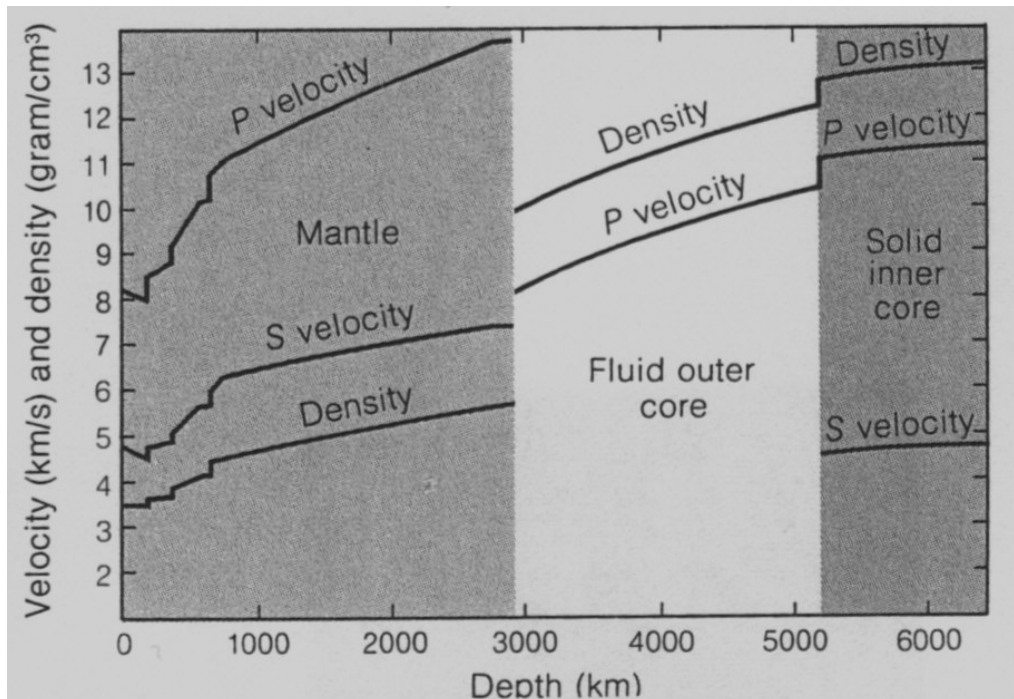
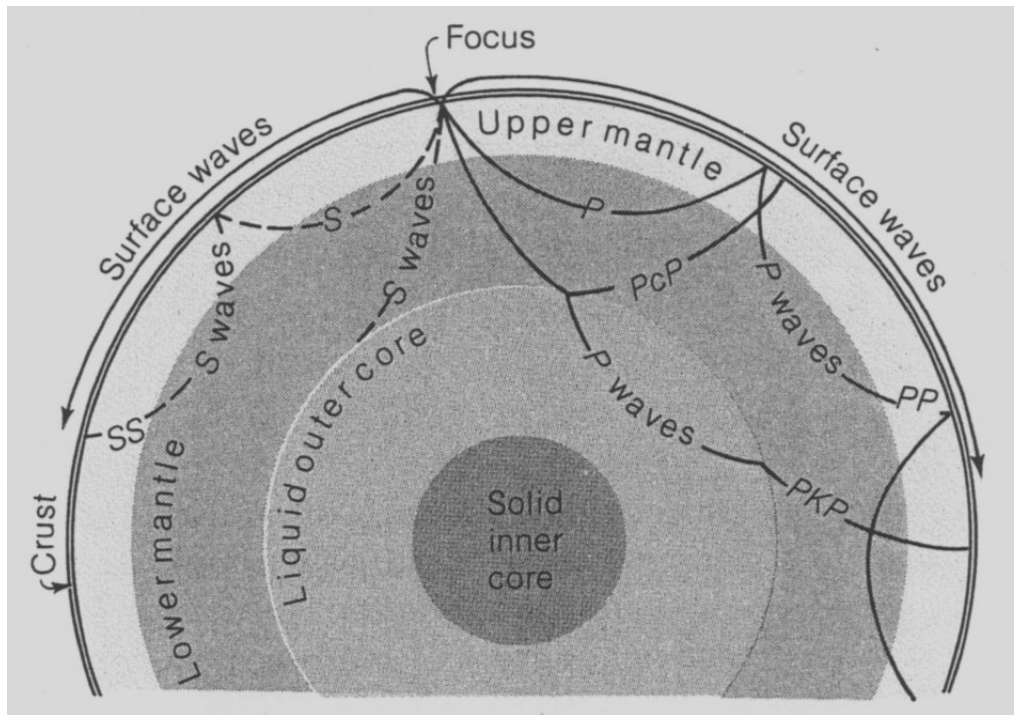
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