

The microstate of the solar corona and the solar wind

- Radial gradients of kinetic temperatures
- Velocity distribution functions
- Ion composition and suprathermal electrons
- Coulomb collisions in the solar wind
- Waves and plasma microinstabilities
- Diffusion and wave-particle interaction

Length scales in the solar wind

Macrostructure - fluid scales

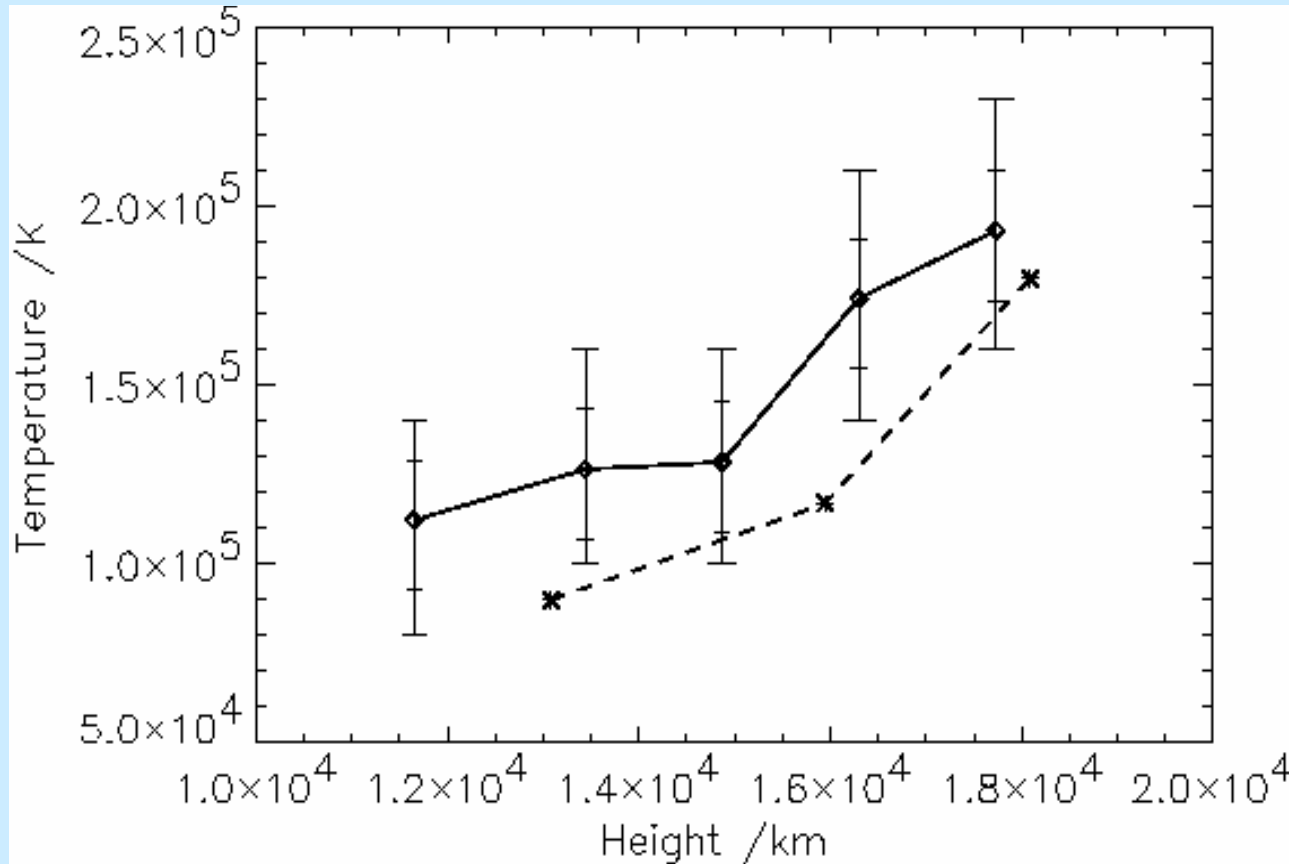
- Heliocentric distance: r 150 Gm (1AU)
- Solar radius: R_s 696000 km (215 R_s)
- Alfvén waves: λ 30 - 100 Mm

Microstructure - kinetic scales

- Coulomb free path: l ~ 0.1 - 10 AU
- Ion inertial length: V_A/Ω_p (c/ω_p) ~ 100 km
- Ion gyroradius: r_L ~ 50 km
- Debye length: λ_D ~ 10 m
- Helios spacecraft: d ~ 3 m

Microscales vary with solar distance!

Proton temperature at coronal base



SUMER/SOHO

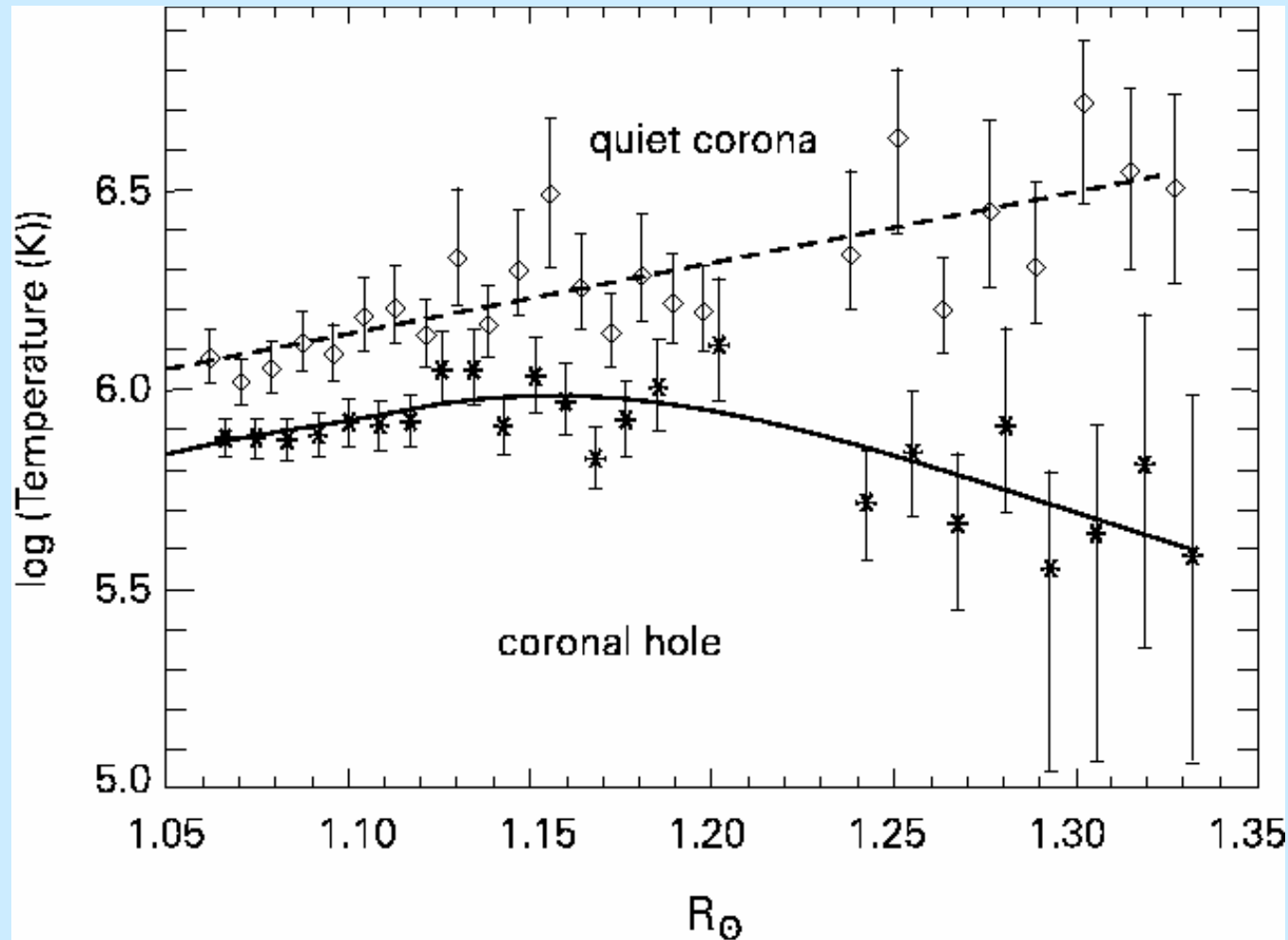
Hydrogen
Lyman series

Off-limb
transition
region at the
base of north
polar CH

Marsch et al., A&A,
359, 381, 2000

Charge-exchange equilibrium: $T_H = T_p$
Turbulence broadening: $\xi = 30 \text{ km s}^{-1}$

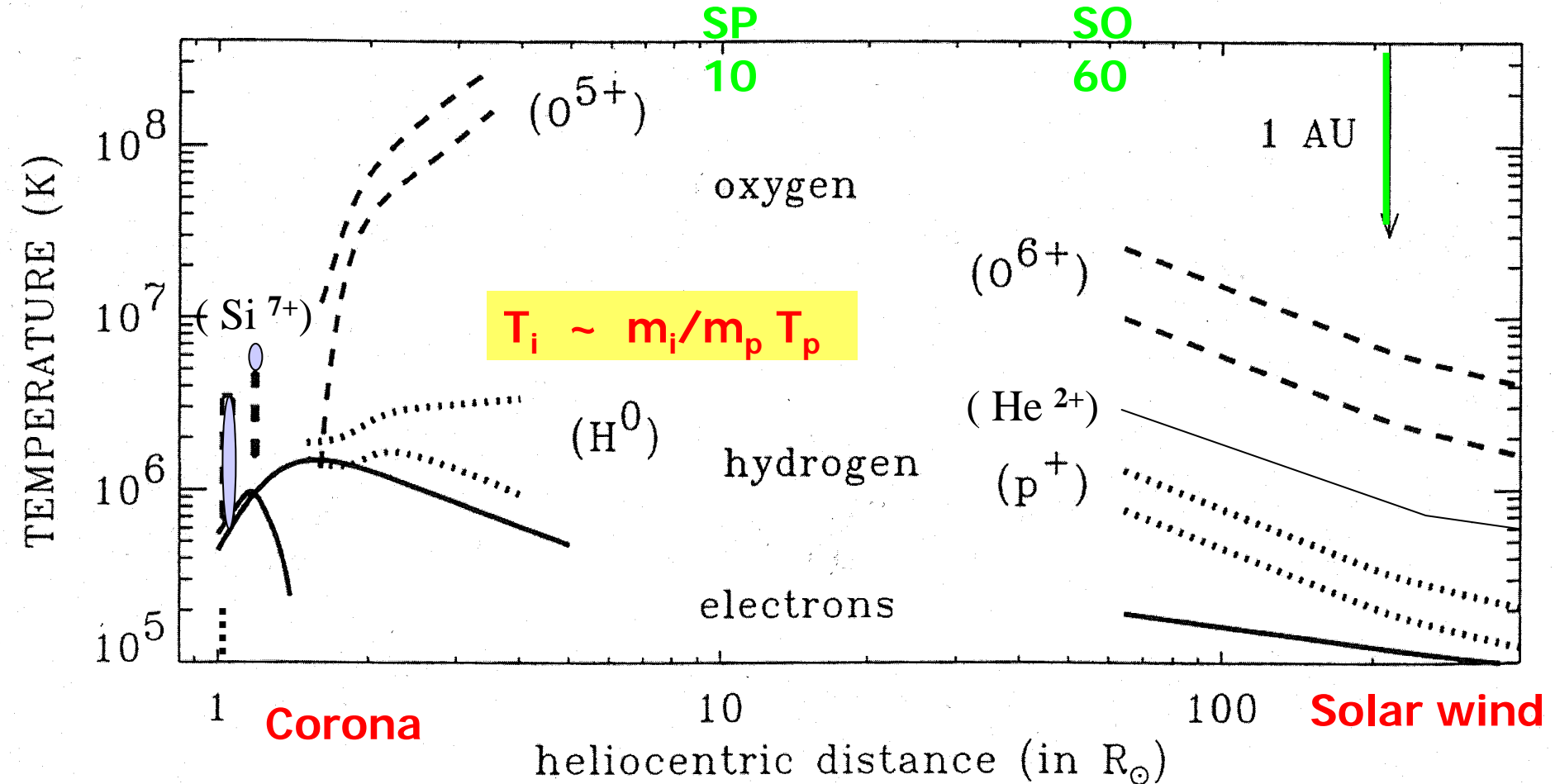
Electron temperature in the corona



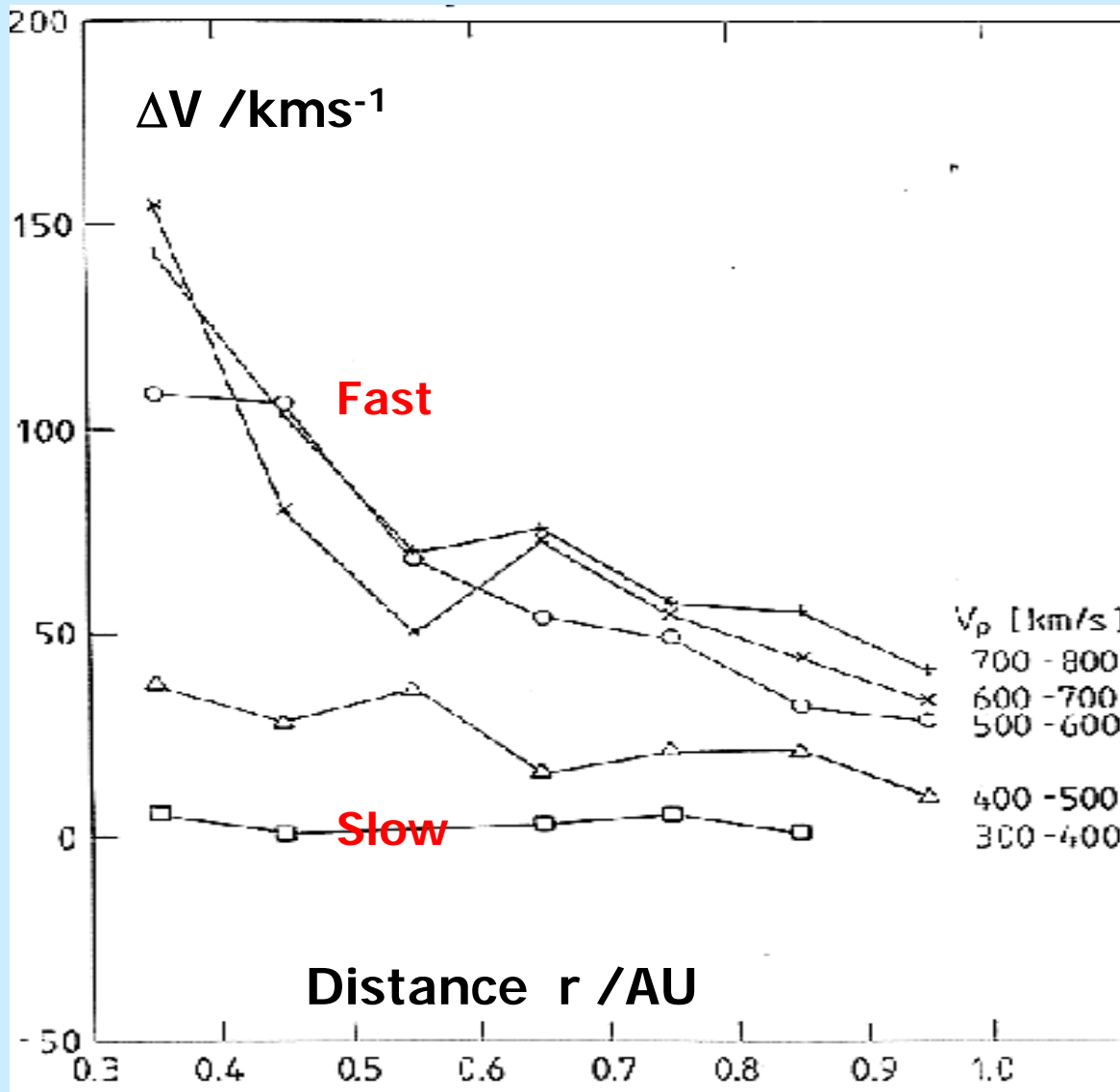
Streamer
belt, closed

Coronal hole,
open
magnetically

Temperature profiles in the corona and fast solar wind



Ion differential streaming



- Helios:

Alpha particles are faster than the protons!

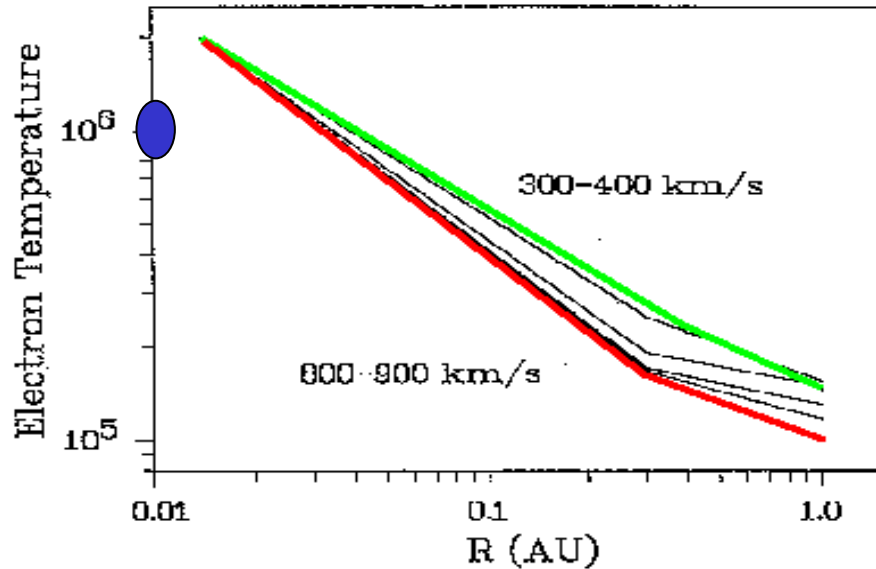
- In fast streams the differential velocity $\Delta V \leq V_A$

- Ulysses:

Heavy ions travel at alpha-particle speed!

Proton and electron temperatures

Electrons
are cool!

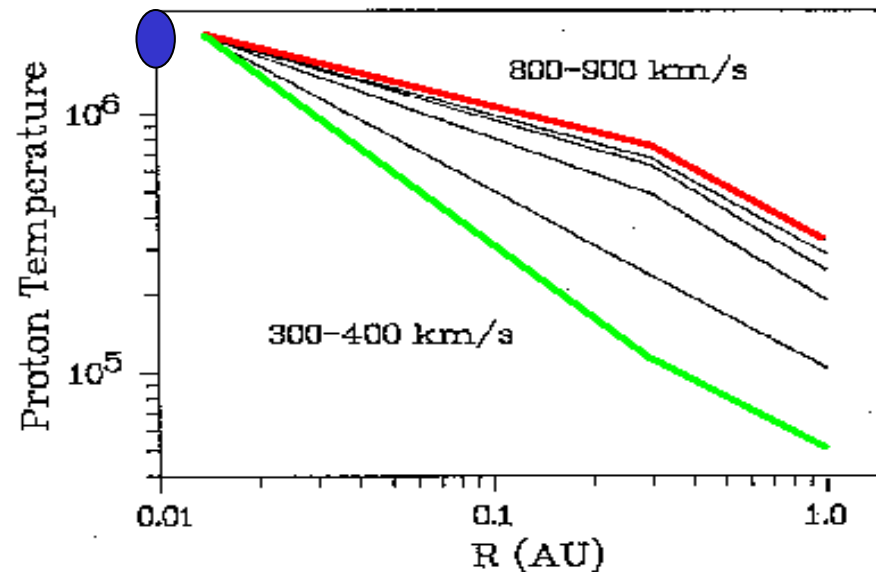


slow wind



fast wind

Protons
are hot!

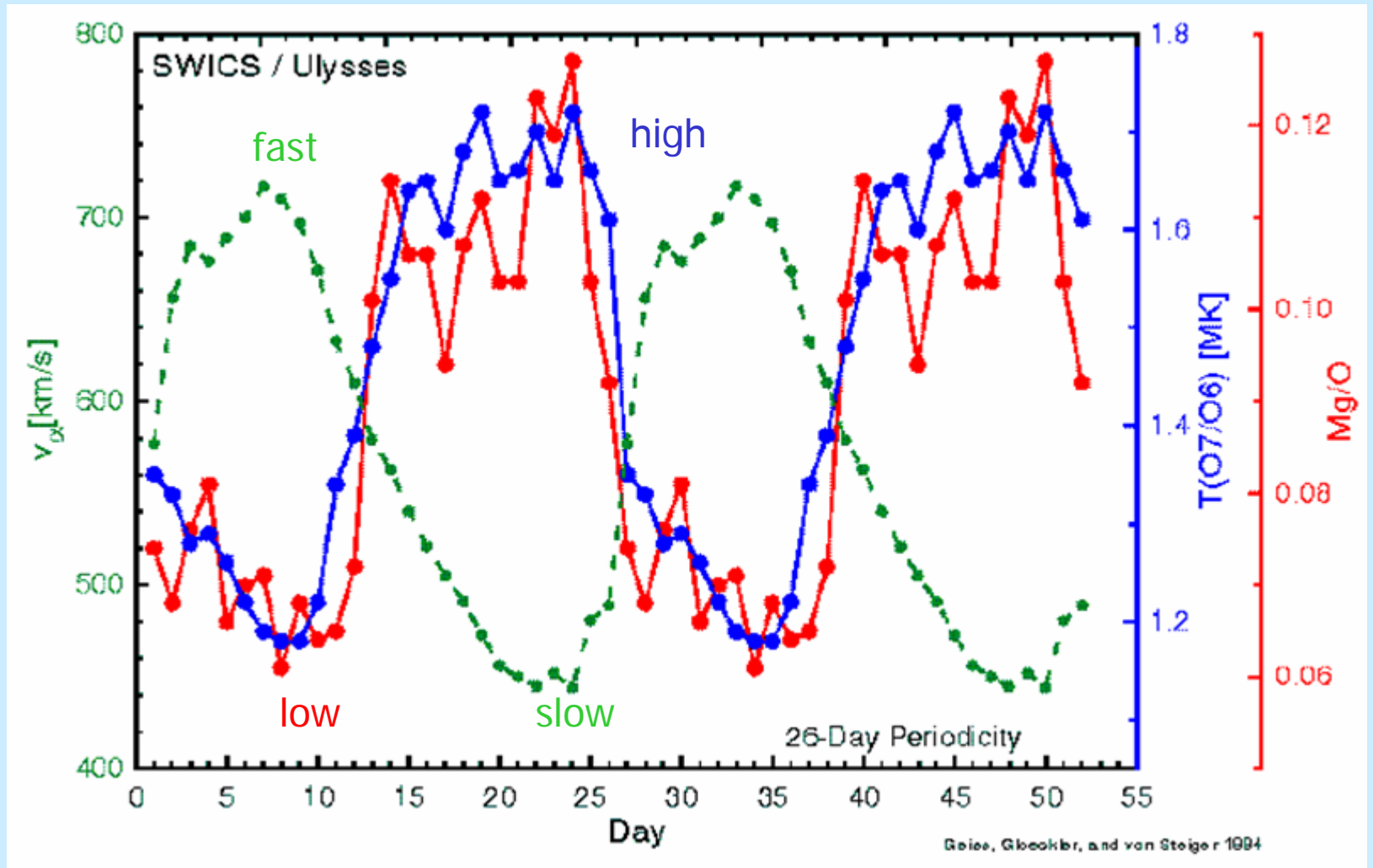


fast wind

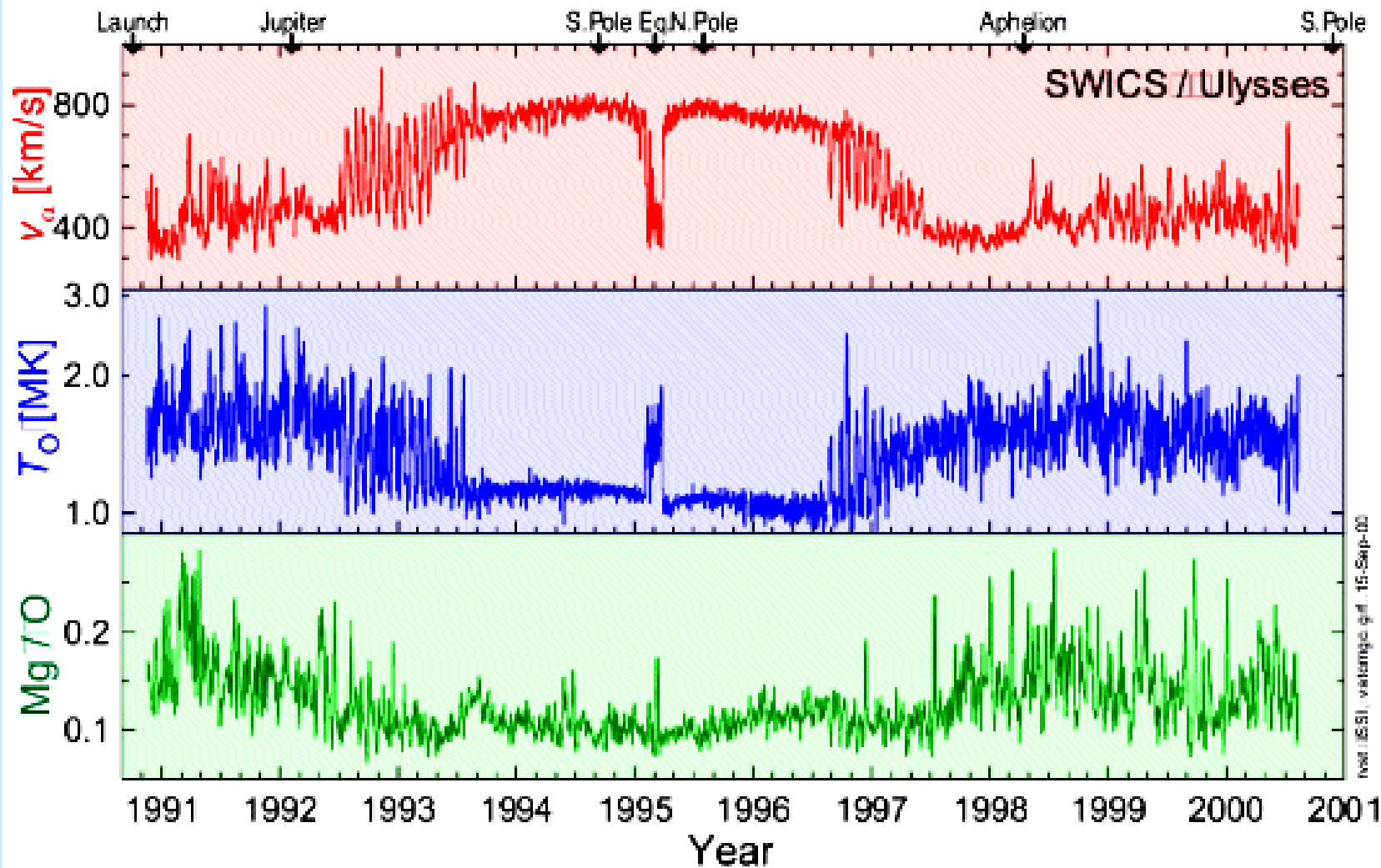


slow wind

Oxygen freeze-in temperature



Correlations between wind speed and corona temperature



Velocity distribution functions

Statistical description: $f_j(\mathbf{x}, \mathbf{v}, t) d^3x d^3v,$

gives the probability to find a particle of species j with a velocity \mathbf{v} at location \mathbf{x} at time t in the 6-dimensional phase space.

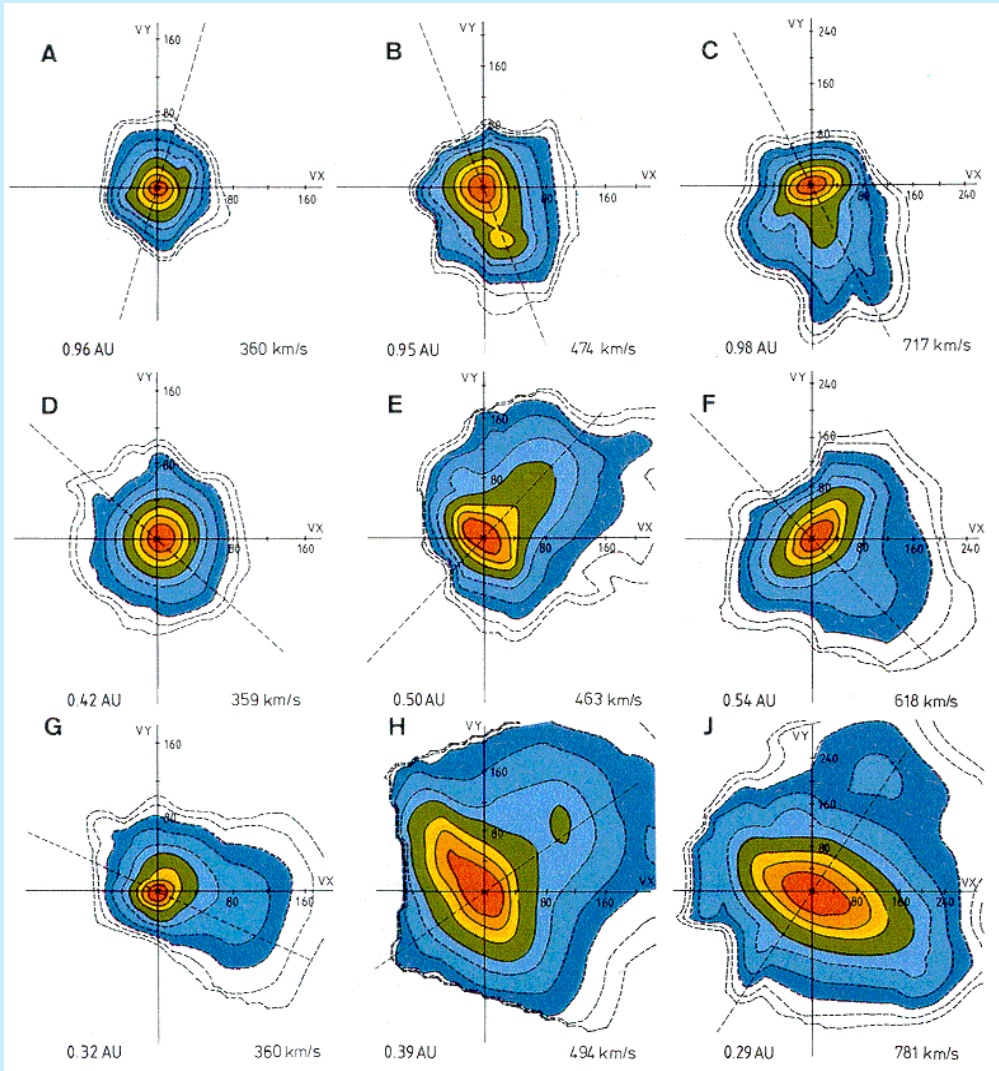
Local thermodynamic equilibrium:

$$f_j^M(\mathbf{x}, \mathbf{v}, t) = n_j (2\pi v_j)^{-3/2} \exp[-(\mathbf{v} - \mathbf{U}_j)^2 / (2v_j)^2],$$

with number density, n_j , thermal speed, $v_j = (k_B T_j / m_j)^{1/2}$, temperature T_j , and bulk velocity, \mathbf{U}_j , of species j .

Dynamics in phase space: Vlasov/Boltzmann kinetic equation

Proton velocity distributions



- Temperature anisotropies
- Ion beams
- Plasma instabilities
- Interplanetary heating

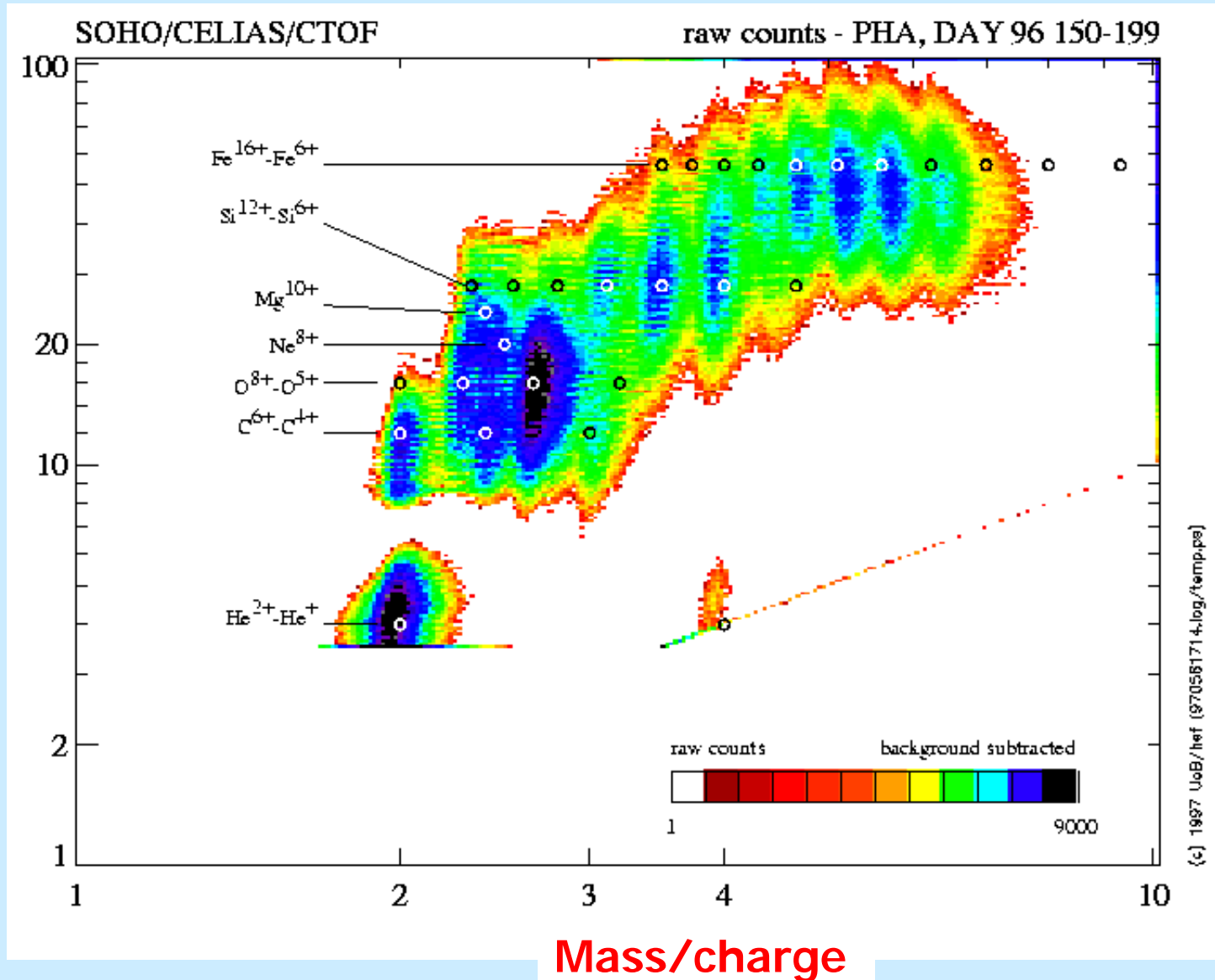
Plasma measurements
made at 10 s resolution
(> 0.29 AU from the Sun)

Helios

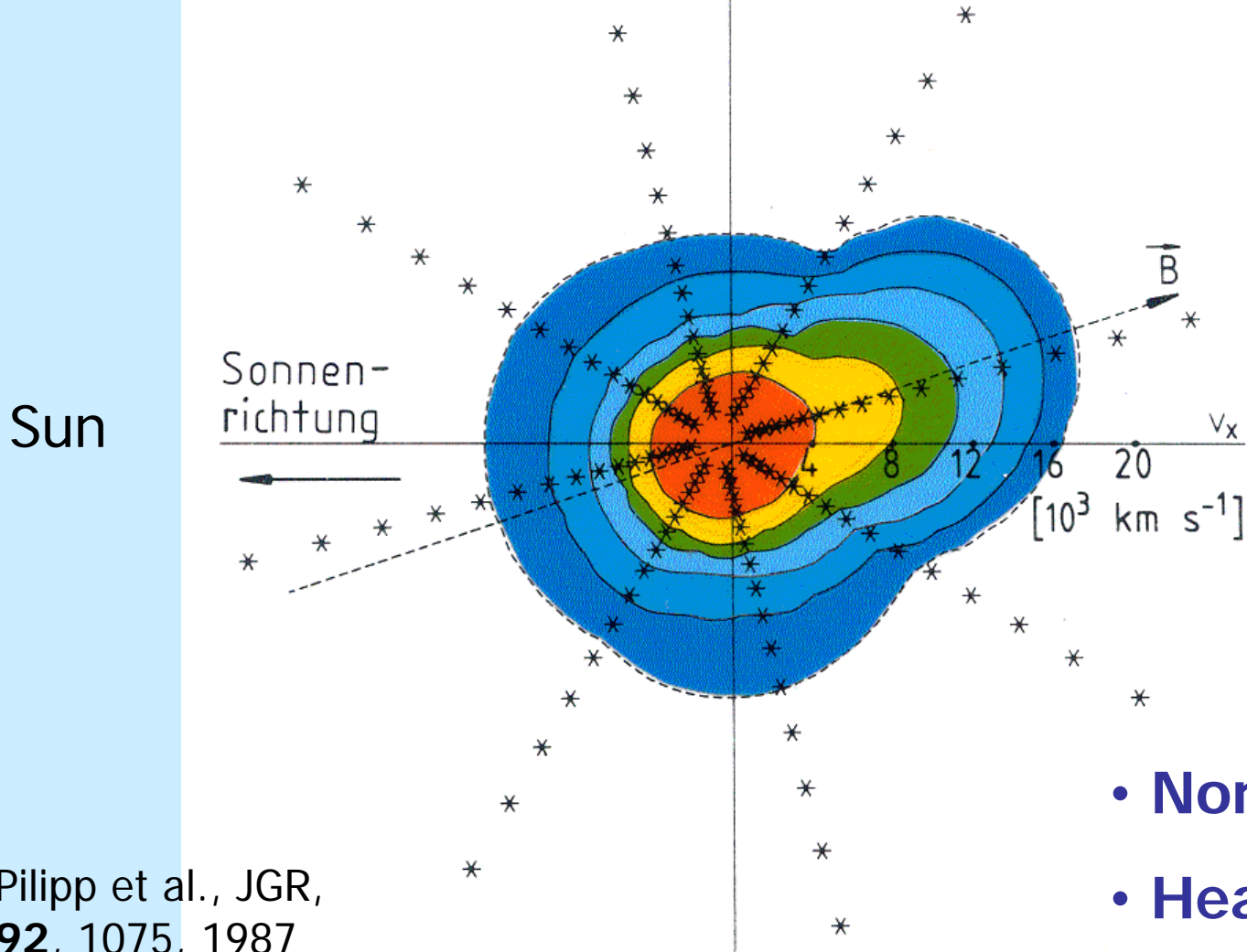
Marsch et al., JGR, **87**, 52, 1982

Ion composition of the solar wind

Ion
mass



Electron velocity distribution function



Helios

- Non-Maxwellian
- Heat flux tail

Kinetic processes in the solar corona and solar wind I

- Plasma is multi-component and nonuniform
 - **complexity**
- Plasma has low density
 - **deviations from local thermal equilibrium**
 - **suprathermal particles (electron strahl)**
 - **global boundaries are reflected locally**

Problem: Thermodynamics of the plasma, which is far from equilibrium.....

Electron energy spectrum

IMP spacecraft

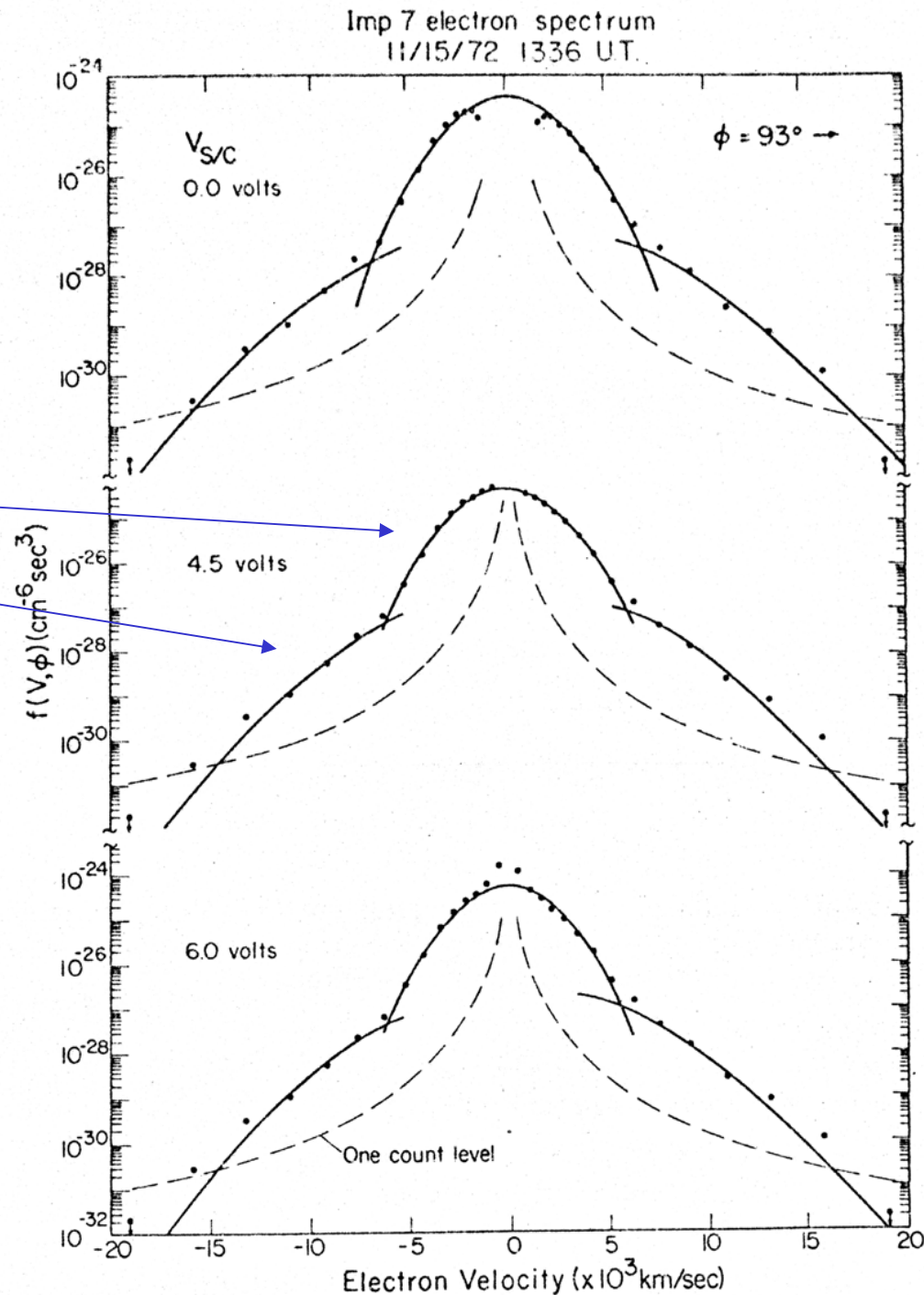
Two solar wind populations:

- Core (96%)
- Halo (4%)

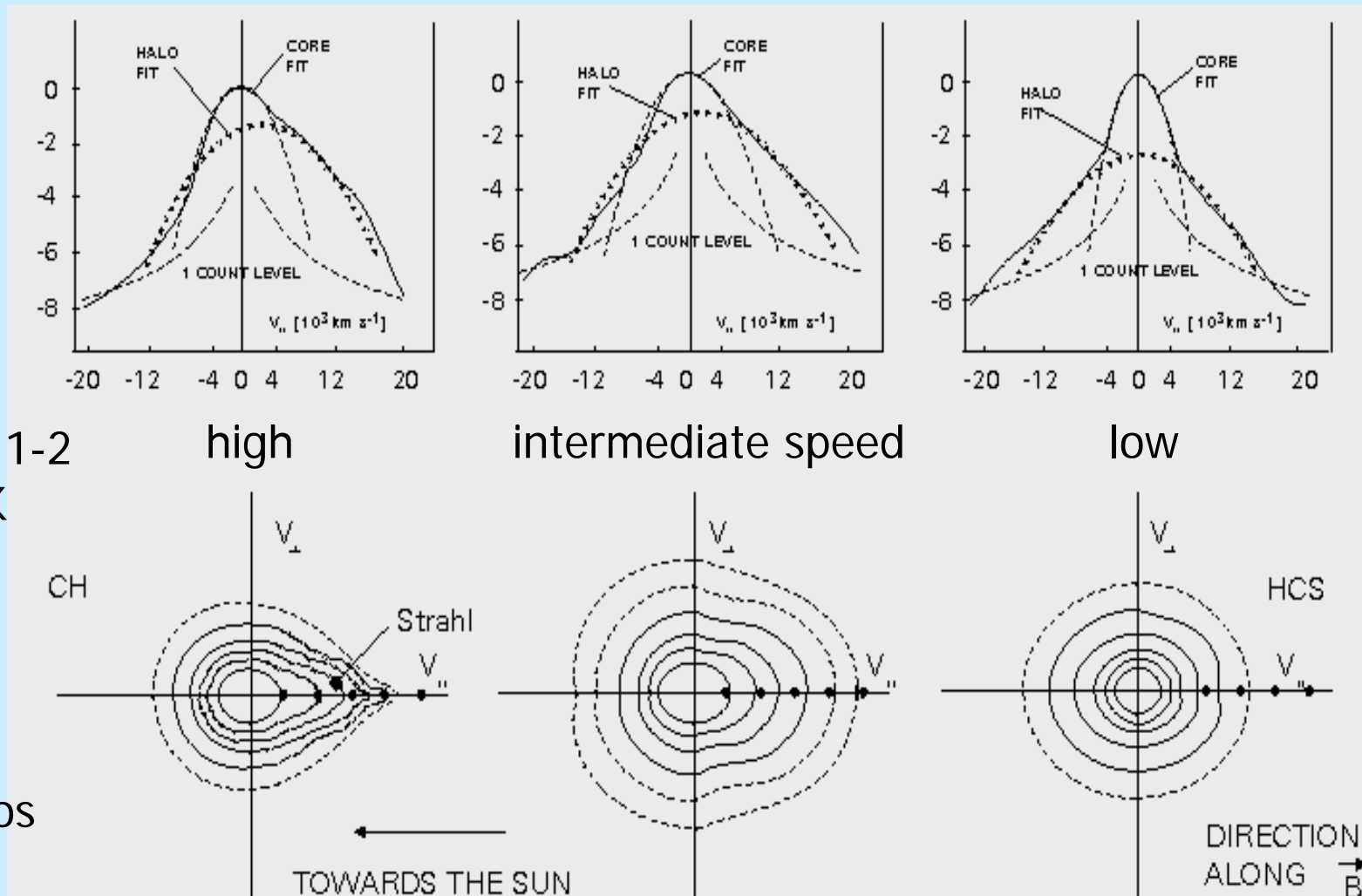
Core: local, collisional, **bound** by electrostatic potential

Halo: global, collisionless, **free** to escape (exospheric)

Feldman et al., JGR, **80**, 4181, 1975



Electron velocity distributions



$T_e = 1-2$
 10^5 K

high

intermediate speed

low

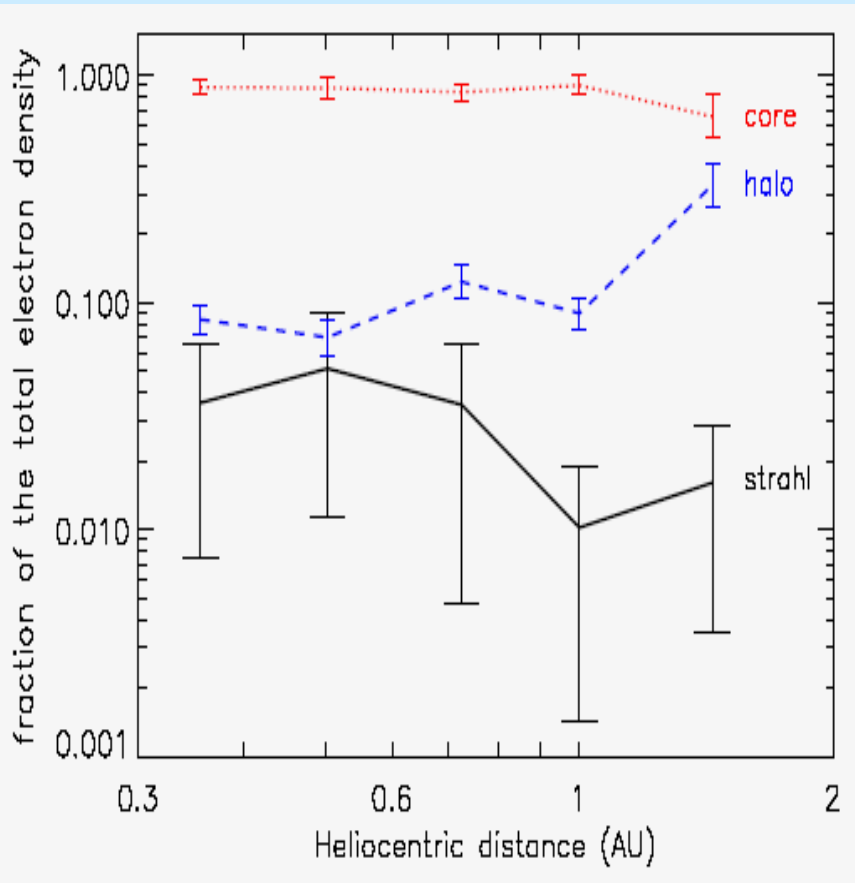
Helios

Pilipp et al., JGR,
92, 1075, 1987

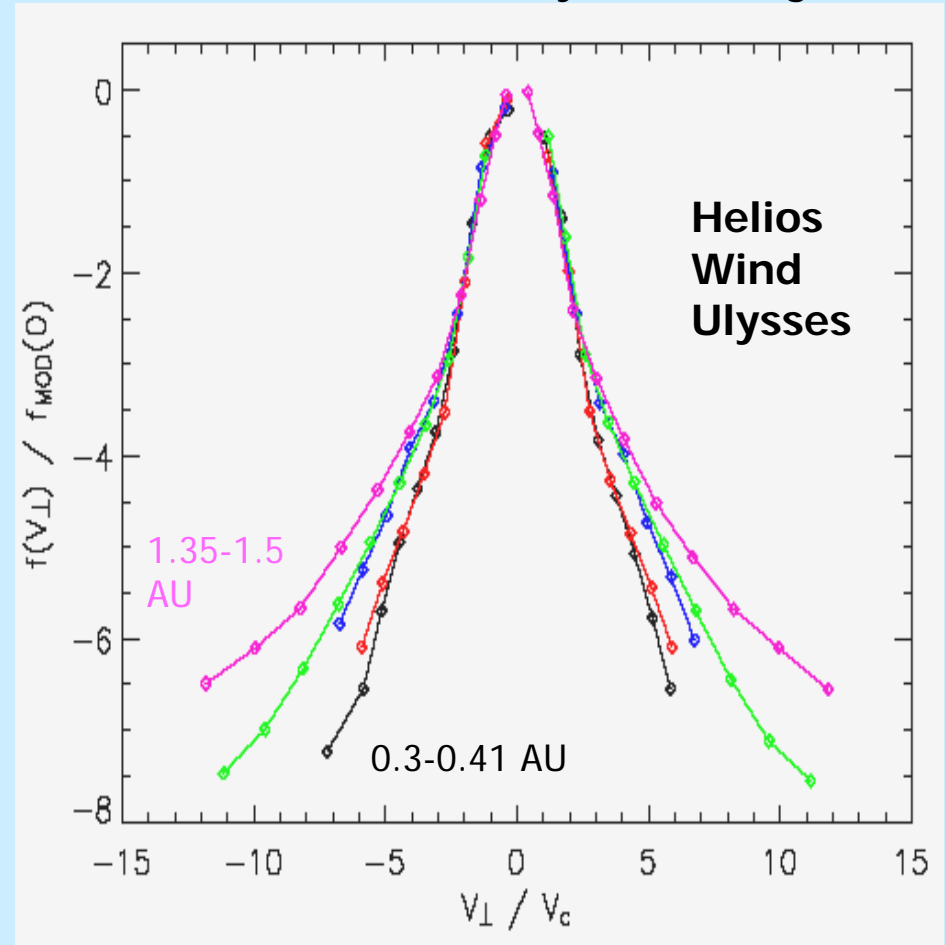
Core (96%), halo (4%) electrons, and „strahl“

Solar wind electrons: Core-halo evolution

Halo is relatively increasing while strahl is diminishing.



Normalized core remains constant while halo is relatively increasing.



Fluid description

Moments of the Vlasov/Boltzmann equation:

Density: $n_j = \int d^3v f_j(\mathbf{x}, \mathbf{v}, t)$

Flow velocity: $\mathbf{U}_j = 1/n_j \int d^3v f_j(\mathbf{x}, \mathbf{v}, t) \mathbf{v}$

Thermal speed: $v_j^2 = 1/(3n_j) \int d^3v f_j(\mathbf{x}, \mathbf{v}, t) (\mathbf{v} - \mathbf{U}_j)^2$

Temperature: $T_j = m_j v_j^2 / k_B$

Heat flux: $Q_j = 1/2 m_j \int d^3v f_j(\mathbf{x}, \mathbf{v}, t) (\mathbf{v} - \mathbf{U}_j) (\mathbf{v} - \mathbf{U}_j)^2$

Electron heat conduction

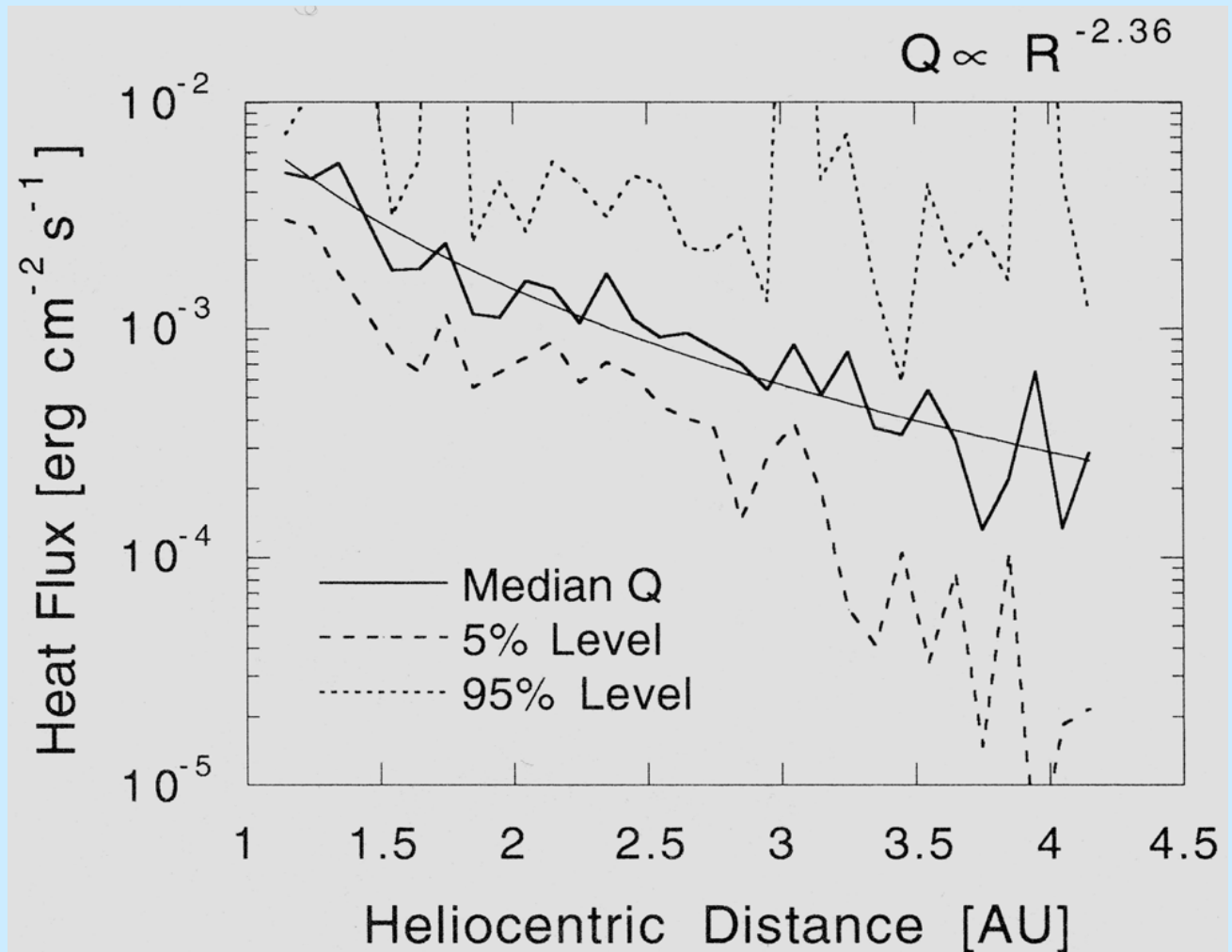
Heat carried
by halo
electrons!

$$T_H = 7 T_C$$

Interplanetary
potential:

$$\Phi = 50-100 \text{ eV}$$

$$\underline{E} = -1/n_e \underline{\nabla} p_e$$



$$\underline{Q}_e \neq -\kappa \underline{\nabla} T_e$$

Theoretical description

Boltzmann-Vlasov kinetic equations for protons, alpha-particles (4%), minor ions and electrons

Distribution functions

Kinetic equations

- + Coulomb collisions (Landau)
- + Wave-particle interactions
- + Micro-instabilities (Quasilinear)
- + Boundary conditions

→ **Particle velocity distributions and field power spectra**

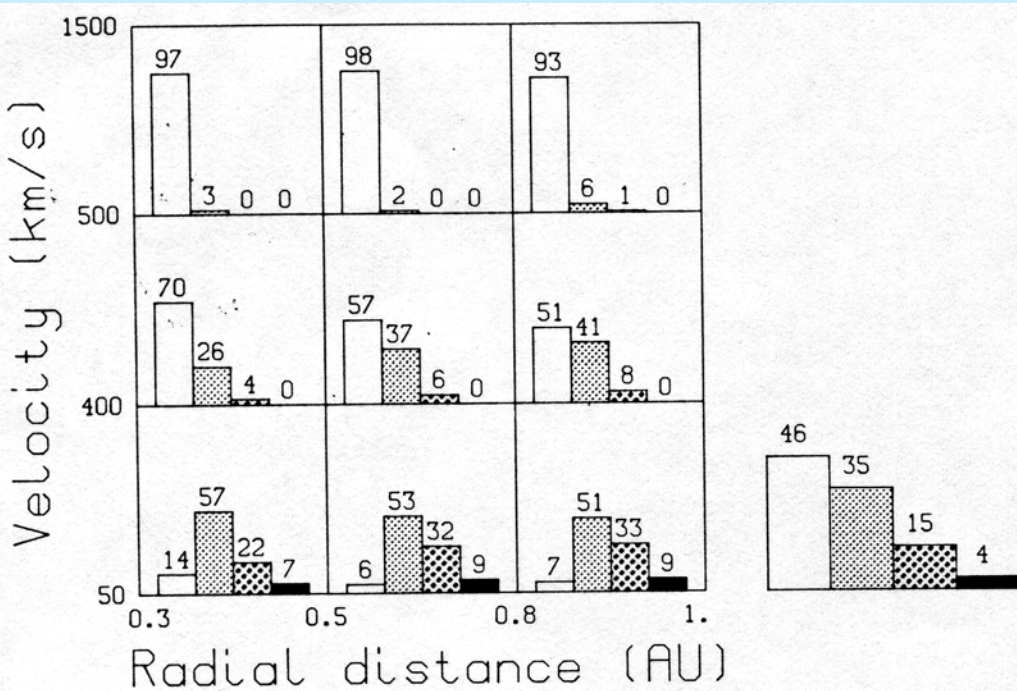
Moments

Multi-Fluid (MHD) equations

- + Collision terms
- + Wave (bulk) forces
- + Energy addition
- + Boundary conditions

→ **Single/multi fluid parameters**

Proton Coulomb collision statistics



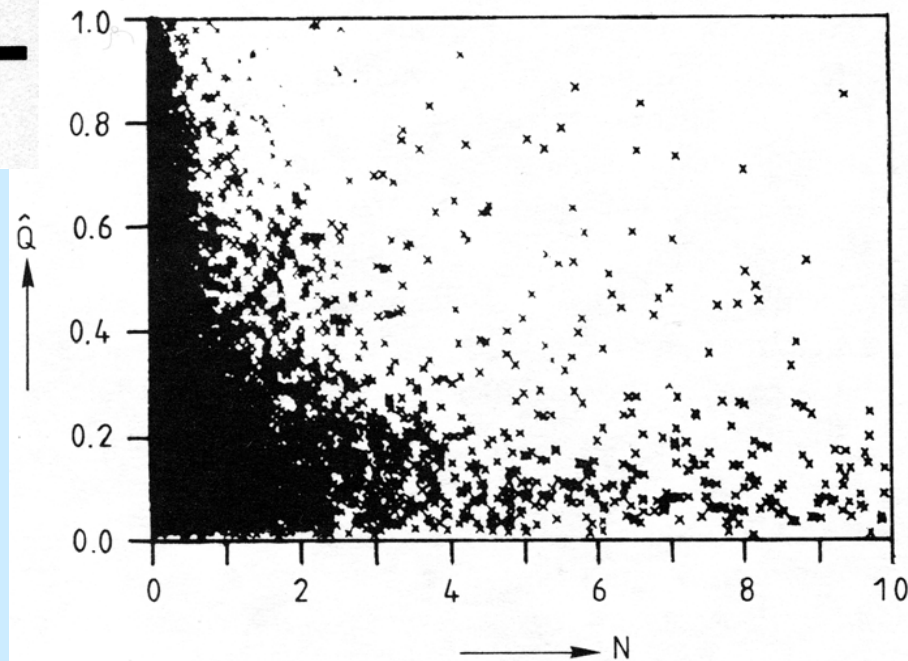
- Fast protons are collisionless!
- Slow protons show collision effects!

Proton heat flux regulation

N : <0.1 ; $[0.1, 1]$; $[1, 5]$; >5

$$N = \tau_{\text{exp}} v_c \sim n_p V^{-1} T_p^{-3/2}$$

Livi et al., JGR, **91**, 8045, 1986



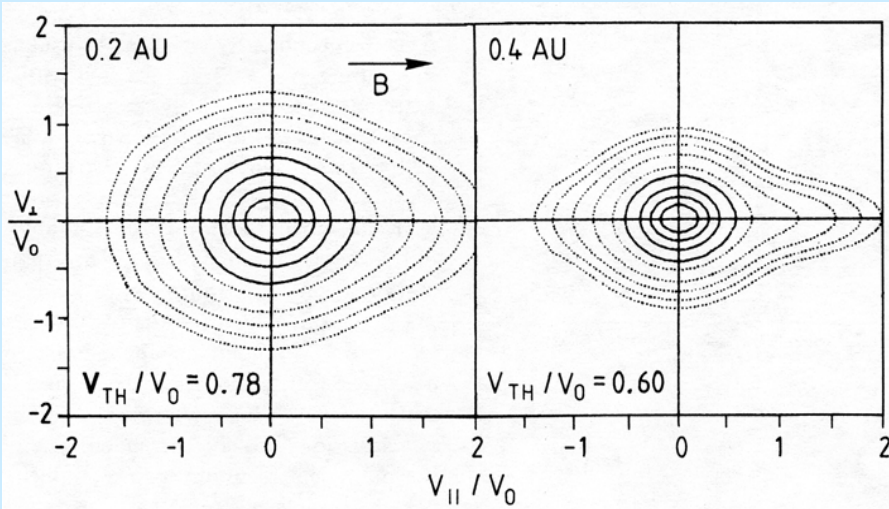
Coulomb collisions and quasilinear wave-particle interactions

$$\frac{\delta f}{\delta t} = \frac{\partial}{\partial \mathbf{v}} \cdot \left(-\mathbf{A} + \mathcal{D} \cdot \frac{\partial}{\partial \mathbf{v}} \right) f$$

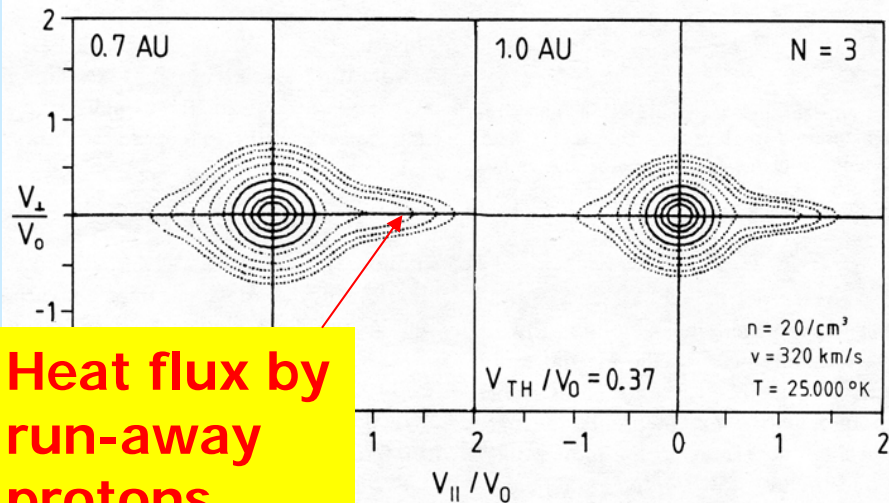
Coulomb collisions and/or **wave-particle interactions** are represented by a second-order differential operator, including the acceleration vector $\mathbf{A}(\mathbf{v})$ and diffusion tensor $\mathcal{D}(\mathbf{v})$:

Parameter	Chromo-sphere	Corona (1R _S)	Solar wind (1AU)
n_e (cm ⁻³)	10 ¹⁰	10 ⁷	10
T_e (K)	6-10 10 ³	1-2 10 ⁶	10 ⁵
λ_e (km)	10	1000	10 ⁷

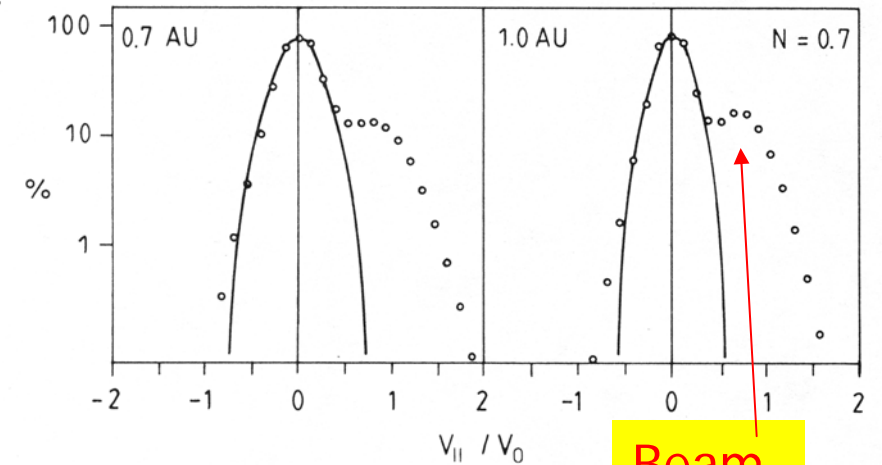
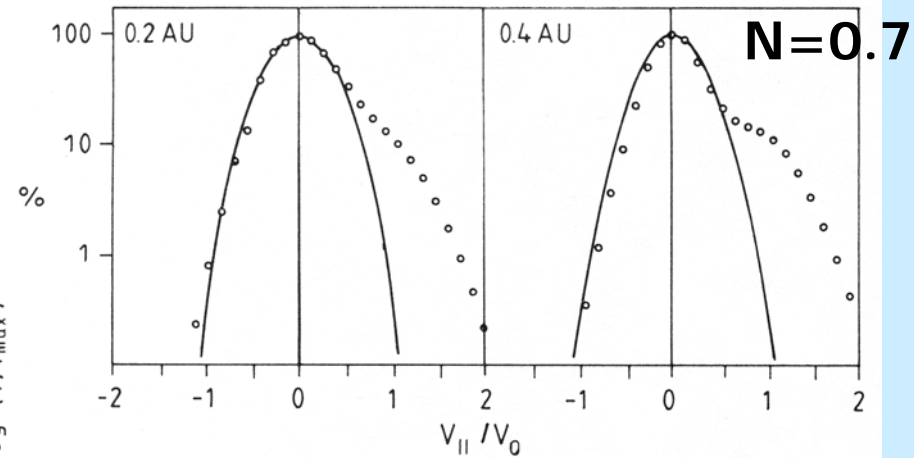
Coulomb collisions in slow wind



N=3



Heat flux by run-away protons



Beam

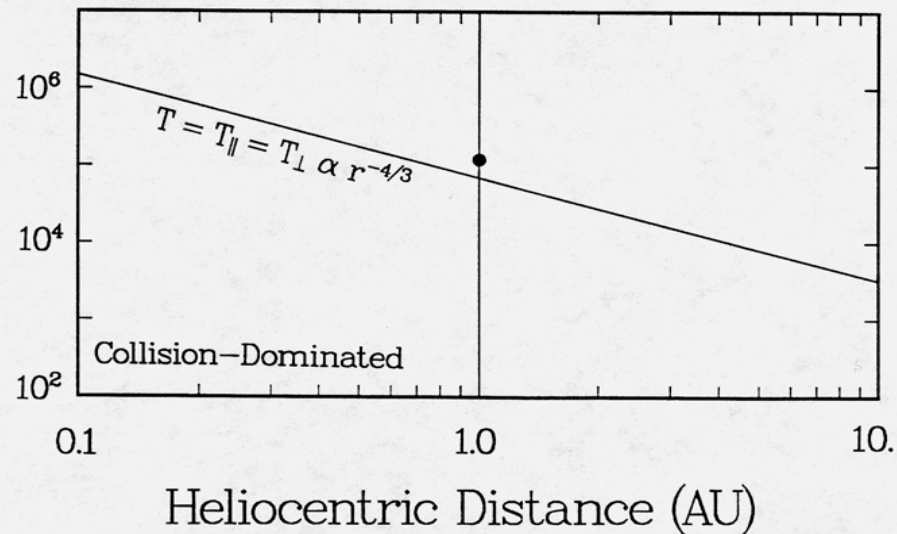
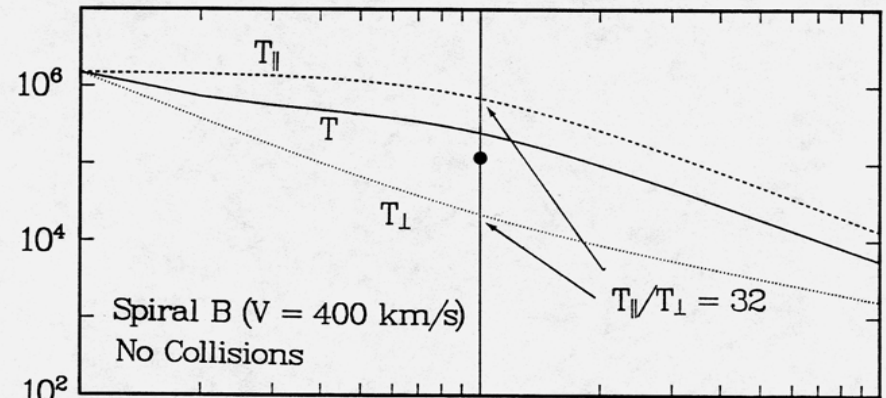
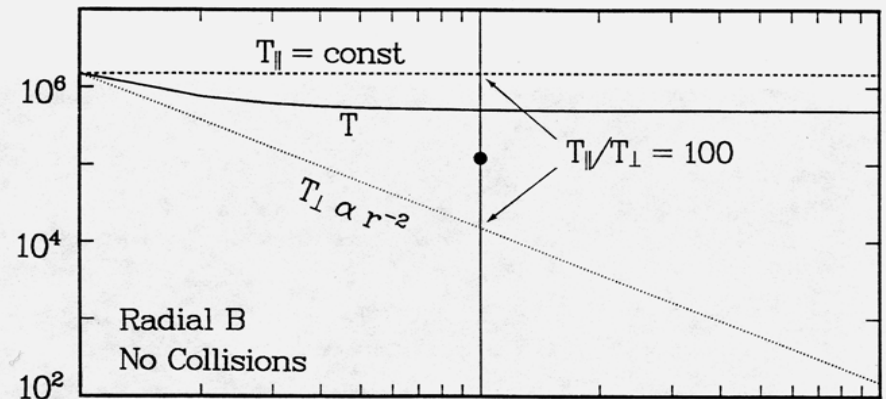
Collisions and geometry

Double adiabatic invariance, \rightarrow extreme anisotropy not observed!

Spiral reduces anisotropy!

Adiabatic collision-dominated \rightarrow isotropy, is not observed!

Core Electron Temperature (K)

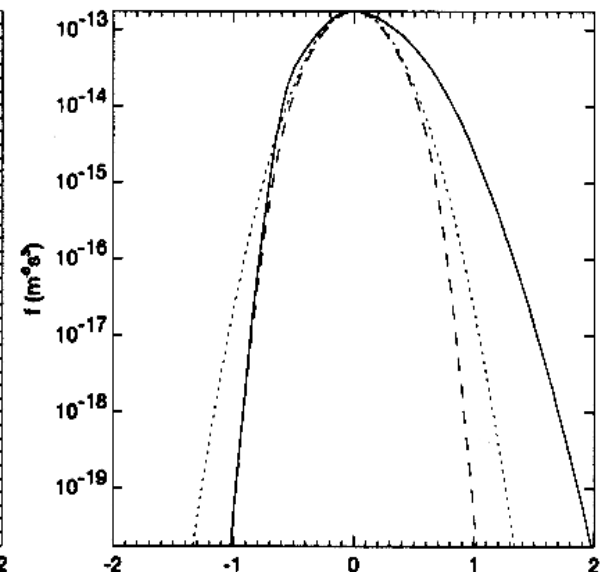
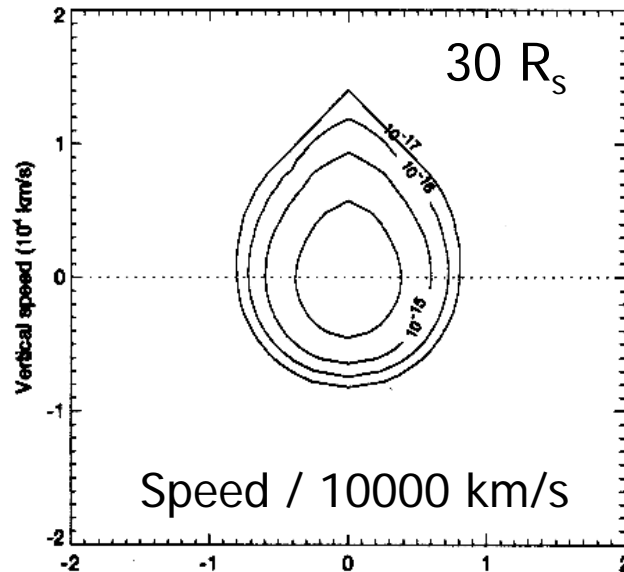
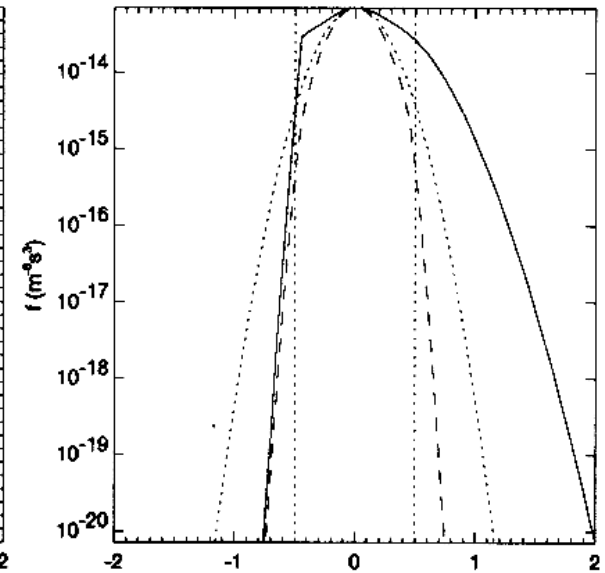
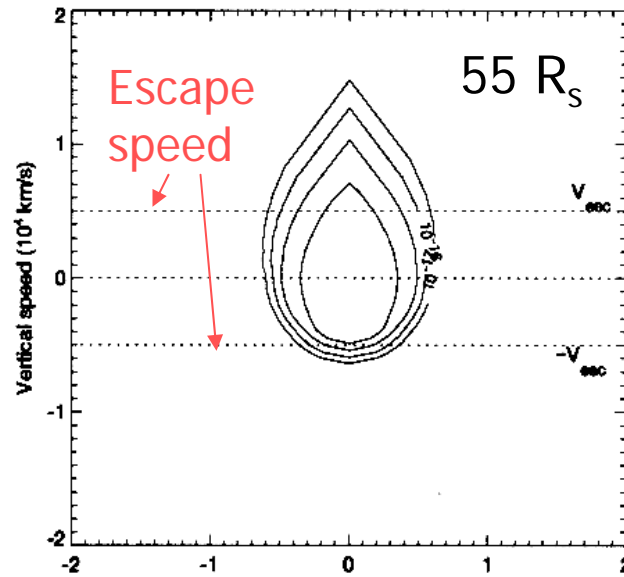


Coulomb collisions and electrons

Integration of Fokker-Planck equation

- Velocity filtration is weak!
- Strahl formation by escape electrons
- Core bound by electric field

Lie-Svendson et al.,
JGR, **102**, 4701, 1997



Wave-particle interactions

Dispersion relation using measured or model distribution functions $f(\underline{v})$, e.g. for electrostatic waves:

$$\varepsilon_L(\underline{k}, \omega) = 0 \rightarrow \omega(\underline{k}) = \omega_r(\underline{k}) + i\gamma(\underline{k})$$

Dielectric constant is functional of $f(\underline{v})$, which may when being non-Maxwellian contain free energy for wave excitation.

$\gamma(\underline{k}) > 0 \rightarrow$ micro-instability.....

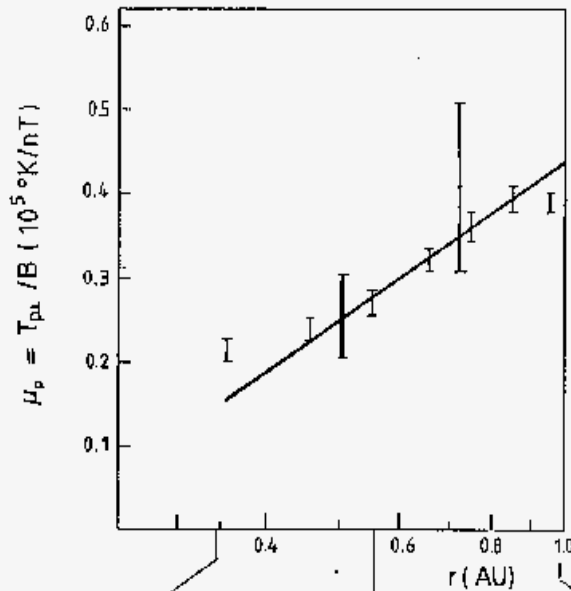
Resonant particles:

$$\omega(\underline{k}) - \underline{k} \cdot \underline{v} = 0 \quad (\text{Landau resonance})$$

$$\omega(\underline{k}) - \underline{k} \cdot \underline{v} = \pm \Omega_j \quad (\text{cyclotron resonance})$$

\rightarrow Energy and momentum exchange between waves and particles. Quasi-linear or non-linear relaxation.....

Heating of protons by cyclotron and Landau resonance

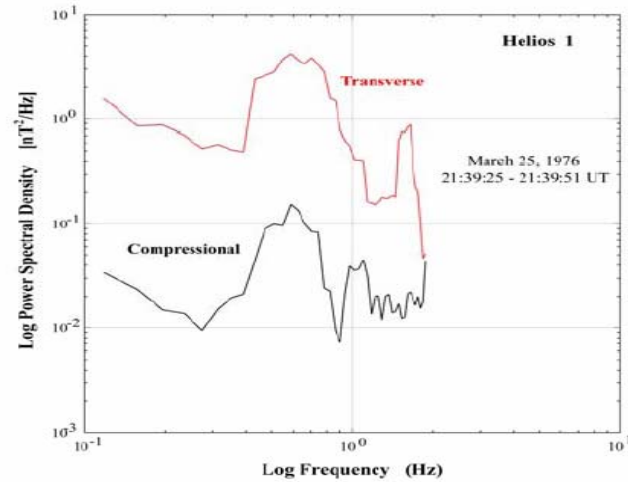
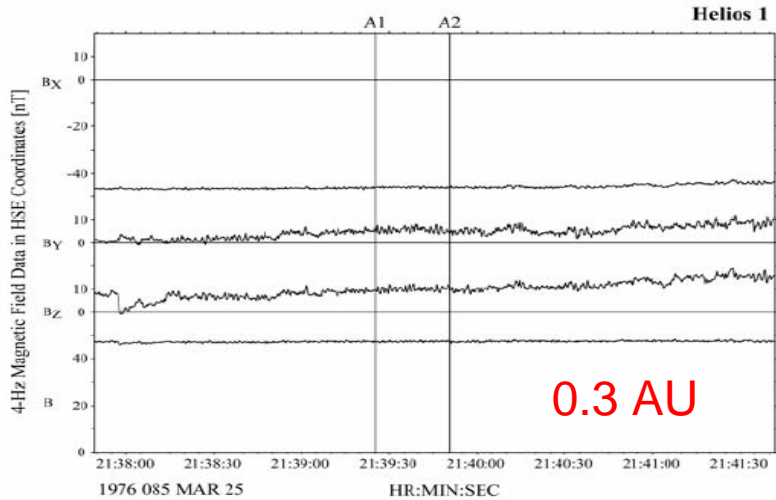


- Increasing magnetic moment
- Decelerating proton/ion beams
- Evolving temperature anisotropy



Velocity
distribution
functions

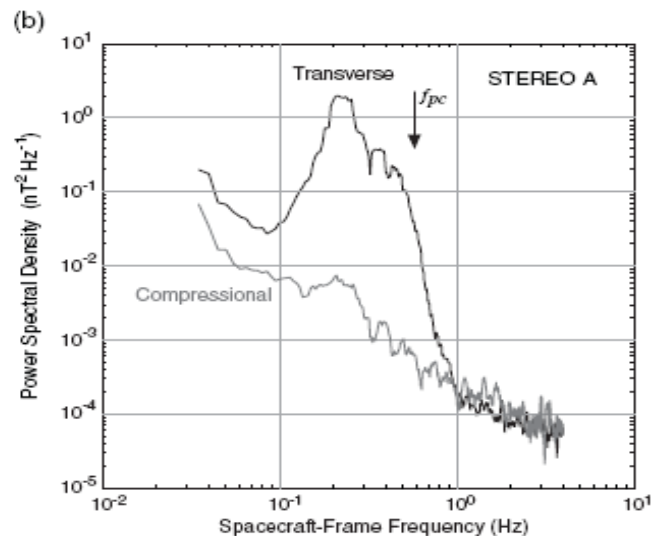
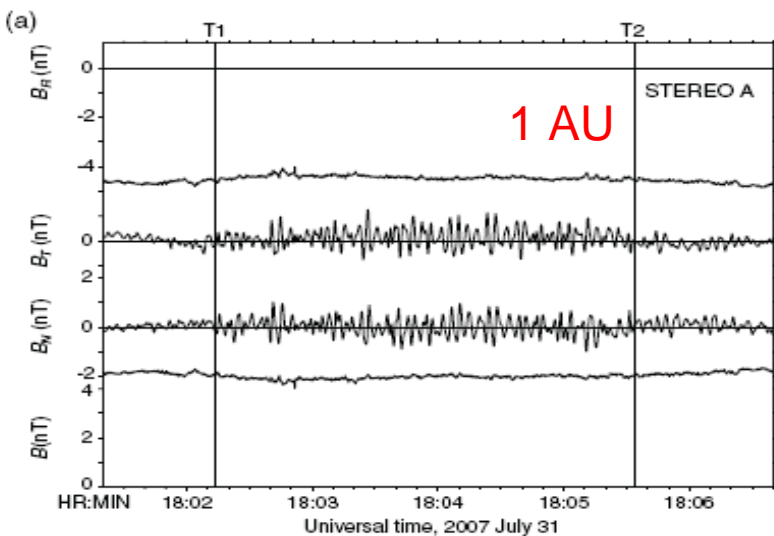
Ion cyclotron waves



Helios

Jian and Russell,
The Astronomy
and Astrophysics
Decadal Survey,
Science White
Papers, no. 254,
2009

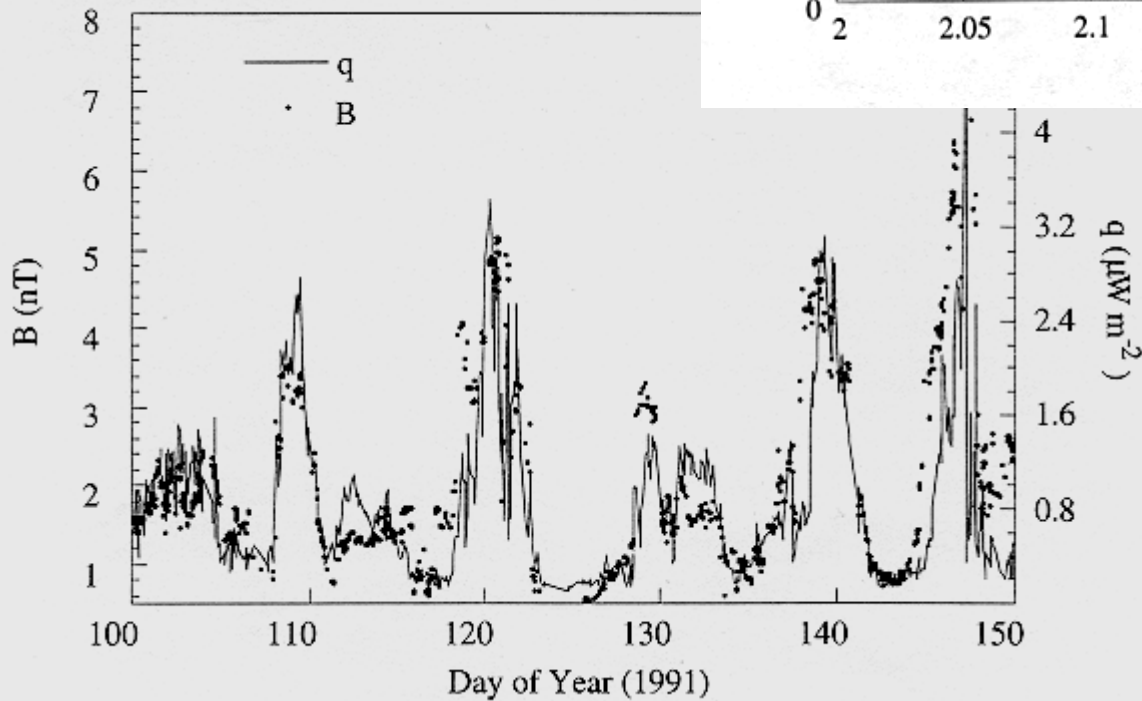
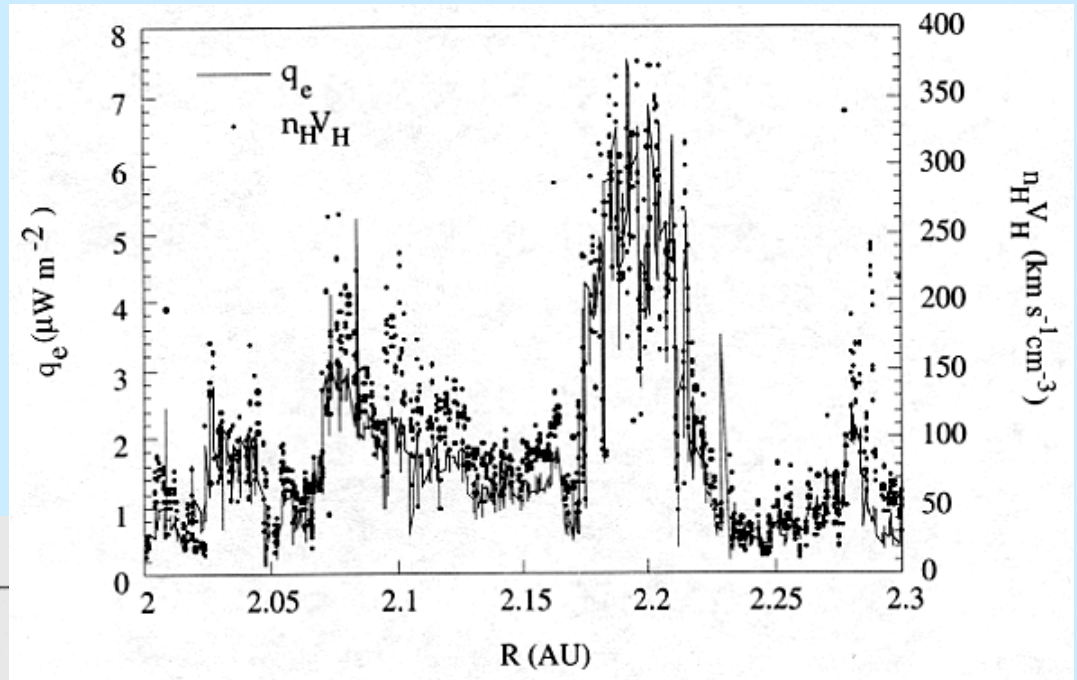
Parallel in- and outward propagation



STEREO

Jian et al.
ApJ, 2009

Whistler regulation of electron heat flux



- Halo electrons carry heat flux
- Heat flux varies with B or V_A
- Whistler instability regulates drift

Sime et al., JGR, 1994

Kinetic processes in the solar corona and solar wind II

- Plasma is multi-component and nonuniform
 - **multi-fluid or kinetic physics is required**
- Plasma is tenuous and turbulent
 - **free energy for micro-instabilities**
 - **resonant wave-particle interactions**
 - **collisions described by Fokker-Planck operator**

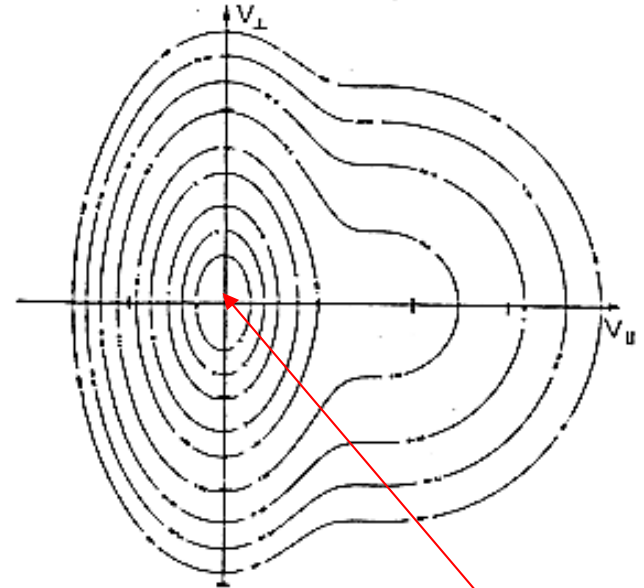
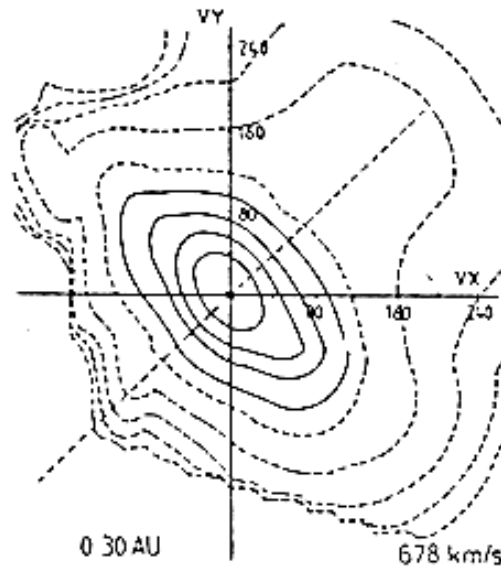
Problem: Transport properties of the plasma, which involves multiple scales.....

Proton temperature anisotropy

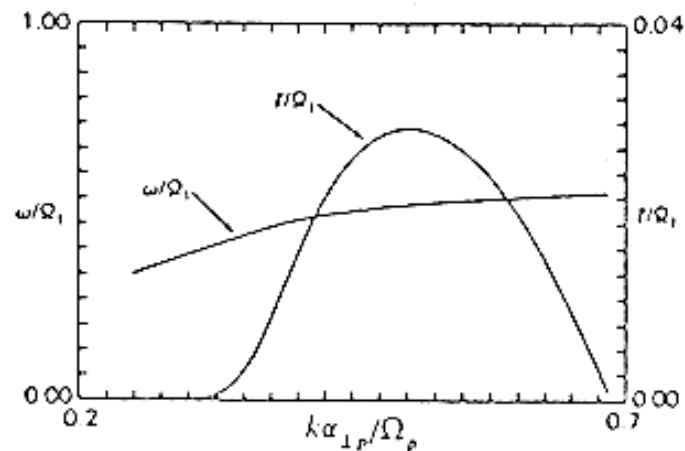
- Measured and modelled proton velocity distribution
- Growth of ion-cyclotron waves!
- Anisotropy-driven instability by large perpendicular T_{\perp}

$$\omega \approx 0.5\Omega_p$$

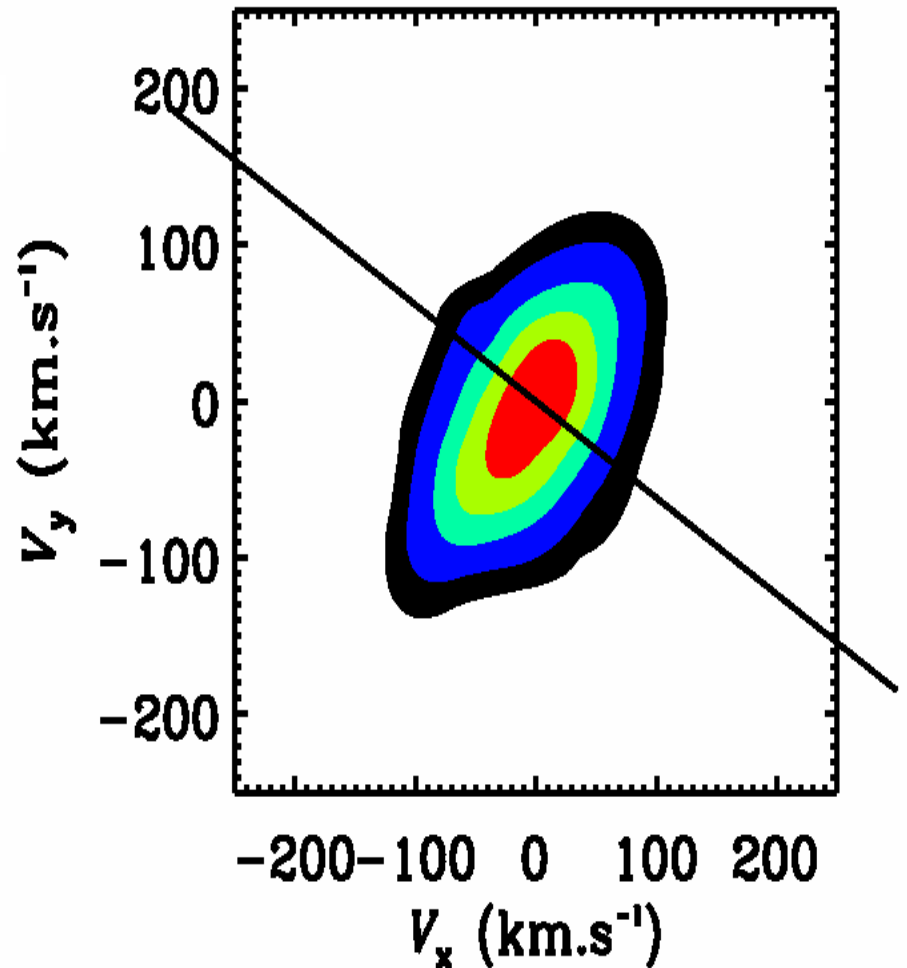
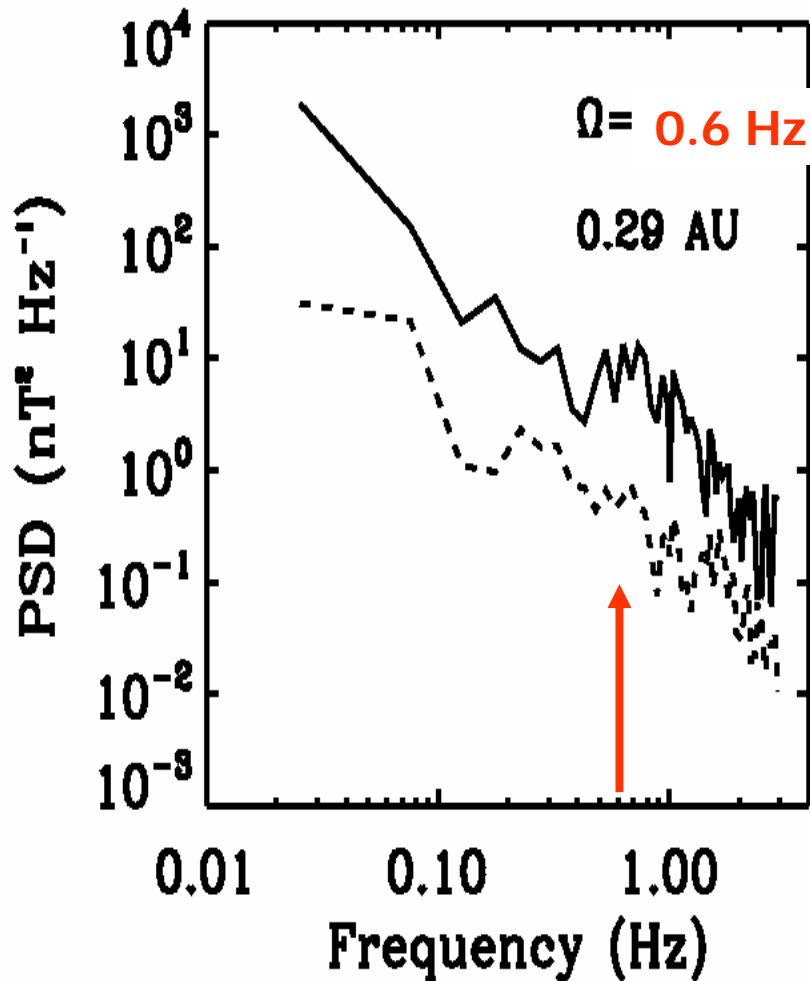
$$\gamma \approx 0.05\Omega_p$$



anisotropy



Correlation of anisotropy with Alfvén-cyclotron wave power

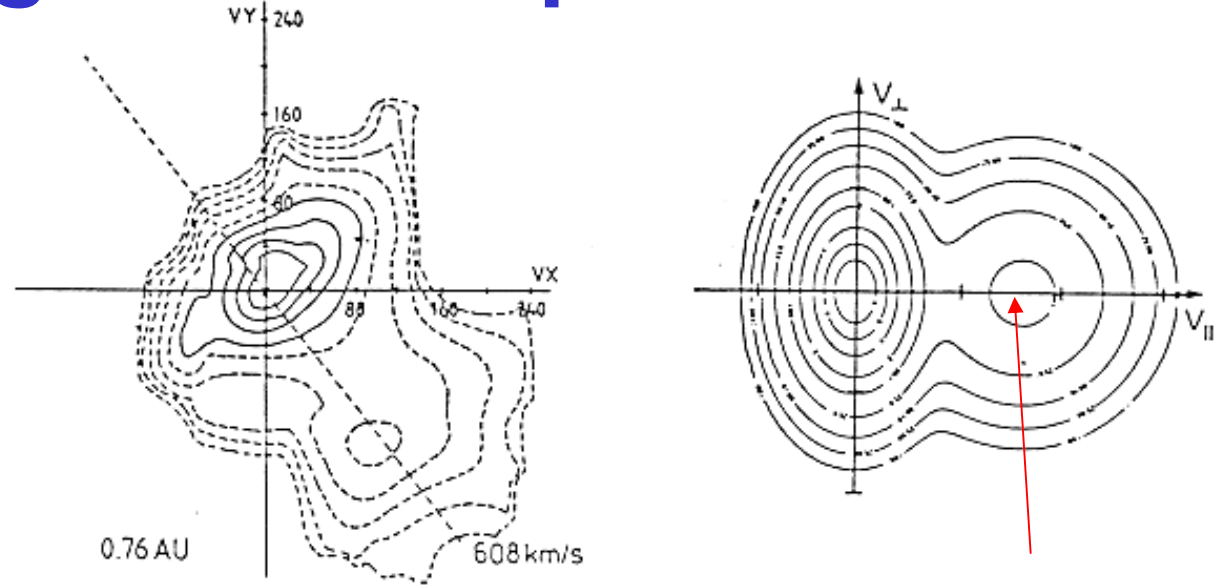


Wave regulation of proton beam

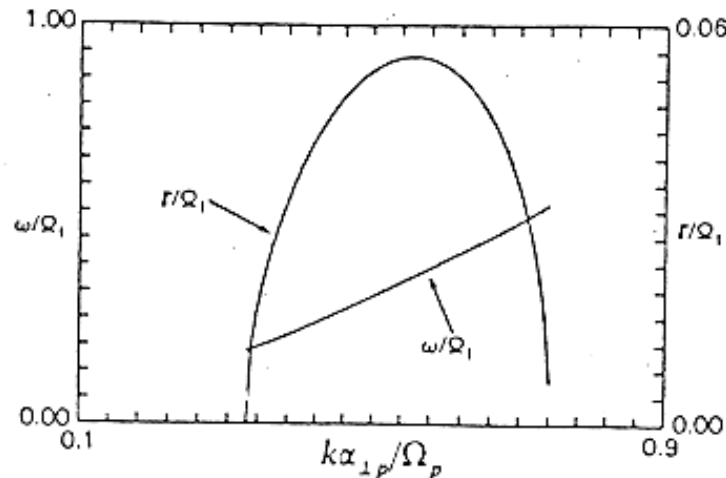
- Measured and modelled velocity distribution
- Growth of fast mode waves!
- Beam-driven instability, large drift speed

$$\omega \approx 0.4\Omega_p$$

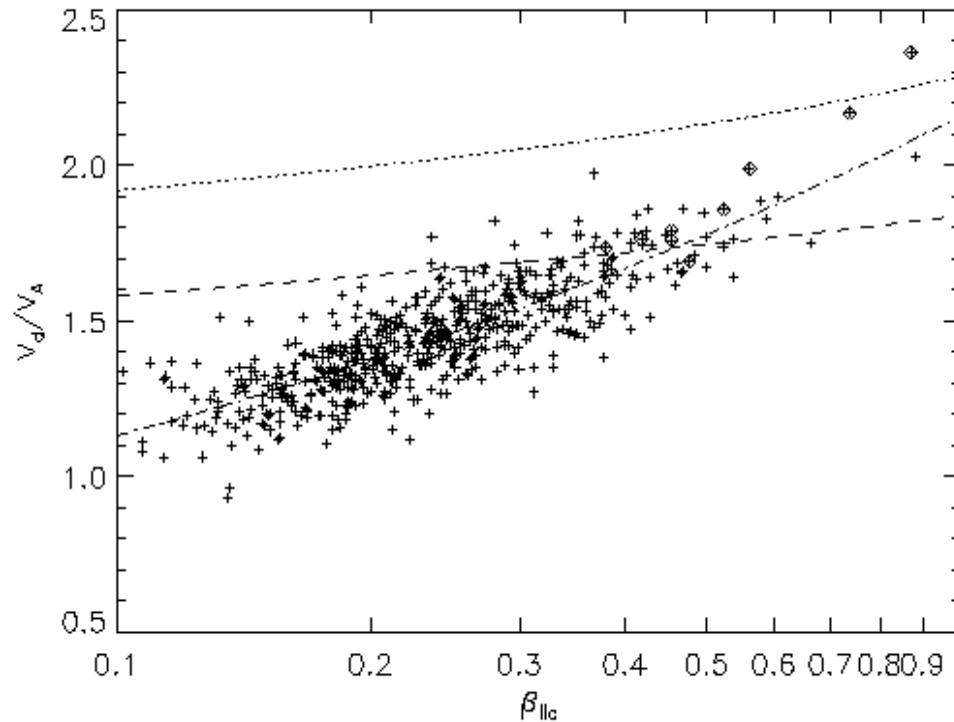
$$\gamma \approx 0.06\Omega_p$$



beam



Beam drift versus core plasma beta



- Origin of the beam not yet explained (mirror force in corona, Coulomb collisions)
- Quasilinear resonant pitch-angle diffusion determines the kinetics (thermodynamics) of solar wind protons.
- Proton beam drift speed is regulated by wave-particle interactions and depends on the plasma beta.

• The dotted (dashed) line corresponds to a relative beam density of 0.05 (0.2).

$$V_d/V_A = (2.16 \pm 0.03) \beta_{\parallel c}^{(0.28 \pm 0.01)}$$

(empirical fit to data given by the dot-dash line)

Kinetic plasma instabilities

- Observed velocity distributions at margin of stability
- Selfconsistent quasi- or non-linear effects not well understood
- Wave-particle interactions are the key to understand ion kinetics in corona and solar wind!

Marsch, 1991; Gary, Space Science Rev., **56**, 373, 1991

Wave mode	Free energy source
Ion acoustic	Ion beams, electron heat flux
Ion cyclotron	Temperature anisotropy
Whistler (Lower Hybrid)	Electron heat flux
Magnetosonic	Ion beams, differential streaming

Quasi-linear (pitch-angle) diffusion

Diffusion
equation

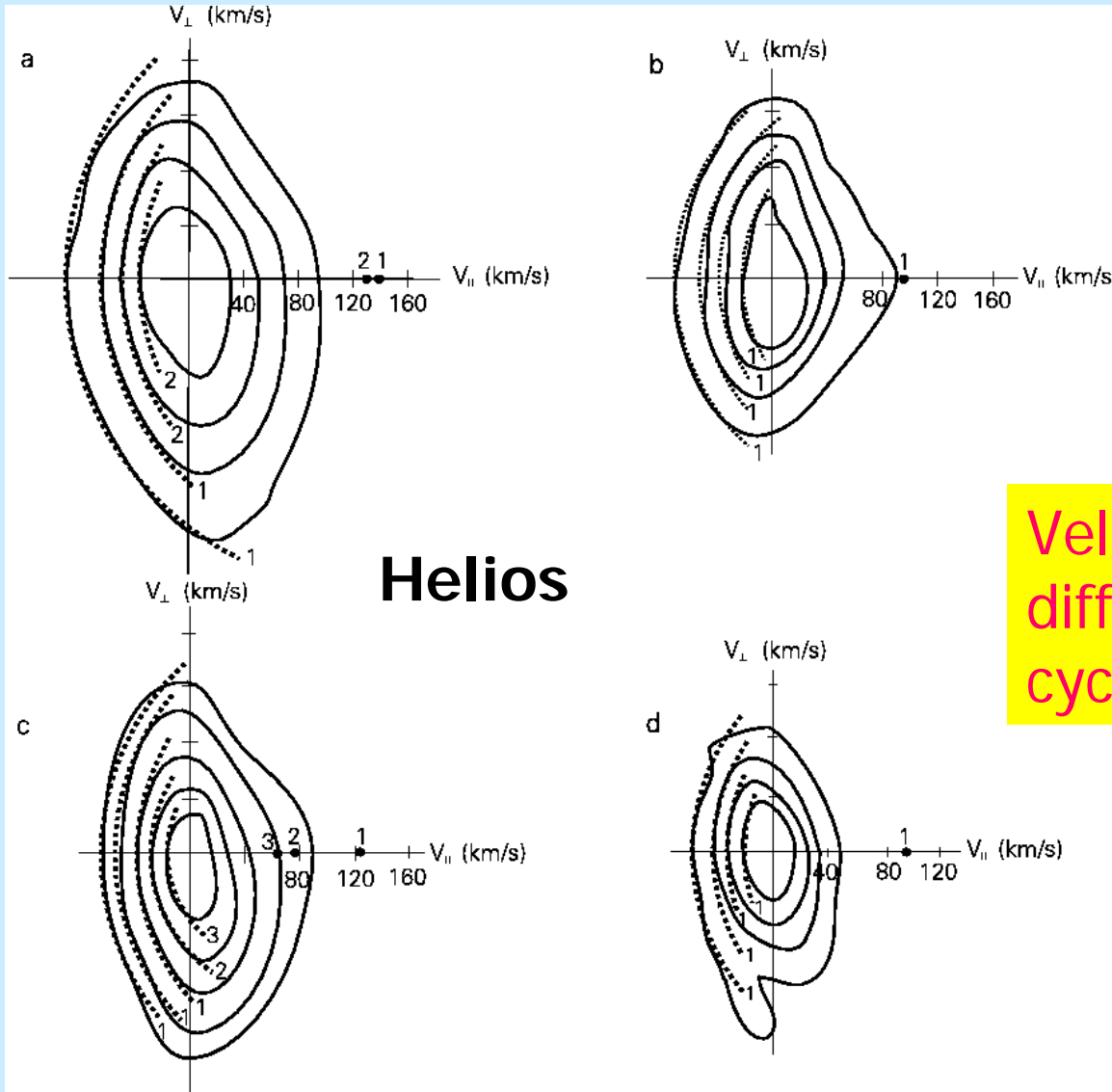
$$\frac{\partial}{\partial t} f_j(V_{\perp}, V_{\parallel}, t) = \sum_M \sum_{s=-\infty}^{+\infty} \frac{1}{(2\pi)^3} \int_{-\infty}^{+\infty} d^3k \hat{\mathcal{B}}_M(\mathbf{k}) \\ \times \frac{1}{V_{\perp}} \frac{\partial}{\partial \alpha} \left(V_{\perp} \nu_j(\mathbf{k}, s; V_{\parallel}, V_{\perp}) \frac{\partial}{\partial \alpha} f_j(V_{\perp}, V_{\parallel}, t) \right)$$

Pitch-angle
gradient in
wave frame

$$\frac{\partial}{\partial \alpha} = V_{\perp} \frac{\partial}{\partial V_{\parallel}} - \left(V_{\parallel} - \frac{\omega_M(\mathbf{k})}{k_{\parallel}} \right) \frac{\partial}{\partial V_{\perp}}$$

Kennel and Engelmann,
Phys. Fluids, **9**, 2377, 1966

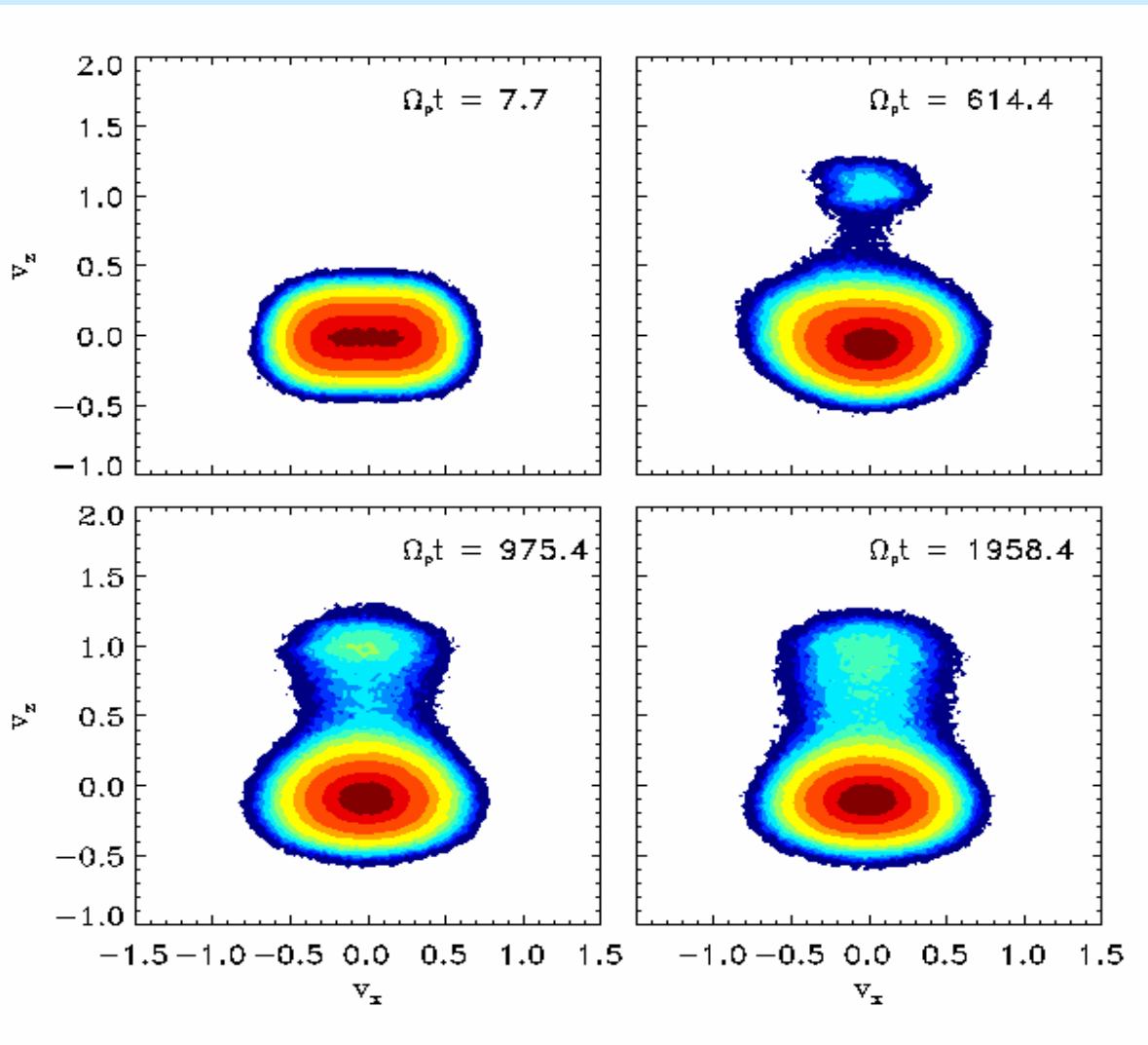
Pitch-angle diffusion of protons



VDF contours are segments of circles centered in the wave frame ($< V_A$)

Velocity-space resonant diffusion caused by the cyclotron-wave field!

Proton core heating and beam formation



VDFs as obtained by numerical simulation of the decay of Alfvén-cyclotron waves and the related ion kinetics

Contour plots of the proton VDF in the v_x - v_z -plane for the dispersive-wave case at four instants of time. The color coding of the contours corresponds, respectively, to 75 (dark red), 50 (red), 10 (yellow) percent of the maximum.

Summary

- **In-situ ion and electron measurements indicate strong deviations from local (collisional) thermal equilibrium**
- **Wave-particle interactions and micro-instabilities regulate the kinetic features of particle velocity distributions**
- **Kinetic models are required to describe the essential features of the plasma in the solar wind**
- **The non-equilibrium thermodynamics in the tenuous solar wind involve particle interaction with micro-turbulent fields**
- **Wave energy transport as well as cascading and dissipation in the kinetic domain are still not well understood**