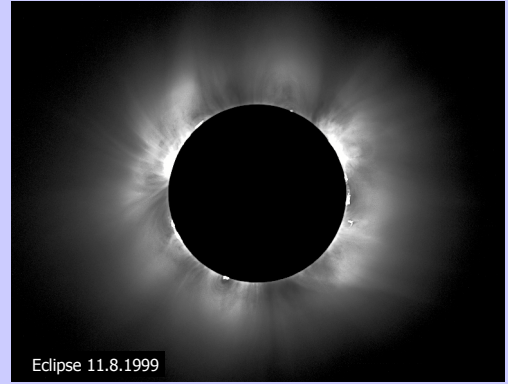


## Coronal expansion and solar wind

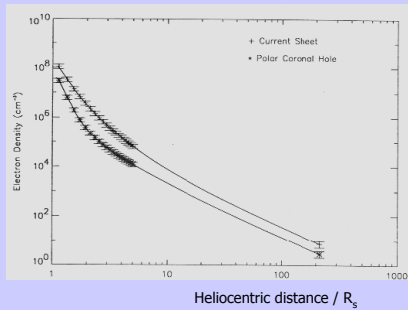
- The large solar corona
- Coronal and interplanetary temperatures
- Coronal expansion and solar wind
- The heliosphere
- Origin of solar wind in magnetic network
- Multi-fluid models of the solar wind

## The visible solar corona



Eclipse 11.8.1999

## Electron density in the corona



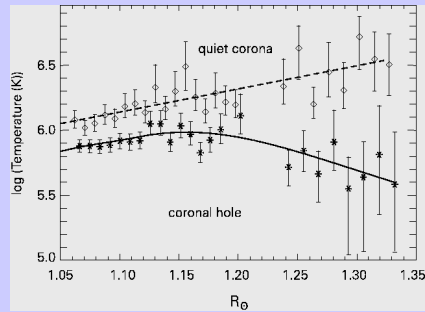
+ Current sheet and streamer belt, closed

• Polar coronal hole, open magnetically

Guthathakurta and Sittler, 1999, Ap.J., 523, 812

Skylab coronagraph/Ulysses in-situ

## Electron temperature in the corona



Streamer belt, closed

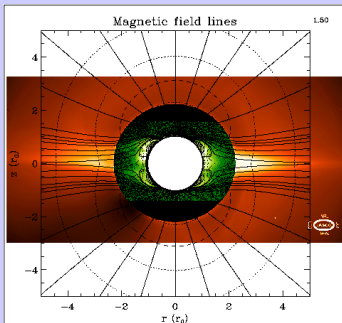
Coronal hole, open magnetically

David et al., A&A, 336, L90, 1998

Heliocentric distance SUMER/CDS SOHO

## Coronal magnetic field and density

Dipolar, quadrupolar, current sheet contributions



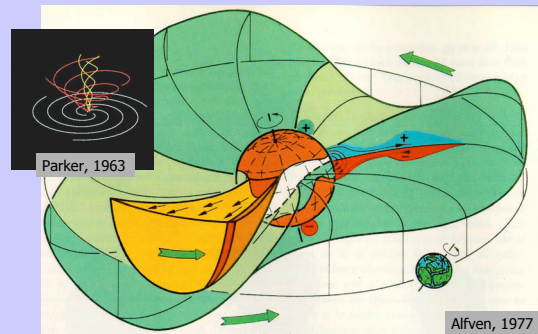
Polar field: B = 12 G

Current sheet is a symmetric disc anchored at high latitudes!

Banaszkiewicz et al., 1998; Schwenn et al., 1997

LASCO C1/C2 images (SOHO)

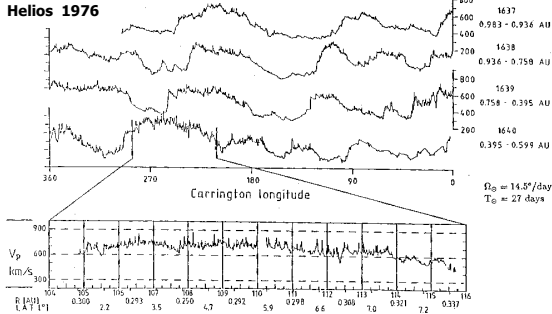
## Solar wind stream structure and heliospheric current sheet



Parker, 1963

Alfvén, 1977

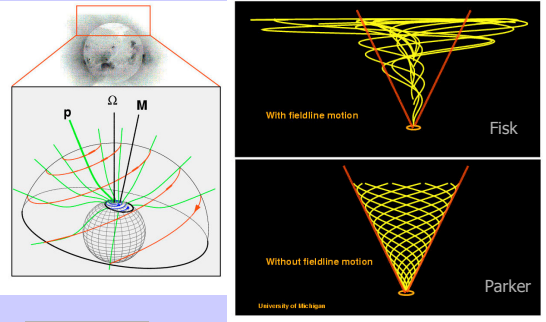
## Solar wind fast and slow streams



Alfvén waves and small-scale structures

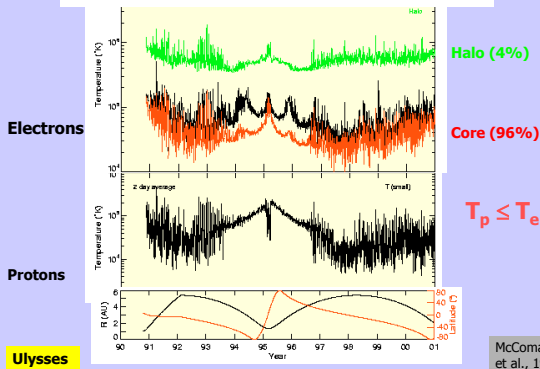
Marsch, 1991

## Model of coronal-heliospheric field



Fisk, JGR, 1996

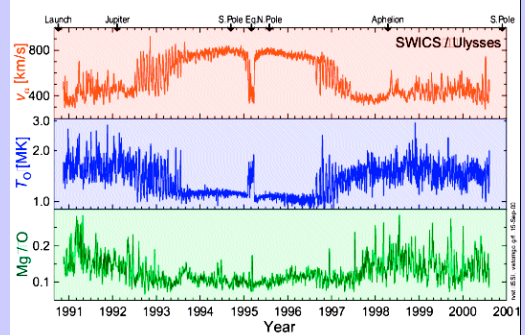
## Heliospheric temperatures



Ulysses

McComas et al., 1998

## Correlations between wind speed and corona temperature



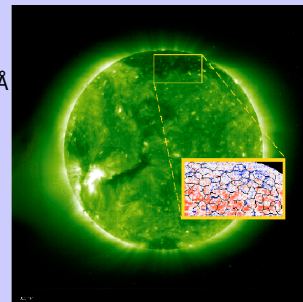
## Fast solar wind parameters

- Energy flux at 1  $R_S$ :  $F_E = 5 \cdot 10^5 \text{ erg cm}^{-2} \text{ s}^{-1}$
- Speed beyond 10  $R_S$ :  $V_p = (700 - 800) \text{ km s}^{-1}$
- Proton flux at 1 AU:  $n_p V_p = 2 \cdot 10^8 \text{ cm}^{-2} \text{ s}^{-1}$
- Density at 1 AU:  $n_p = 3 \text{ cm}^{-3}$ ;  $n_\alpha/n_p = 0.04$
- Temperatures at 1 AU:  
 $T_p = 3 \cdot 10^5 \text{ K}$ ;  $T_\alpha = 10^6 \text{ K}$ ;  $T_e = 1.5 \cdot 10^5 \text{ K}$
- Heavy ions:  $T_i \approx m_i / m_p T_p$ ;  $V_i - V_p = V_A$

Schwenn and Marsch, 1990, 1991

## On the source regions of the fast solar wind in coronal holes

Image: EIT  
 Corona in  
 Fe XII 195 Å  
 at 1.5 M K



Insert: SUMER  
 Ne VIII 770 Å  
 at 630 000 K

Chromospheric  
 network  
 Doppler shifts  
 Red: down  
 Blue: up

Outflow at  
 lanes and  
 junctions

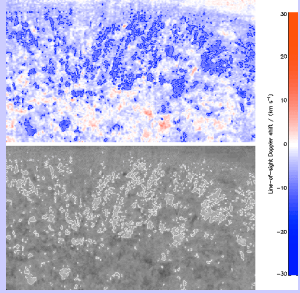
Hassler et al.,  
 Science 283,  
 811-813, 1999

# The source regions of the fast solar wind in polar coronal holes

Doppler-shift

map of solar polar cap: 520" x 300" gliding step size: 3"

Wilhelm et al., A&A, 353, 749, 2000



SUMER Ne VIII 770 Å at 630 000 K

Encircled contours: Doppler shift > 5 km/s

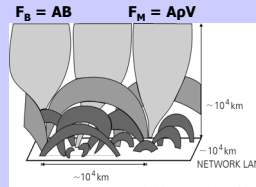
Radiance

September 21, 1996

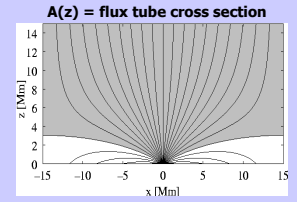
# Magnetic network loops and funnels

Structure of transition region

Magnetic field of coronal funnel

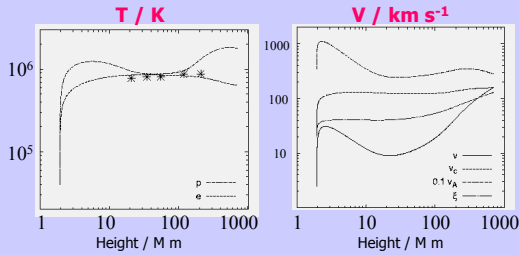


Dowdy et al., Solar Phys., 105, 35, 1986



Hackenberg et al., Space Sci. Rev., 87, 207, 1999

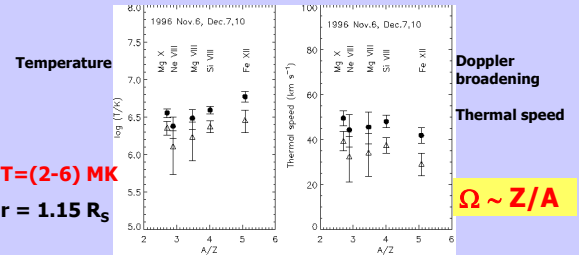
# Height profiles in funnel flows



- Heating by wave sweeping
- Steep temperature gradients

• Critical point at 1 R<sub>S</sub>  
Hackenberg, Marsch, Mann, A&A, 360, 1139, 2000

# Heating and acceleration of ions by cyclotron and Landau resonance



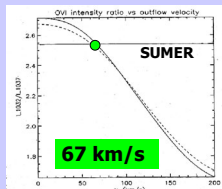
T = (2-6) MK  
r = 1.15 R<sub>S</sub>

Tu et al., Space Sci. Rev., 87, 331, 1999

Ion heating ~ mass/charge

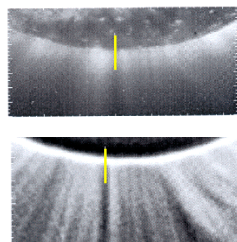
$$\Omega \sim Z/A$$

# Outflow speed in interplume region at the coronal base



O VI 1031.9 Å / 1037.2 Å line ratio; Doppler dimming

T<sub>e</sub> = T<sub>i</sub> = 0.9 MK, n<sub>e</sub> = 1.8 10<sup>7</sup> cm<sup>-3</sup>

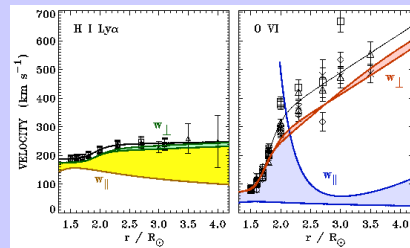


1.05 R<sub>S</sub>  
EIT  
FeIX/X

Eclipse  
26/02  
1998  
18:33 UT

Patsourakos and Vial, A&A, 359, L1, 2000

# Oxygen and hydrogen thermal speeds in coronal holes

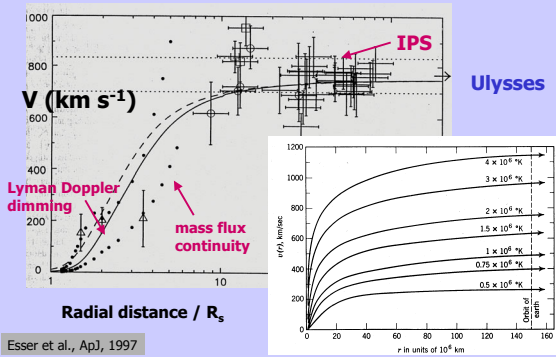


Very Strong perpendicular heating of Oxygen !

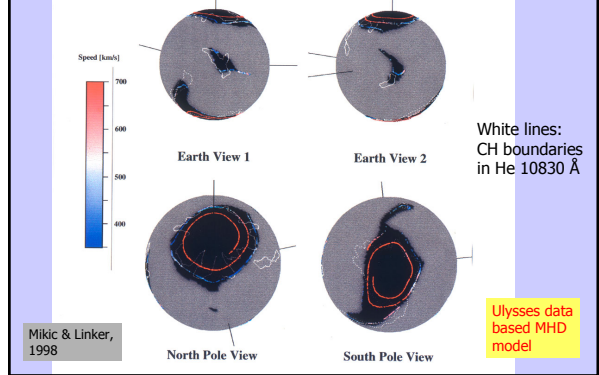
Cranmer et al., Ap. J., 511, 481, 1998

Large anisotropy: T<sub>⊥</sub>/T<sub>∥</sub> ≥ 10

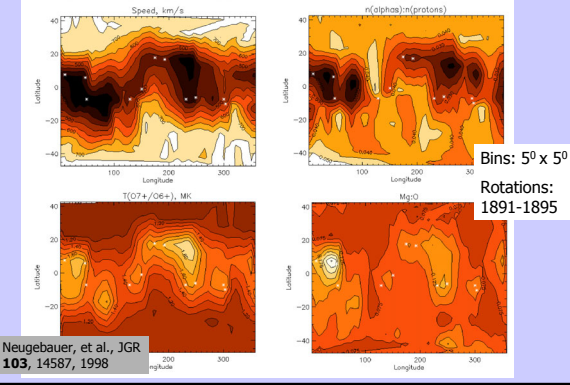
## Fast solar wind speed profile



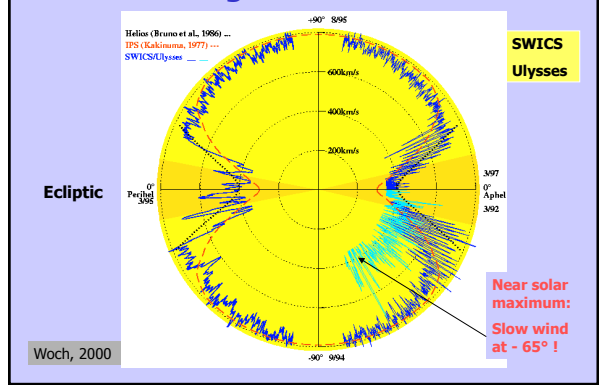
## Boundaries of coronal holes



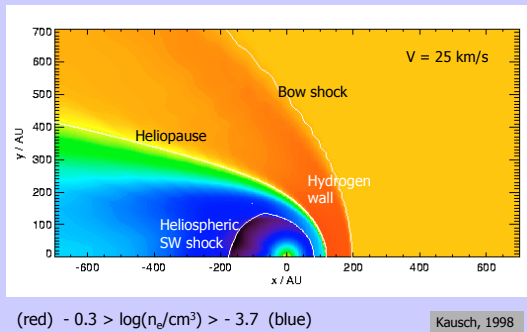
## Solar wind in Carrington longitude



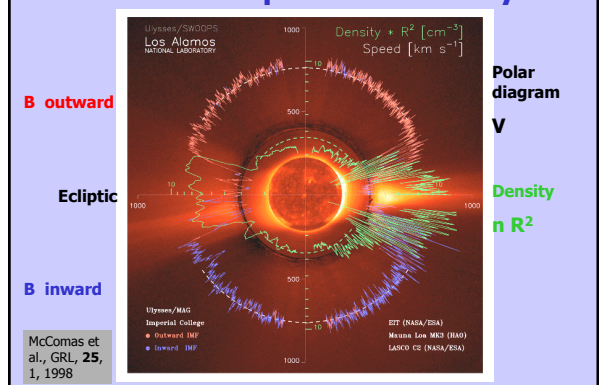
## Polar diagram of solar wind



## Heliosphere and local interstellar medium



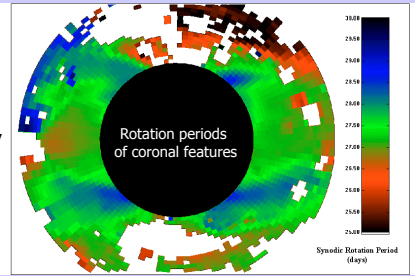
## Solar wind speed and density



## Rotation of solar corona

Fe XIV  
5303 Å

Time series:  
1 image/day  
(24-hour  
averages)



LASCO  
/SOHO

Stenborg et al., 1999

Long-lived coronal patterns exhibit uniform rotation at the equatorial rotation period!

## Sun's loss of angular momentum carried by the solar wind I

Induction equation:

$$\nabla \times (\mathbf{V} \times \mathbf{B}) = 0 \quad \rightarrow \quad r (V_r B_\phi - B_r V_\phi) = -r_0 B_0 \Omega_0 r_0$$

Momentum equation:

$$\rho \mathbf{V} \cdot \nabla V_\phi = 1/4\pi \mathbf{B} \cdot \nabla B_\phi \quad \rightarrow \quad r (\rho V_r V_\phi - B_r B_\phi) = 0$$

$$L = \Omega_0 r_A^2 \quad (\text{specific angular momentum})$$

$$V_\phi = \Omega_0 r (M_A^2 (r_A/r)^2 - 1) / (M_A^2 - 1)$$

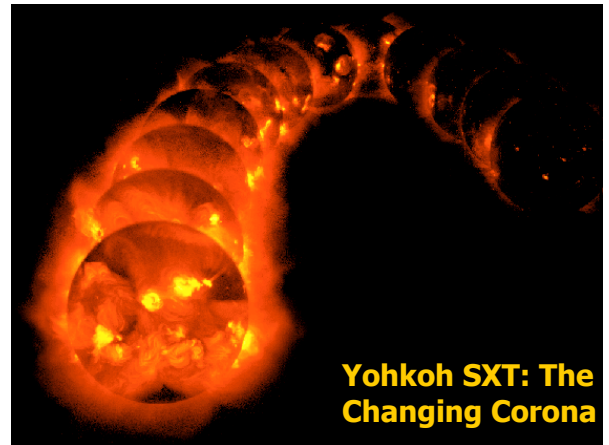
$$M_A = V_r (4\pi\rho)^{1/2} / B_r$$

Alfvén Machnumber

Weber & Davis, ApJ, 148, 217, 1967

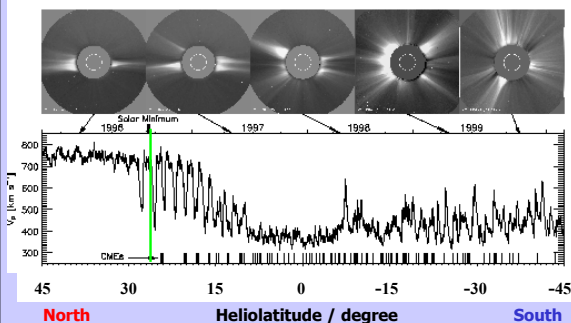
Helios:  $r_A = 10-20 R_\odot$

## Sun's loss of angular momentum carried by the solar wind II



Yohkoh SXT: The Changing Corona

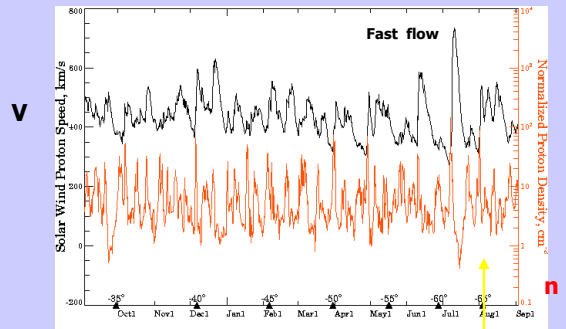
## Changing corona and solar wind



McComas et al., 2000

LASCO/Ulysses

## New solar wind data from Ulysses



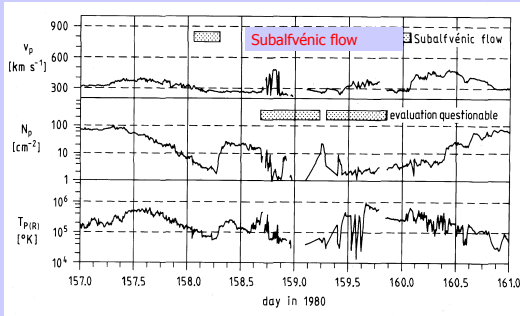
McComas et al., 2000

September 3, 1999 - September 2, 2000

Latitude: -65°



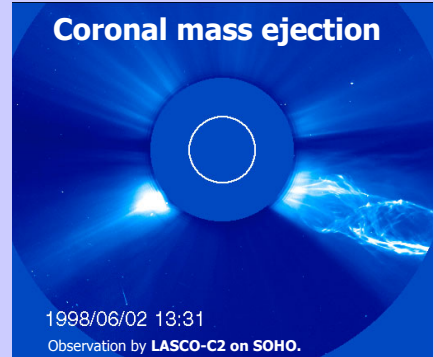
## Solar wind dropout



Schwenn, 1980

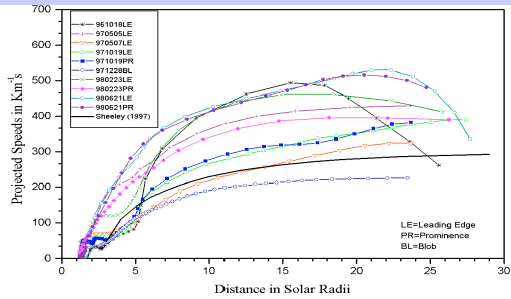
Helios 1 at 0.3 AU

## Coronal mass ejection



Note the helical structure of the prominence filaments!

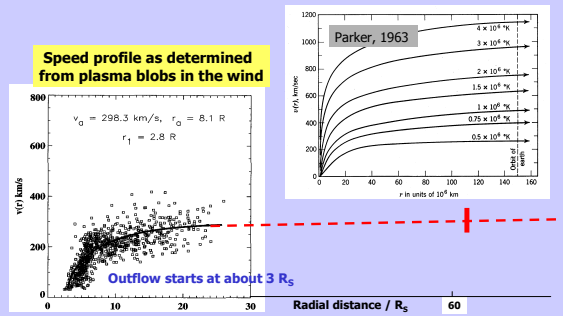
## Speed profile of balloon-type CMEs



Srivastava et al., 1999

Wide range of initial acceleration: 5-25  $\text{ms}^{-2}$

## Speed profile of the slow solar wind

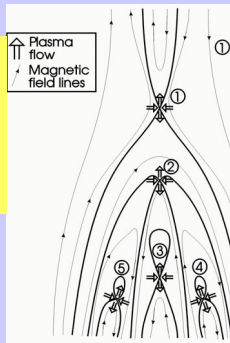


Sheeley et al., Ap.J., 484, 472, 1998

Consistent with Helios data

## Non-stationary slow solar wind

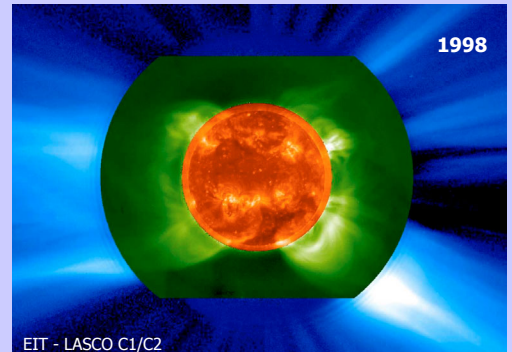
"...small eruptions at the helmet streamer cusp may incessantly accelerate small amounts of plasma without significant changes of the equilibrium configuration and might thus contribute to the non-stationary slow solar wind...."



- Acceleration of slow wind above cusp (1)
- Coronal eruptions by magnetic reconnection inside the streamer (1)
- Interaction of three smaller streamers forming a dome (2)
- Plasmoids form by reconnection (3,4,5)

Wiegmann et al., Solar Phys., 1999

## Corona of the active sun



## Solar wind models I

Assume heat flux,  $Q_e = -\rho\kappa\nabla T_e$ , is free of divergence and thermal equilibrium:  $T = T_p = T_e$ . Heat conduction:  $\kappa = \kappa_0 T^{5/2}$  and  $\kappa_0 = 8 \cdot 10^8$  erg/(cm s K); with  $T(\infty) = 0$  and  $T(0) = 10^6$ K and for spherical symmetry:

$$4\pi r^2 \kappa(T) dT/dr = \text{const} \quad \rightarrow \quad T = T_0 (R/R_0)^{2/7}$$

Density:  $\rho = n_p m_p + n_e m_e$ , quasi-neutrality:  $n = n_p = n_e$ , thermal pressure:  $p = n_p k_B T_p + n_e k_B T_e$ , then with hydrostatic equilibrium and  $p(0) = p_0$ :

$$dp/dr = -GMm_p n/r^2$$

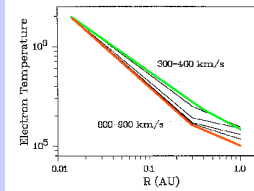
$$p = p_0 \exp\left[\frac{7GMm_p}{5k_B T_0 R} \left(\left(\frac{R}{R_0}\right)^{5/7} - 1\right)\right]$$

Problem:  $p(\infty) > 0$ , therefore corona must expand!

Chapman, 1957

## Proton and electron temperatures

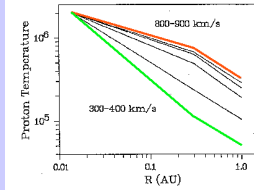
Electrons are cool!



slow wind

fast wind

Protons are hot!



fast wind

slow wind

Marsch, 1991

## Solar wind models II

Density:  $\rho = n_p m_p + n_e m_e$ , quasi-neutrality:  $n = n_p = n_e$ , ideal-gas thermal pressure:  $p = n_p k_B T_p + n_e k_B T_e$ , thermal equilibrium:  $T = T_p = T_e$ , then with hydrodynamic equilibrium:

$$mn_p v \frac{dv}{dr} = -dp/dr - GMm_p n/r^2$$

Mass continuity equation:

$$mn_p v r^2 = J$$

Assume an isothermal corona, with sound speed  $c_0 = (k_B T_0/m_p)^{1/2}$ , then one has to integrate the DE:

$$\left[\left(\frac{v}{c_0}\right)^2 - 1\right] \frac{dv}{v} = 2 \left(\frac{1-r_c/r}{r}\right) \frac{dr}{r}$$

With the critical radius,  $r_c = GMm_p/(2k_B T_0) = (V_\infty/2c_0)^2$ , and the escape speed,  $V_\infty = 618$  km/s, from the Sun's surface.

Parker, 1958

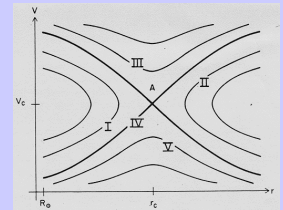
## Solar wind models III

Introduce the sonic Mach number as,  $M_s = v/c_0$ , then the integral of the DE (C is an integration constant) reads:

$$(M_s)^2 - \ln(M_s)^2 = 4 \left( \ln(r/r_c) + r_c/r \right) + C$$

For large distances,  $M_s \gg 1$ ; and  $v \sim (\ln r)^{1/2}$ , and  $n \sim r^{-2}/v$ , reflecting spherical symmetry.

Only the „wind“ solution IV, with  $C=-3$ , goes through the critical point  $r_c$  and yields:  $n \rightarrow 0$  and thus  $p \rightarrow 0$  for  $r \rightarrow \infty$ . This is Parker's famous solution: the solar wind.



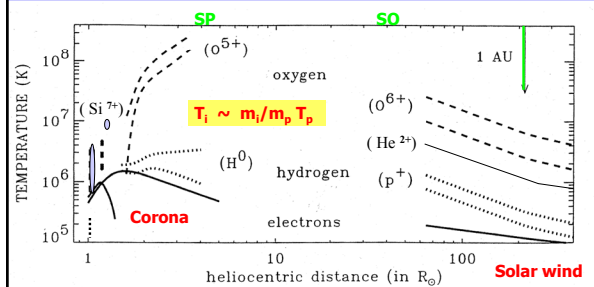
$V_\infty$ , solar breeze; III accretion flow

Parker, 1958

## Fluid equations

- Mass flux:  $F_M = \rho v A$      $\rho = n_p m_p + n_i m_i$
- Magnetic flux:  $F_B = B A$
- Total momentum equation:  $v \frac{dv}{dr} = -\frac{1}{\rho} \frac{dp}{dr} - \frac{GM_s}{r^2} + a_w$
- Thermal pressure:  $p = n_p k_B T_p + n_e k_B T_e + n_i k_B T_i$
- MHD wave pressure:  $p_w = (\delta B)^2 / (8\pi)$
- Kinetic wave acceleration:  $a_w = (\rho_p a_p + \rho_i a_i) / \rho$
- Stream/flux-tube cross section:  $A(r)$

## Temperature profiles in the corona and fast solar wind



Cranmer et al., Ap.J., 2000; Marsch, 1991

## Energy equations

Parallel thermal energy

$$\frac{d}{dr} v_{\parallel j}^2 = -2v_{\parallel j}^2 \left( \frac{1}{u_j} \frac{du_j}{dr} \right) + \frac{2q_{\parallel j}}{u_j} + (Q_{\parallel j} + S_{\parallel j})/u_j$$

w-p terms + sources + sinks

Perpendicular thermal energy

$$\frac{d}{dr} v_{\perp j}^2 = -v_{\perp j}^2 \left( \frac{1}{A} \frac{dA}{dr} \right) + \frac{q_{\perp j}}{u_j} + (Q_{\perp j} + S_{\perp j})/u_j$$

Heating functions:  $q_{\perp j}$  .... ?

Wave energy absorption/emission by wave-particle interactions !

Conduction/collisional exchange of heat + radiative losses

## Heating and acceleration of ions by cyclotron and Landau resonance

$$\begin{aligned} \left( \begin{array}{l} \frac{\partial U_{j\parallel}}{\partial t} \\ \frac{\partial v_{j\parallel}^2}{\partial t} \\ \frac{\partial v_{j\perp}^2}{\partial t} \end{array} \right) &= a_j \quad \text{acceleration} \\ &= 2 q_{\parallel j} \quad \text{parallel heating} \\ &= q_{\perp j} \quad \text{perpendicular heating} \end{aligned}$$

$$= \frac{1}{(2\pi)^3} \int_{-\infty}^{+\infty} d^3k \sum_M \mathcal{E}_M(k) \left( \frac{\Omega_j}{k} \right)^2 \frac{1}{1 - |\mathbf{k} \cdot \mathbf{e}_M(k)|^2}$$

$$\times \sum_{s=-\infty}^{+\infty} \mathcal{R}_j(k, s) \left( \begin{array}{l} k_{\parallel} \\ 2k_{\parallel} w_j(k, s) \\ s\Omega_j \end{array} \right)$$

Marsch and Tu, JGR, 106, 227, 2001

Wave spectrum ?  
Wave dispersion ?  
Resonance function ?

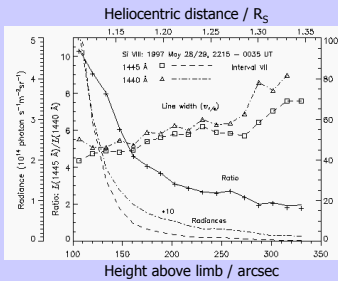
## Height profile of turbulence amplitude

SUMER

Silicon VIII

$\lambda\lambda 1440, 1445$

North polar coronal hole



Doppler velocity

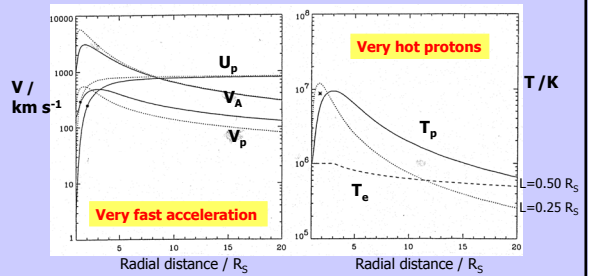
$V_{1/e} / \text{km s}^{-1}$

At 1.33  $R_{\odot}$ :

$T_{\text{Si}} \approx 10^7 \text{ K}$   
 $\xi \approx 70 \text{ km s}^{-1}$   
 $n_e \approx 10^6 \text{ cm}^{-3}$

Wilhelm et al., Ap.J., 500, 1023, 1998

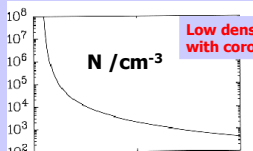
## Rapid acceleration of the high-speed solar wind



McKenzie et al., A&A, 303, L45, 1995

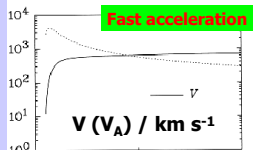
Heating:  $Q = Q_0 \exp(-(r - R_{\odot})/L)$ ; Sonic point:  $r \approx 2 R_{\odot}$

## Model of the fast solar wind

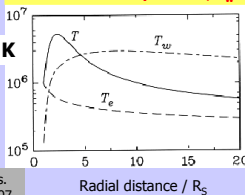


Low density,  $n \approx 10^8 \text{ cm}^{-3}$ , consistent with coronagraph measurements

- hot protons,  $T_{\text{max}} \approx 5 \text{ M K}$
- cold electrons
- small wave temperature,  $T_w$



T / K



McKenzie et al., Geophys. Res. Lett., 24, 2877, 1997

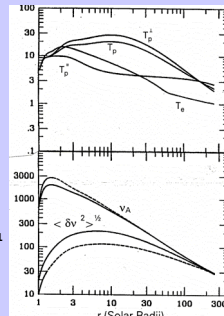
## Anisotropic two-fluid model of the fast solar wind

$T_{\perp 1} = 3 \cdot 10^6 \text{ K}$

T / 10<sup>5</sup> K

- Anisotropic heat deposition in 1-D two-fluid model
- Alfvén wave pressure gradient

v / km s<sup>-1</sup>



- Anisotropy weakly influences dynamics
- Anisotropy needed for perpendicular ion-cyclotron heating and thermodynamics

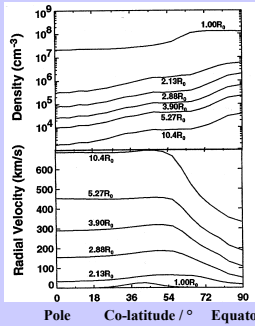
Coronal base:  
 $\delta v \approx 10\text{-}20 \text{ km s}^{-1}$   
 $\xi \approx 20\text{-}30 \text{ km s}^{-1}$

Hu et al., JGR, 102, 14661, 1997



## Two-dimensional two-fluid MHD model of the solar corona

- Time-dependent 2-D model MHD with separate  $T_e$  and  $T_p$  equations
- Slow outflow at equator, fast over poles after 1 day
- Heating functions  $Q_e$  and  $Q_p$  latitude-dependent



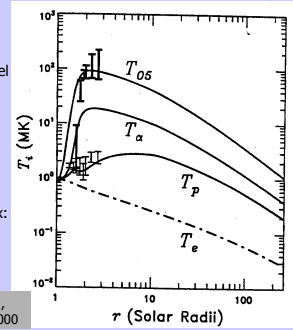
Heating function:  
 $Q_{e,p} = Q_0 f_{ep}(r, \theta) \exp(-0.1(r-R_s)/R_s)$   
 $Q_0 = 5 \cdot 10^{-8} \text{ erg cm}^{-3} \text{ s}^{-1}$

**Poles:**  
 $T_e < T_p \quad \beta \ll 1$   
**Equator:**  
 $T_e = T_p \quad \beta \geq 1$

Suess et al., JGR, **104**, 4697, 1999

## Four-fluid model for turbulence driven heating of coronal ions

- Four-fluid 1-D corona/wind model
- Quasi-linear heating and acceleration by dispersive ion-cyclotron waves
- Rigid power-law spectra with index:  $-2 \leq \gamma \leq -1$

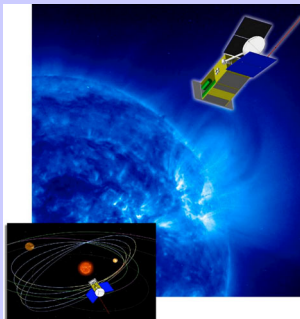


- No wave absorption
- Turbulent spectra not self-consistent

**Preferential heating of heavy ions by waves**

Hu, Esser & Habbal, JGR, **105**, 5093, 2000

## The future: Solar Orbiter

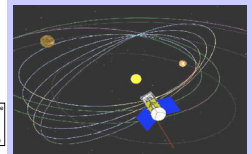
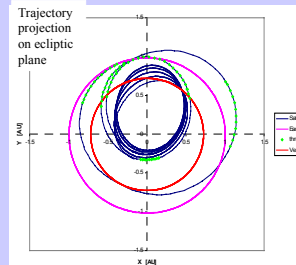


**A high-resolution mission to the Sun and inner heliosphere**

ESA

2011 - 2012

## Solar Orbiter's novel orbital design



- Closer to the Sun!
- Out of the ecliptic!

- Venus gravity assist
- Solar electric propulsion