

Vorlesung "Physik des Sonnensystems"
Univ. Göttingen, 2. Juni 2008

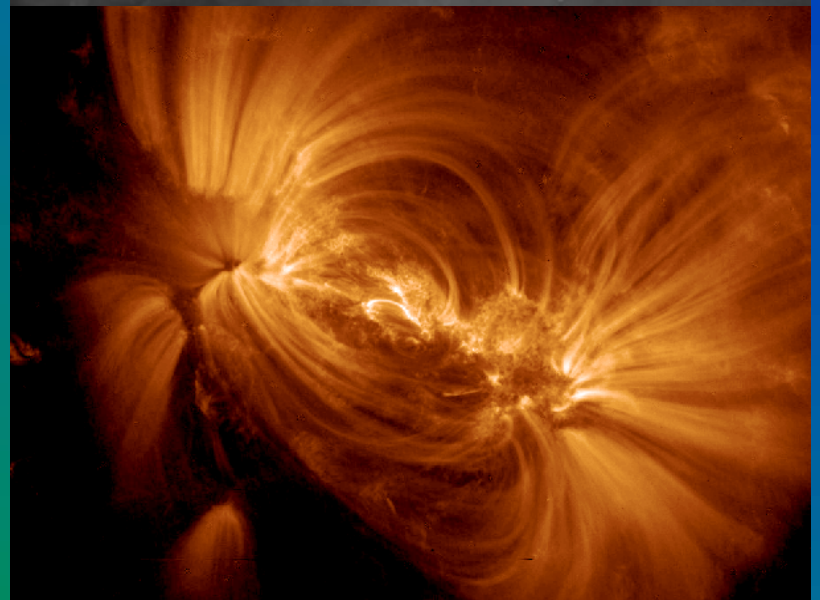
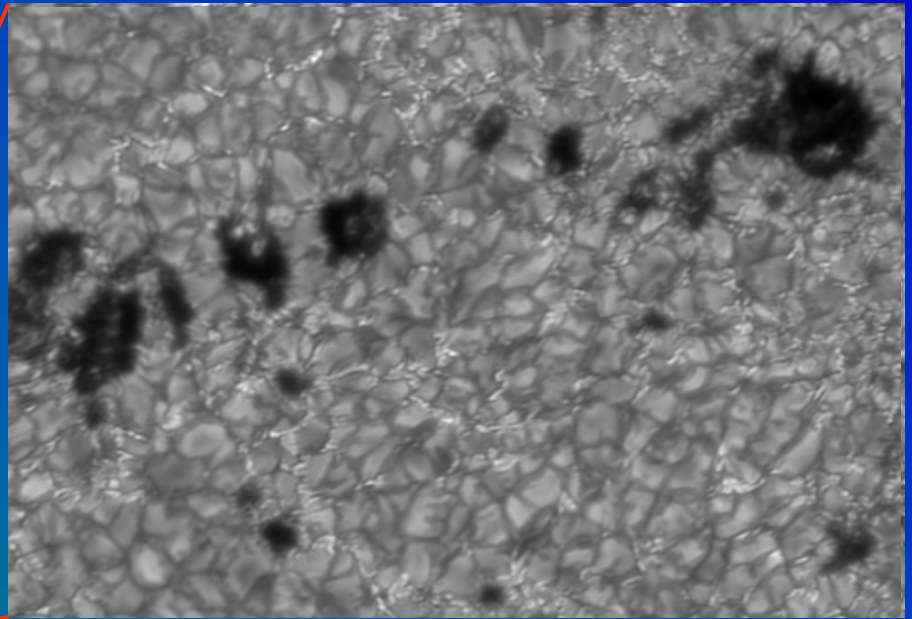
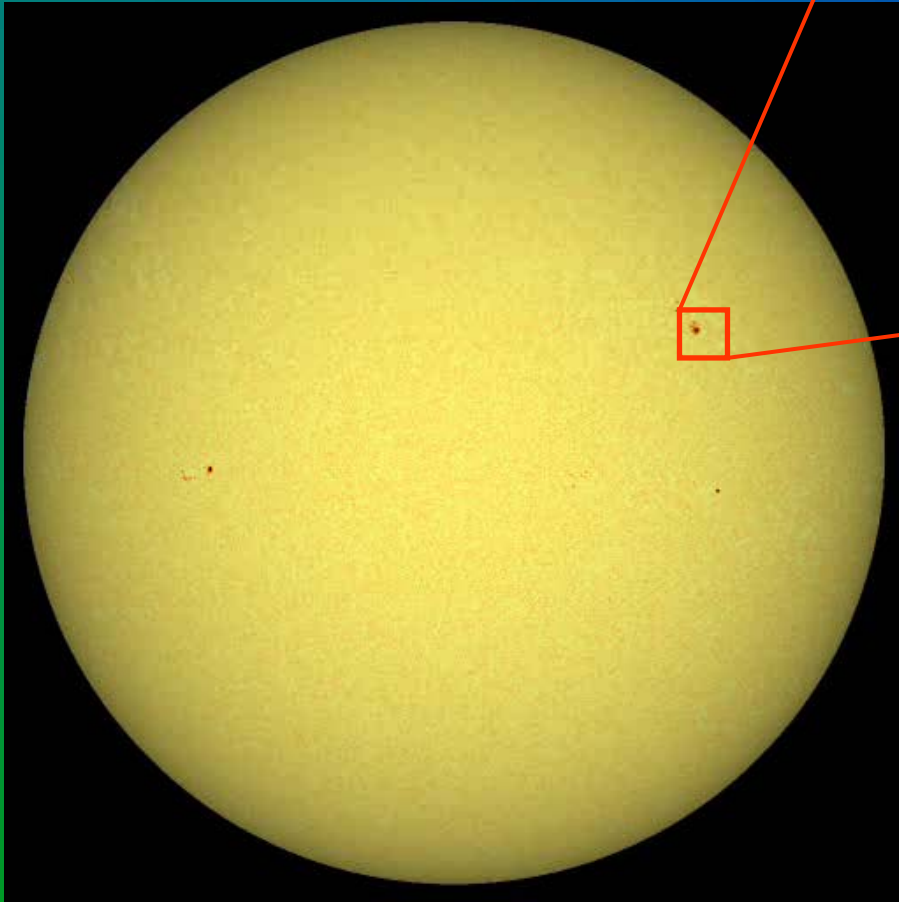
Konvektion und solares Magnetfeld

Manfred Schüssler

Max-Planck-Institut für Sonnensystemforschung

Katlenburg-Lindau

Convection & magnetism: closely related

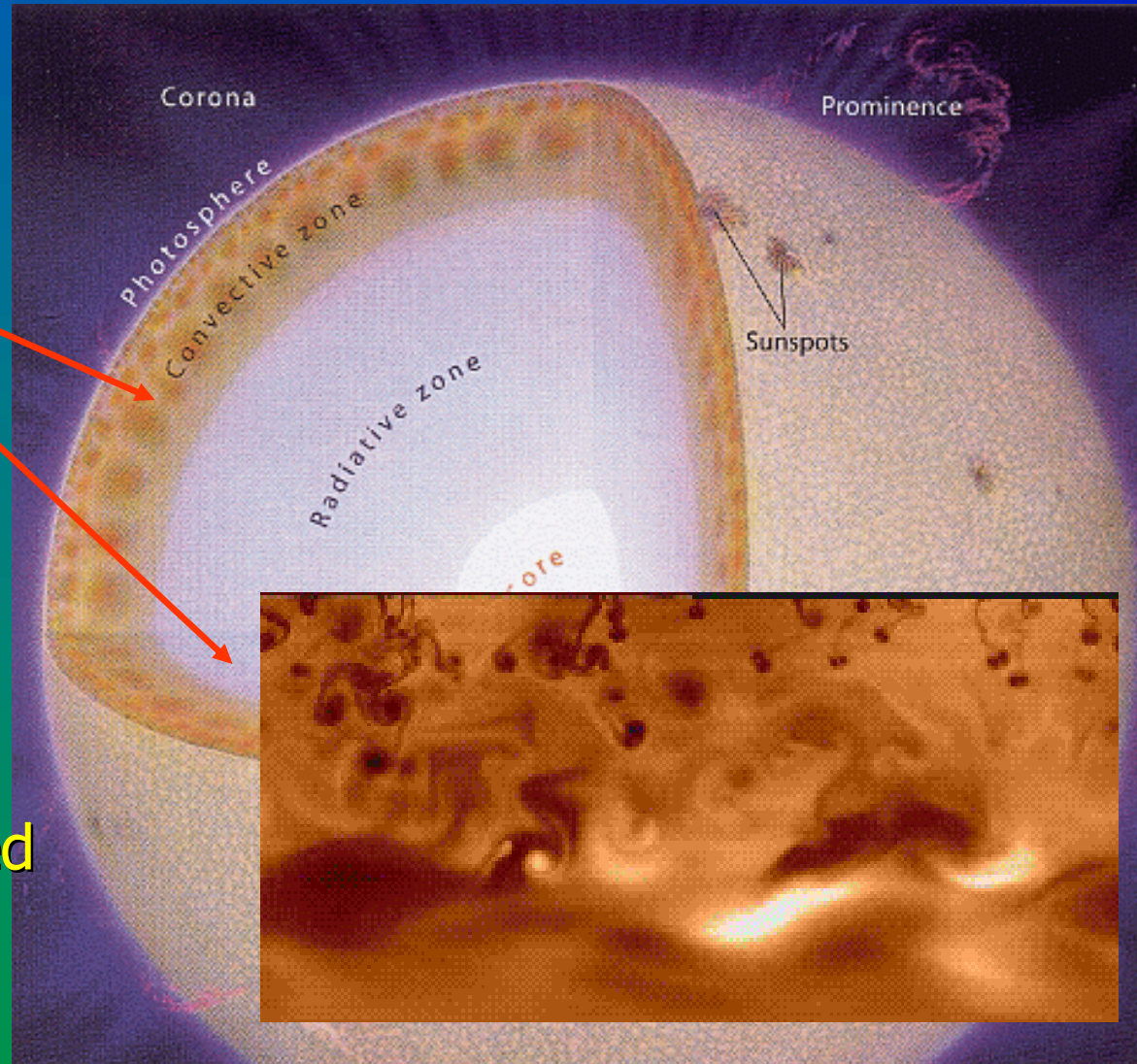


Outline

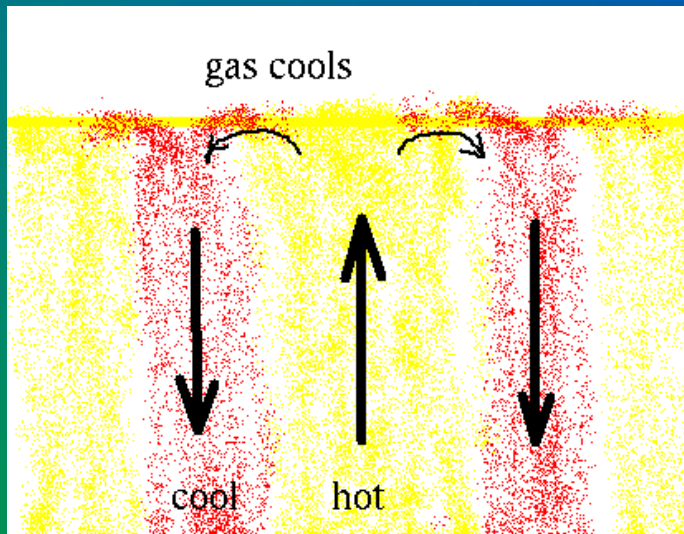
- 1) Basic physics of convection
- 2) Numerical simulation of convection
- 3) Overview of solar magnetism
- 4) Surface magneto-convection
- 5) Deep convection zone field & dynamo

The solar convection zone

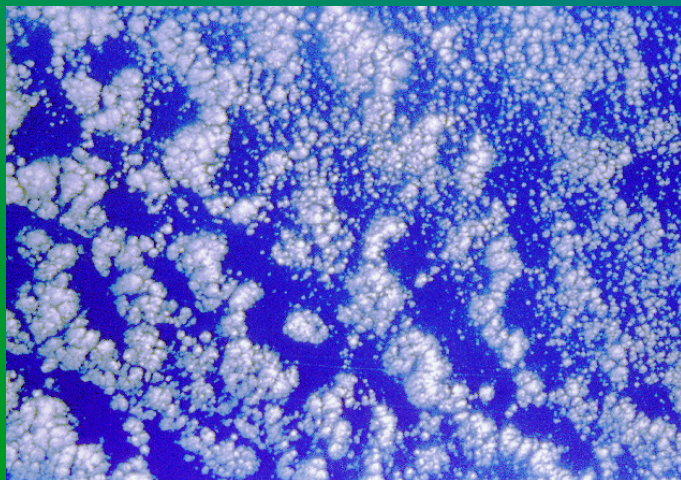
- 200 Mm thick layer in turbulent motion
- Velocities range from 100 m/s (bottom) to 10 km/s (top)
- Energy flux nearly completely transported by convective motion



What is convection?

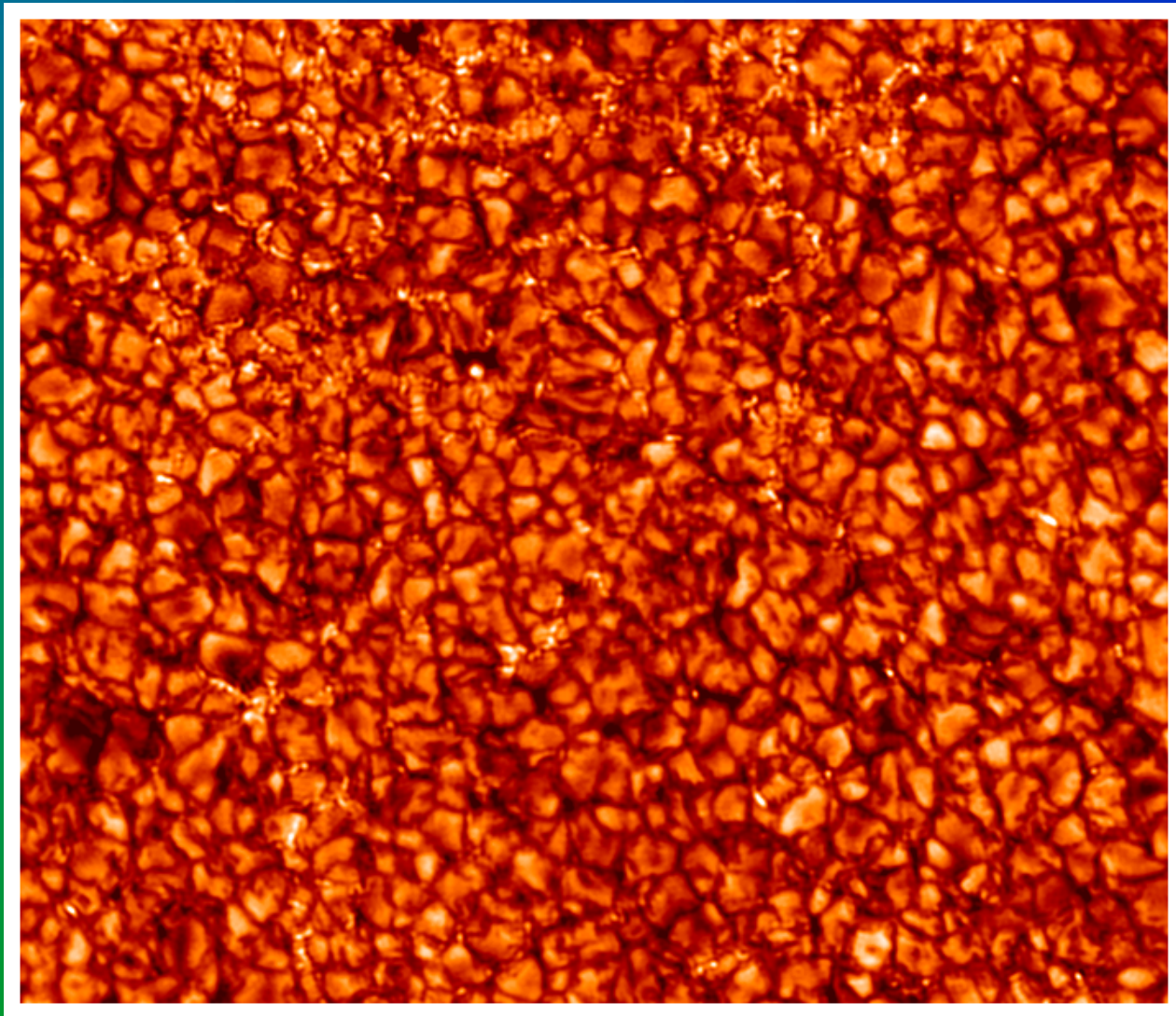


- Flow driven by thermal buoyancy
- Convective instability

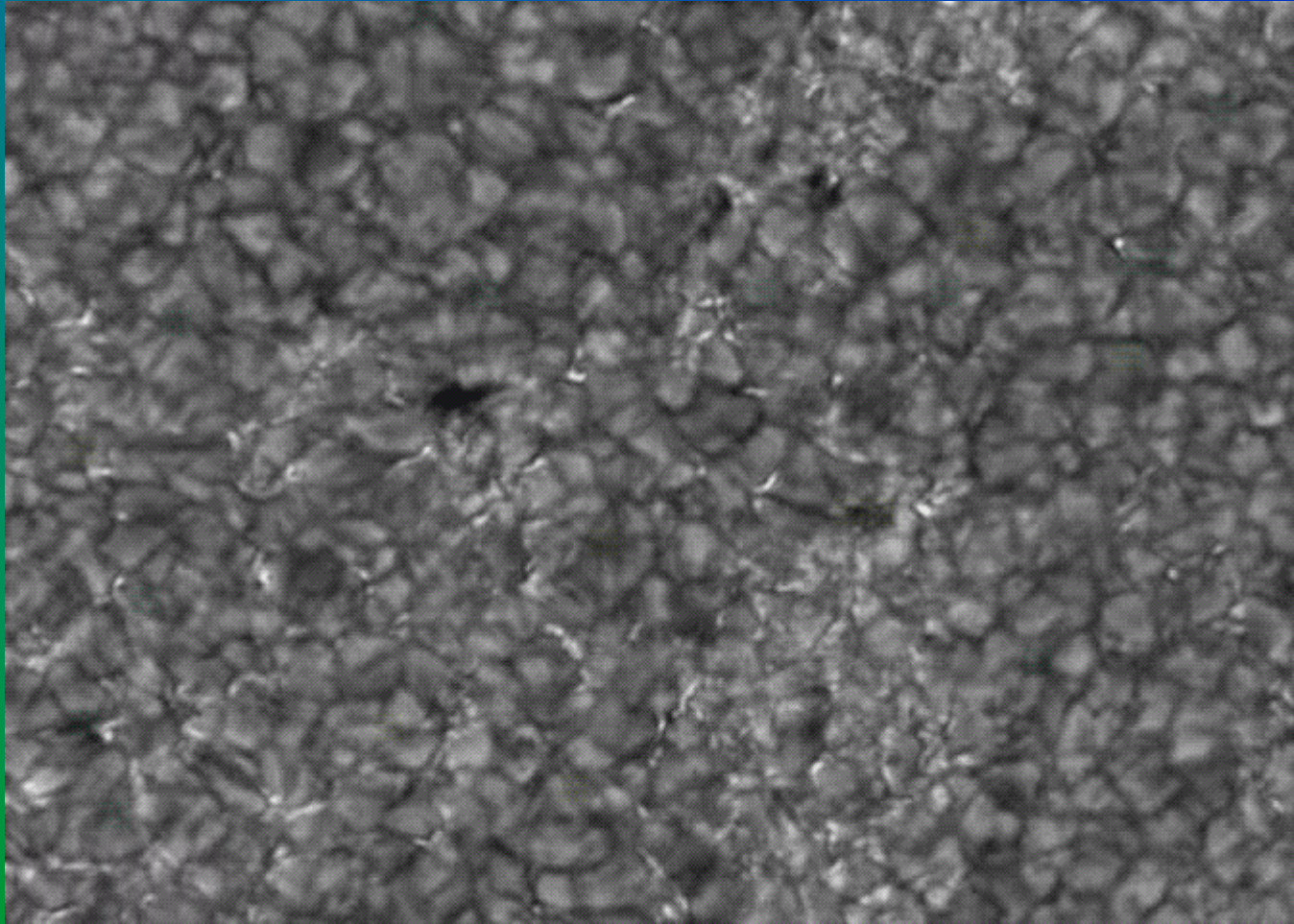


→ Viewgraphs...

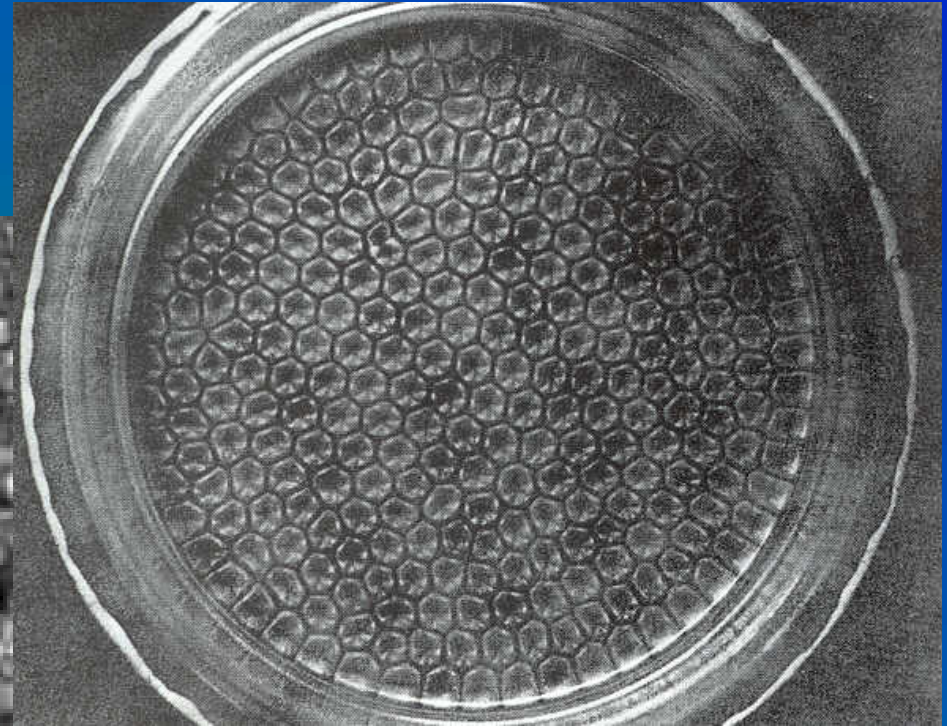
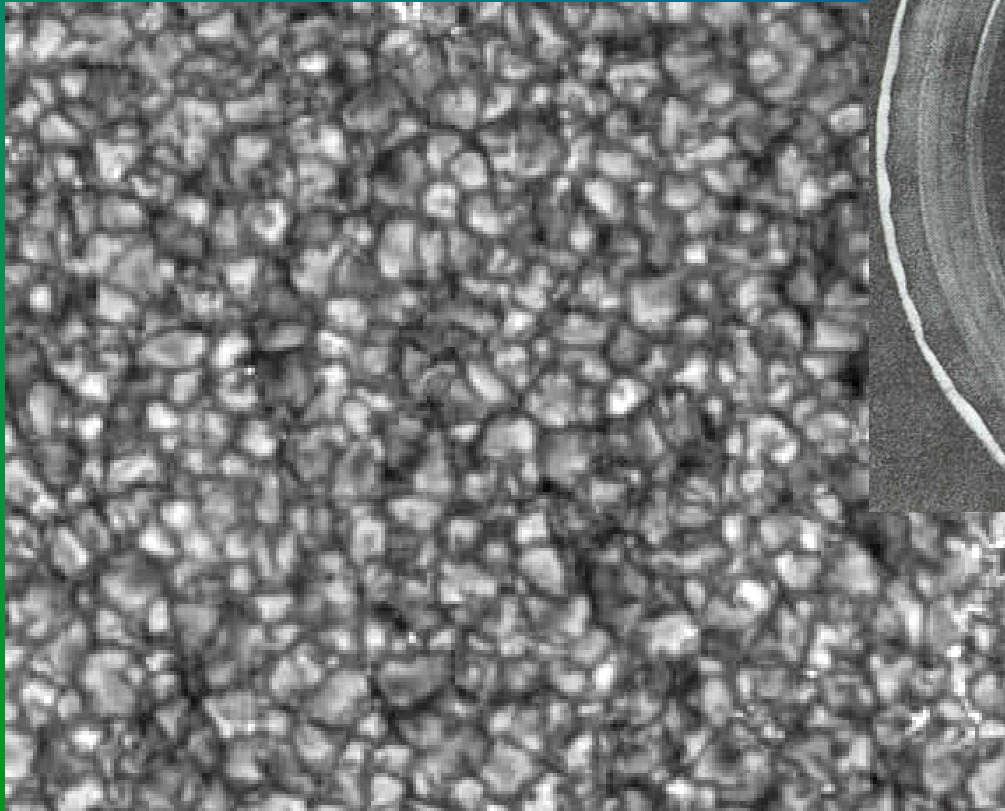
Granulation: Solar surface convection



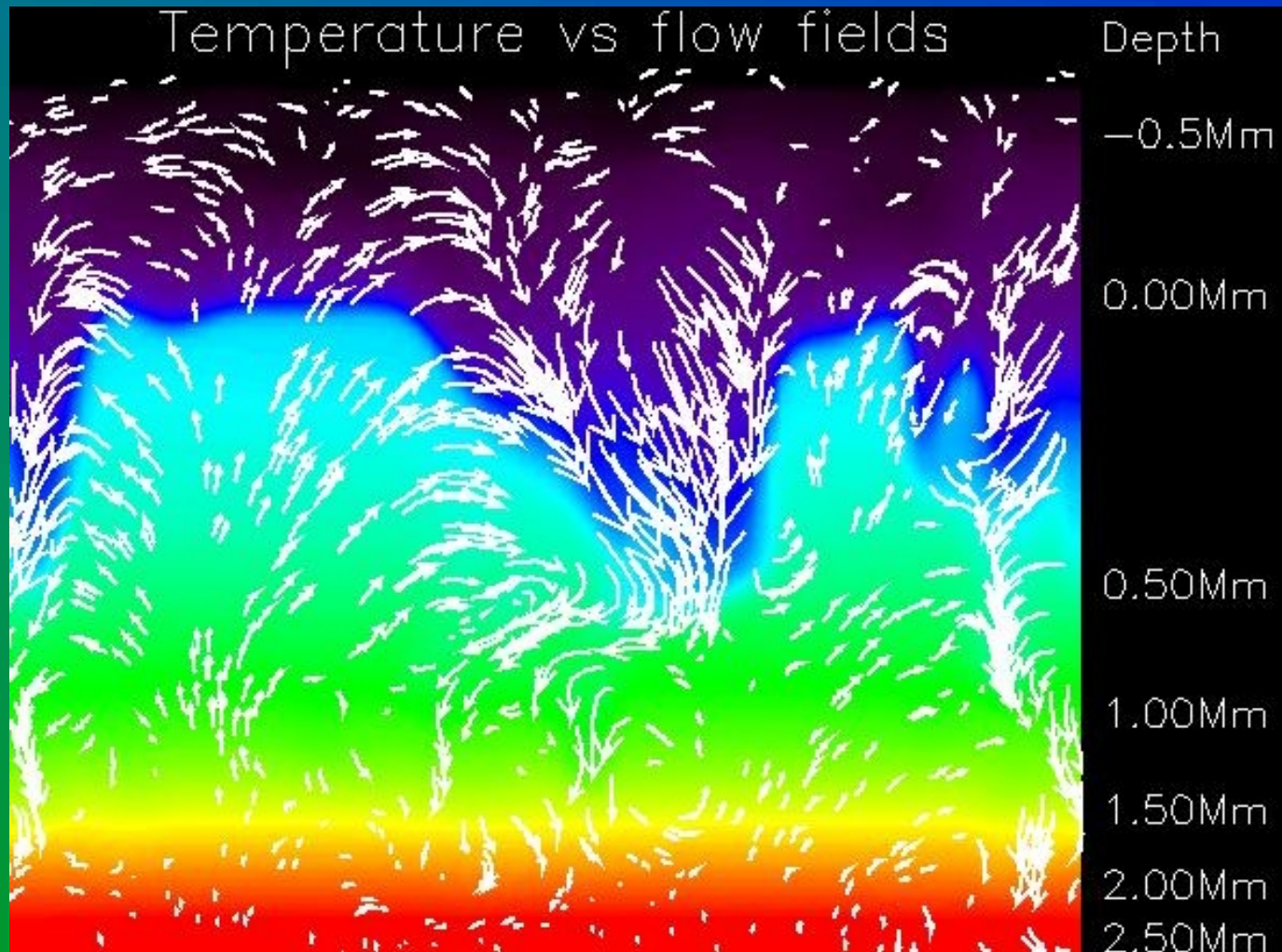
Solar granulation



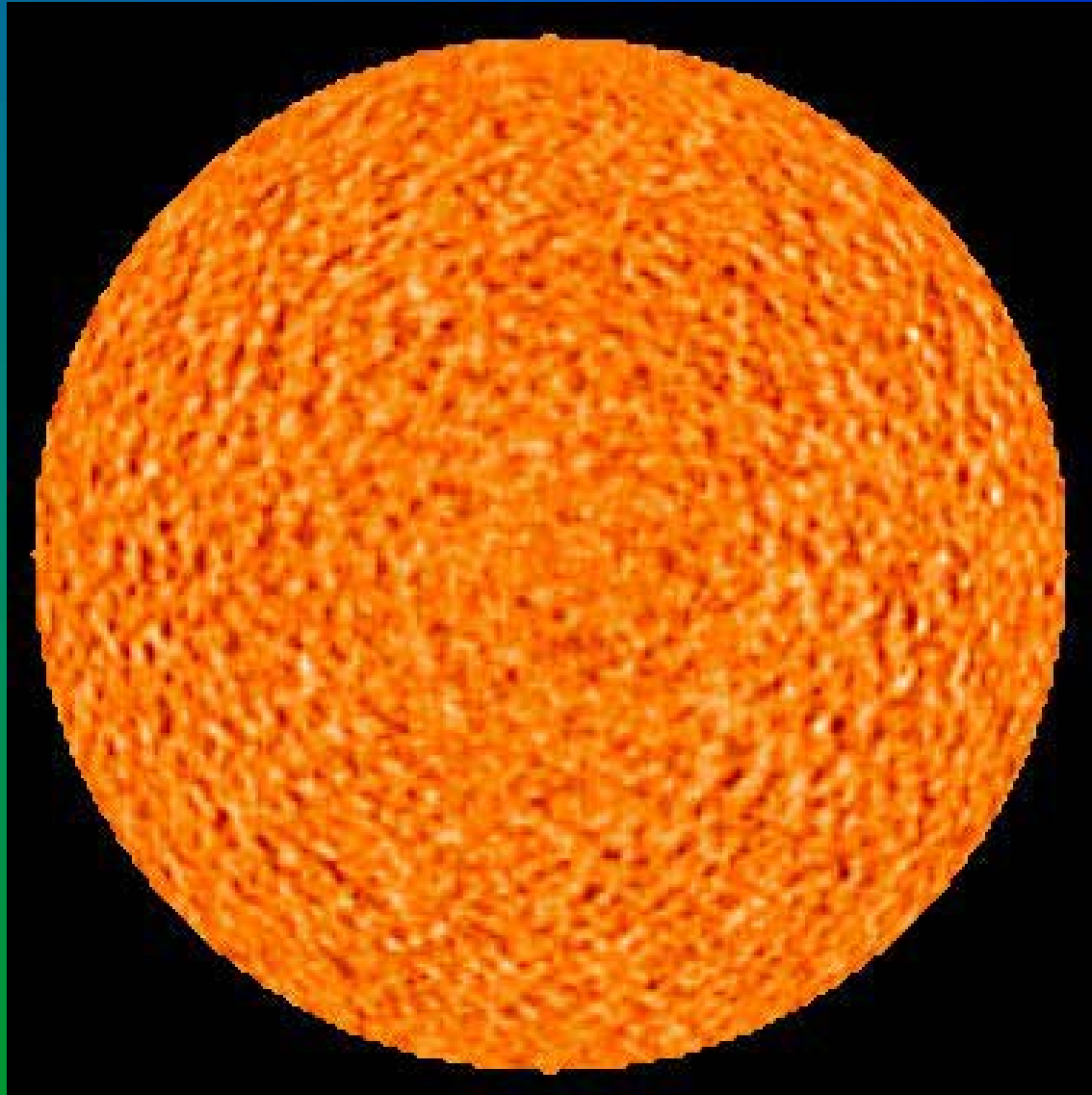
Granulation und laboratory convection



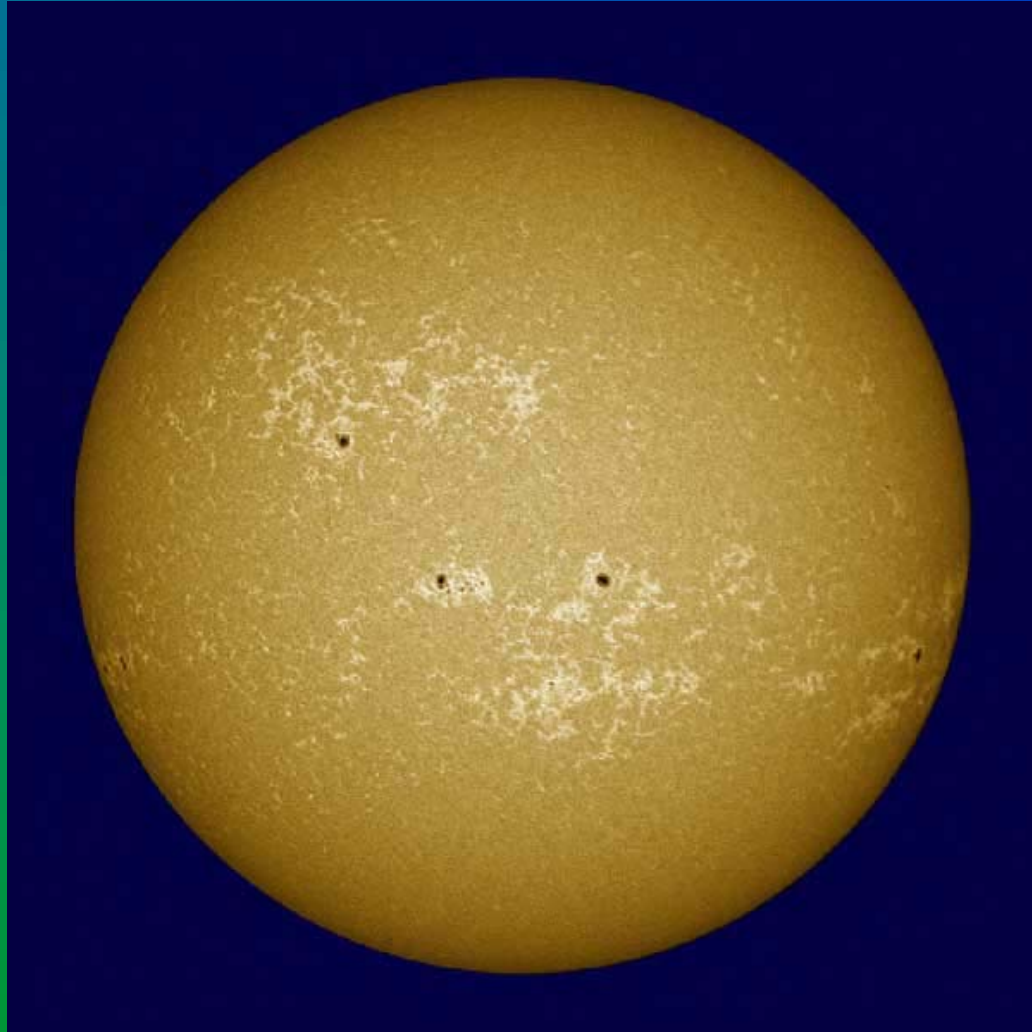
Granulation as a convective phenomenon



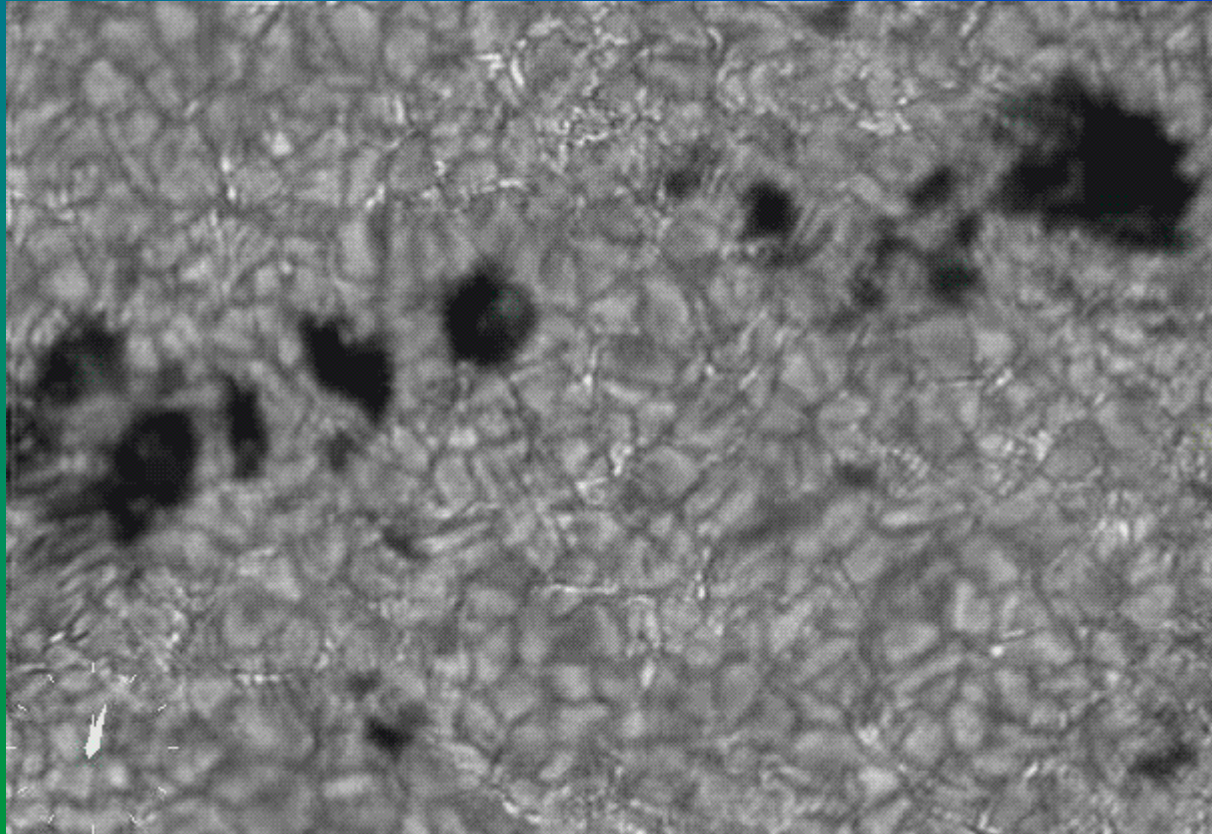
Supergranulation



Supergranulation and magnetic field: the Ca^+ network



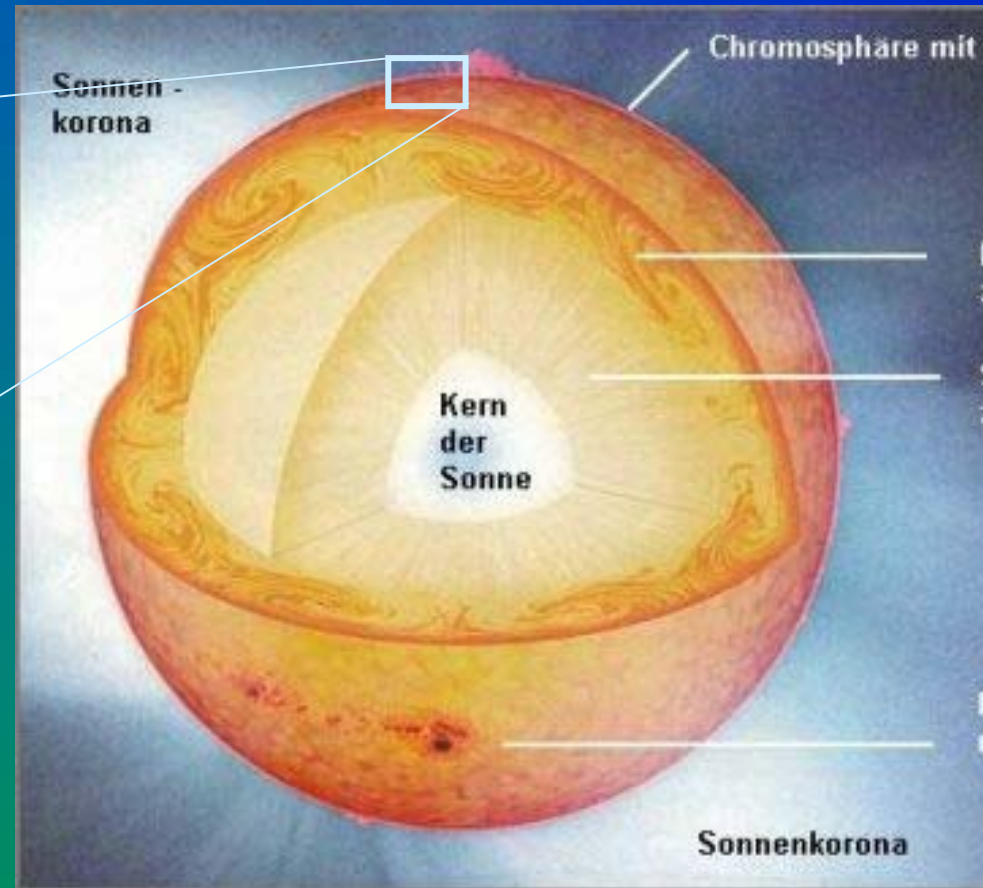
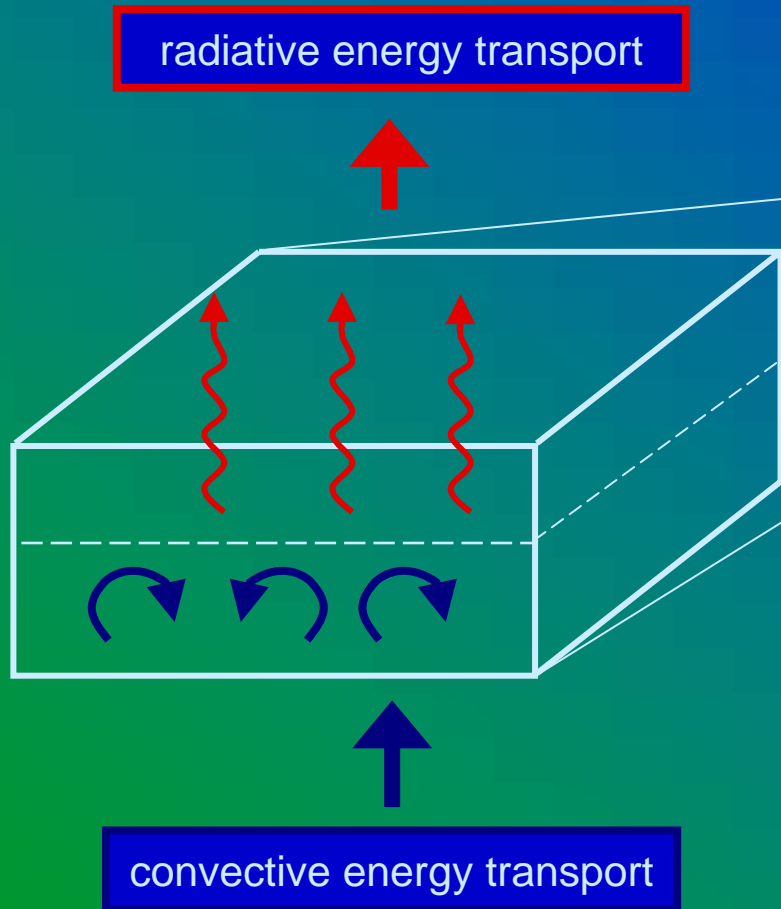
Granulation, sunspots, & small-scale magnetic field



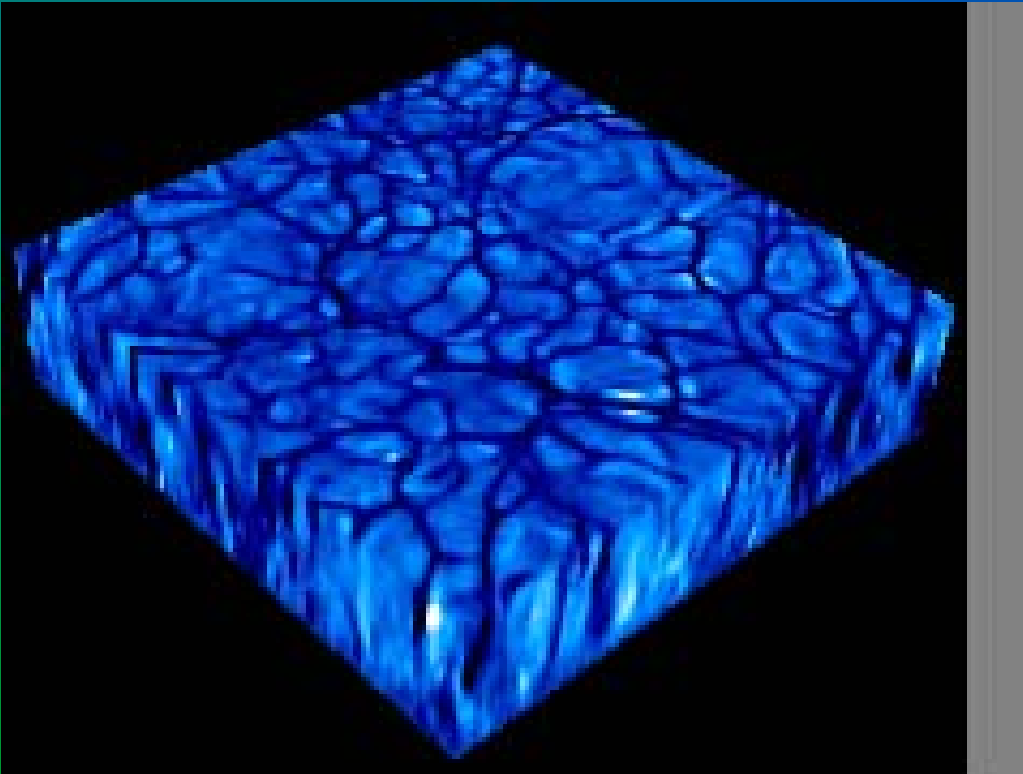
'Realistic' solar simulations

- elaborate physics: partial ionization, radiation, compressible, open box, transmitting boundaries, spectral line diagnostics (Stokes profiles)
- + : approximation to solar conditions
- + : direct comparison with observations
- – : computational restrictions (box size, resolution)
- – : Reynolds numbers much below solar values

Approach: Local simulation box including photosphere



Computer-simulated convection



- Boussinesq model
- Rayleigh number: $5 \cdot 10^5$
- 3D, $512 \times 512 \times 97$ mesh
- wide box, aspect ratio: 10
- “(meso)granulation”

Cattaneo & Emonet (2001)

The MURaM code: equations

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) = 0$$

Continuity equation

$$\frac{\partial \rho \mathbf{u}}{\partial t} + \nabla \cdot \left(\rho \mathbf{u} \mathbf{u} + \left(p + \frac{|\mathbf{B}|^2}{8\pi} \right) \mathbf{1} - \frac{\mathbf{B}\mathbf{B}}{4\pi} \right) = \rho \mathbf{g} + \nabla \cdot \underline{\underline{\tau}}$$

Momentum equation

$$\begin{aligned} \frac{\partial e}{\partial t} + \nabla \cdot \left(\mathbf{u} \left(e + p + \frac{|\mathbf{B}|^2}{8\pi} \right) - \frac{1}{4\pi} \mathbf{B}(\mathbf{u} \cdot \mathbf{B}) \right) \\ = \frac{1}{4\pi} \nabla \cdot (\mathbf{B} \times \eta \nabla \times \mathbf{B}) + \nabla \cdot (\mathbf{u} \cdot \underline{\underline{\tau}}) + \nabla \cdot (\chi \rho \nabla \frac{e}{\rho}) \\ + \rho(\mathbf{g} \cdot \mathbf{u}) - Q_{rad}, \end{aligned}$$

Energy equation

$$Q_{rad} = -\nabla \cdot \mathbf{F} = 4\pi\rho \int \kappa_\nu (J_\nu - S_\nu) d\nu$$

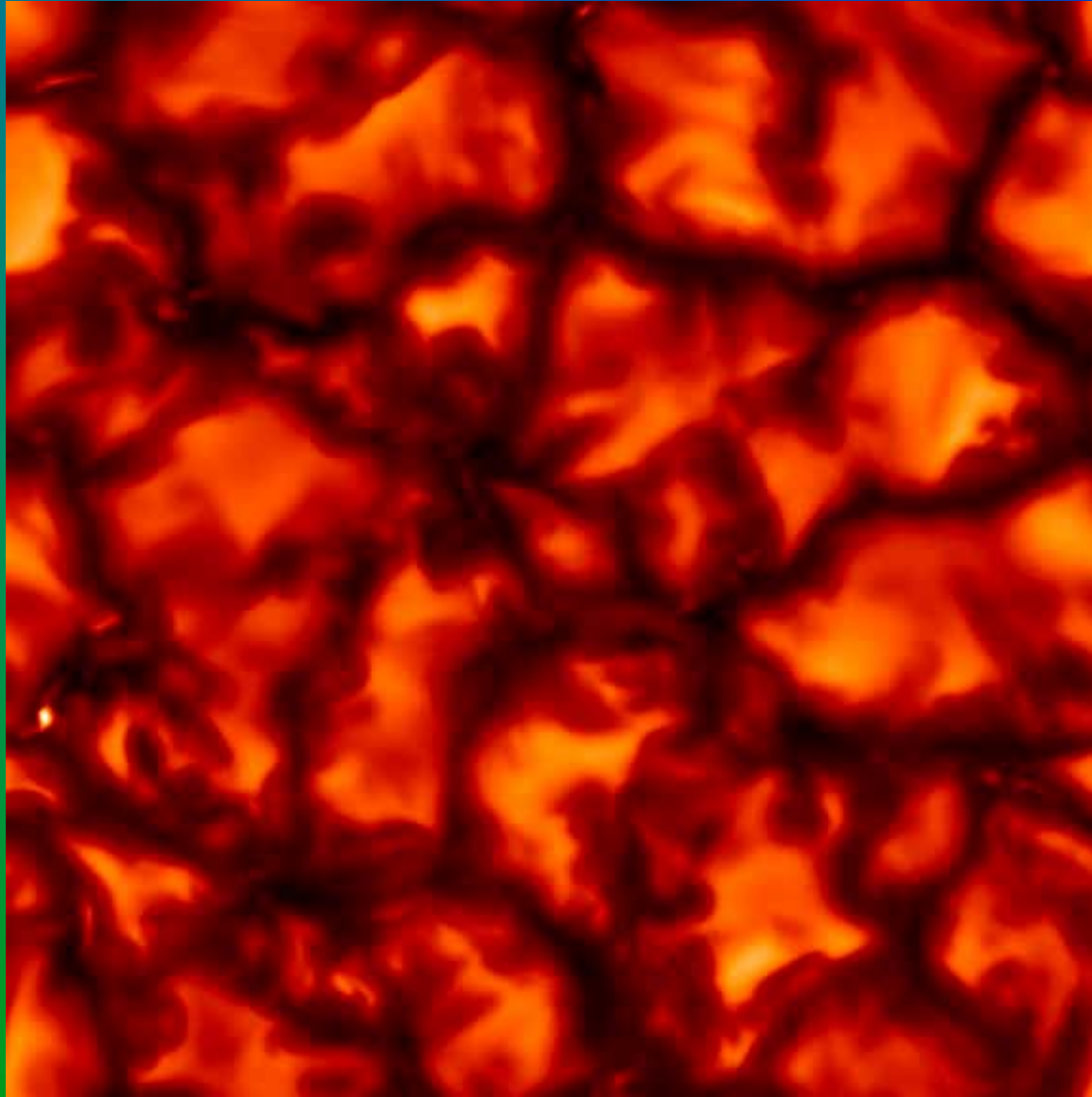
$$\frac{\partial \mathbf{B}}{\partial t} + \nabla \cdot (\mathbf{u}\mathbf{B} - \mathbf{B}\mathbf{u}) = -\nabla \times (\eta \nabla \times \mathbf{B}).$$

Induction equation

$$\frac{dI_\nu}{ds} = -\kappa_\nu \rho (I_\nu - S_\nu)$$

Radiative Transfer Equation

Computer-simulated convection

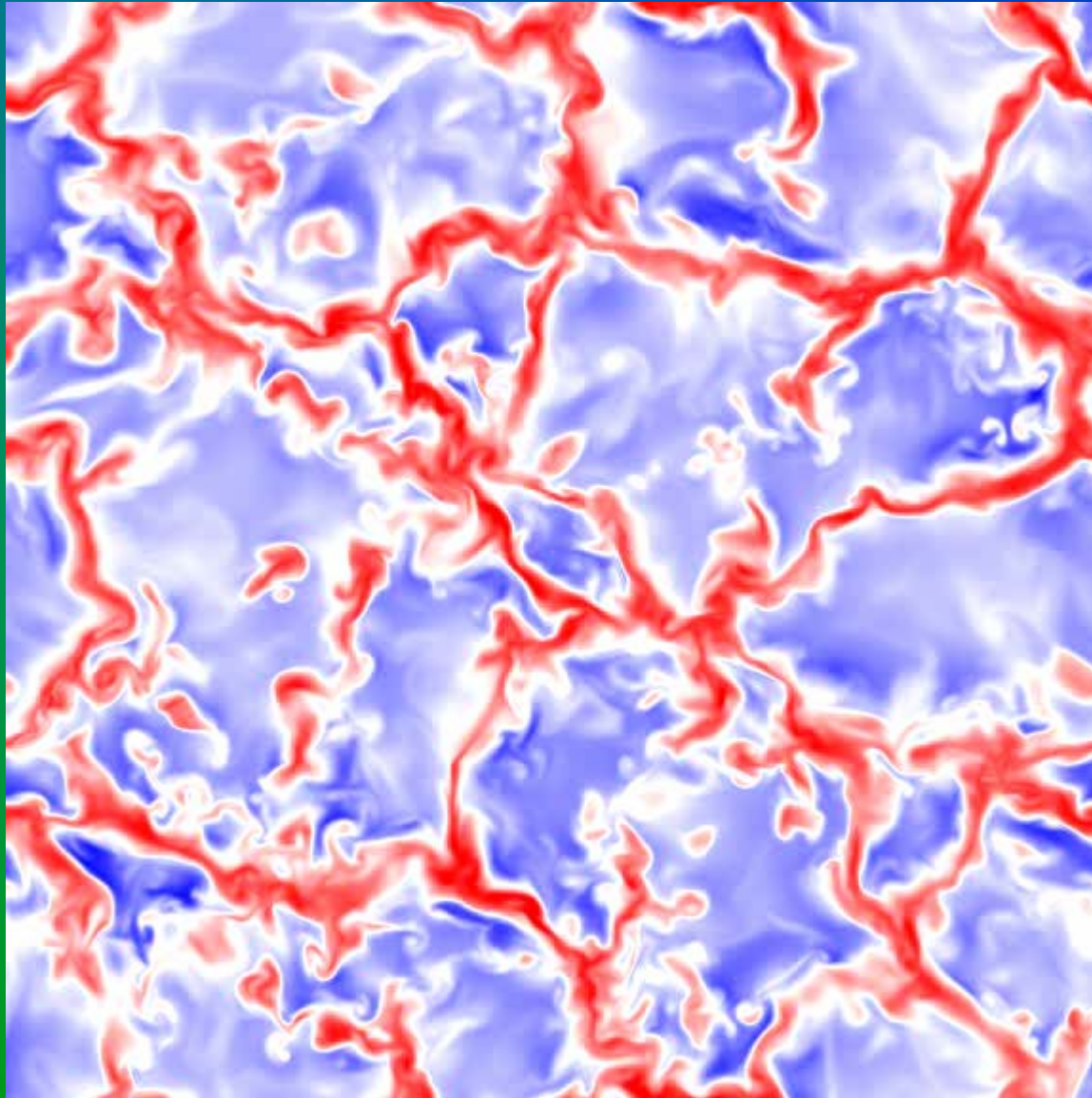


- realistic simulation
- ionization, rad. transfer
- 3D, 576×576×100 mesh
- 6 Mm × 6 Mm × 1.4 Mm
- granulation

Vögler et al. (2005)

Emerging intensity

Computer-simulated convection

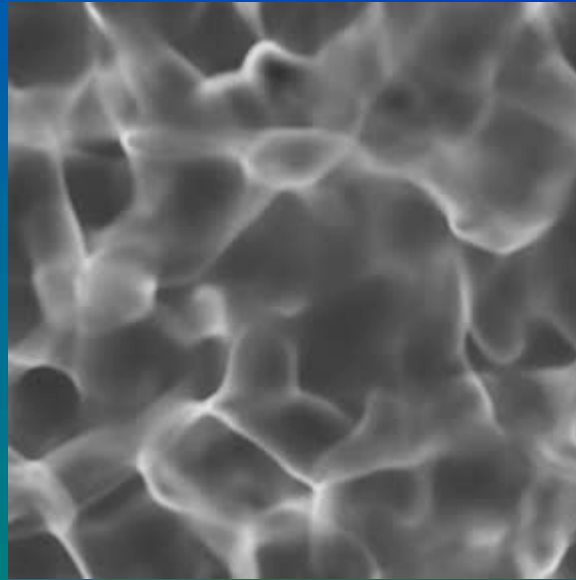
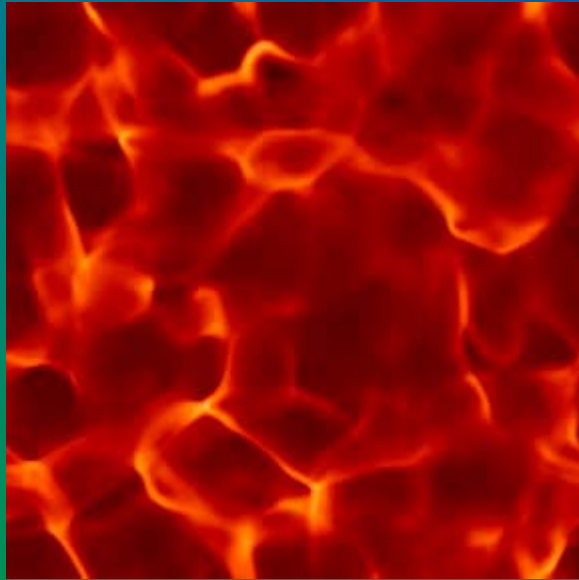


- realistic simulation
- ionization, rad. transfer
- 3D, 576×576×100 mesh
- 6 Mm × 6 Mm × 1.4 Mm
- granulation

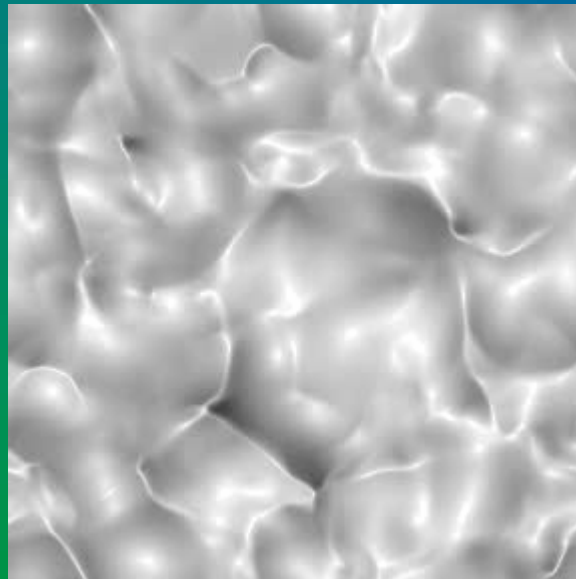
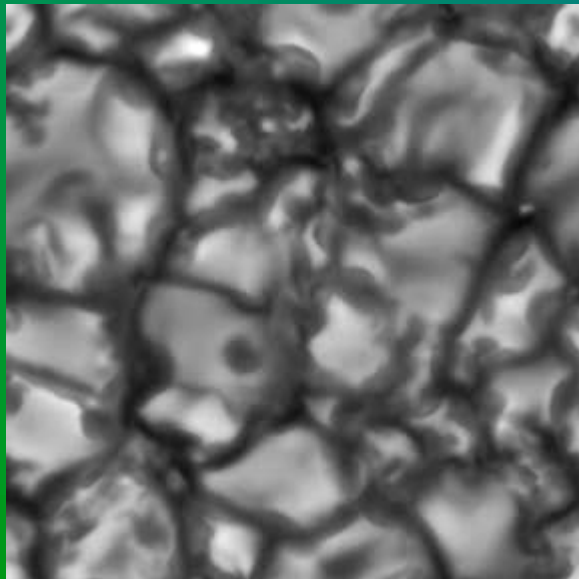
Vögler et al. (2005)

Vertical velocity (red: down, blue: up)

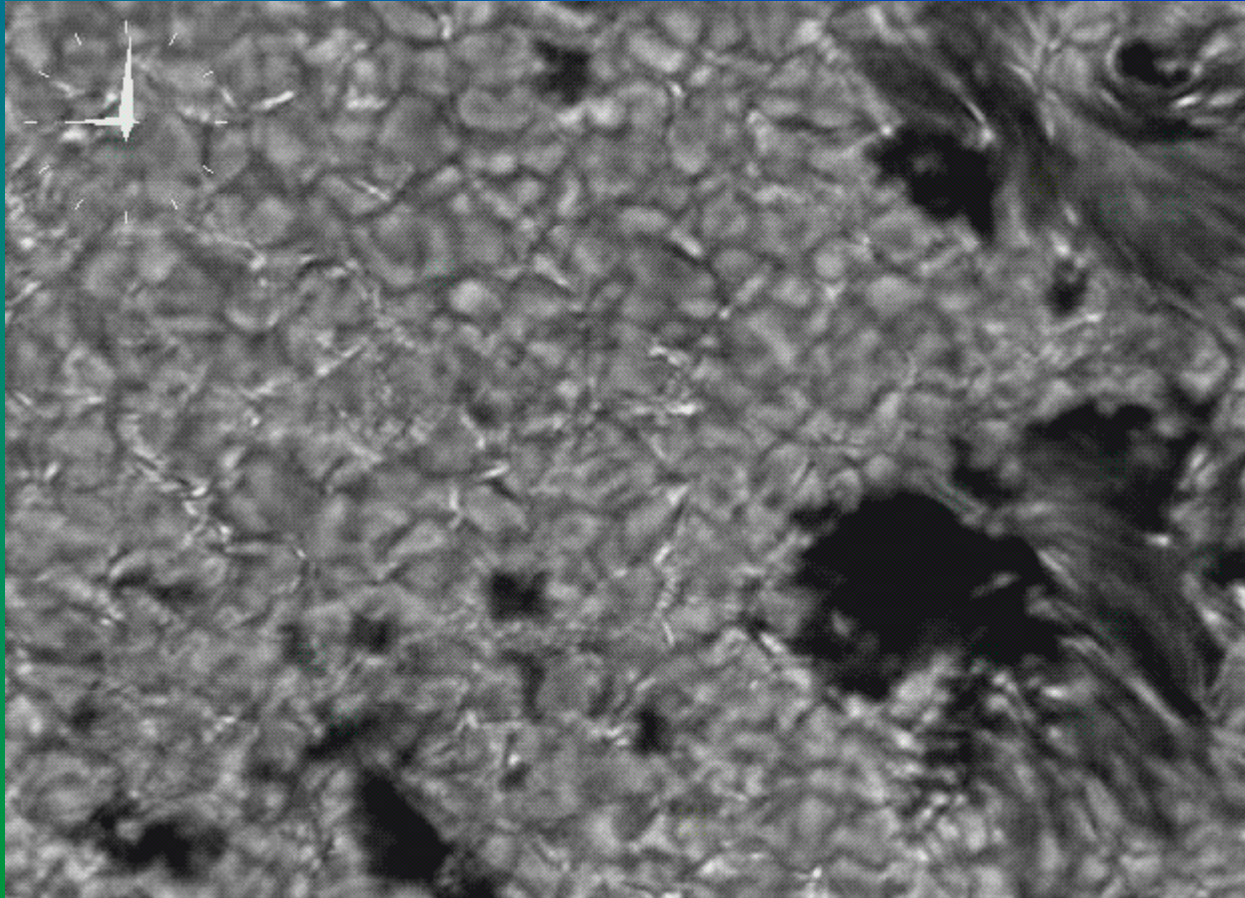
Computer-simulated convection



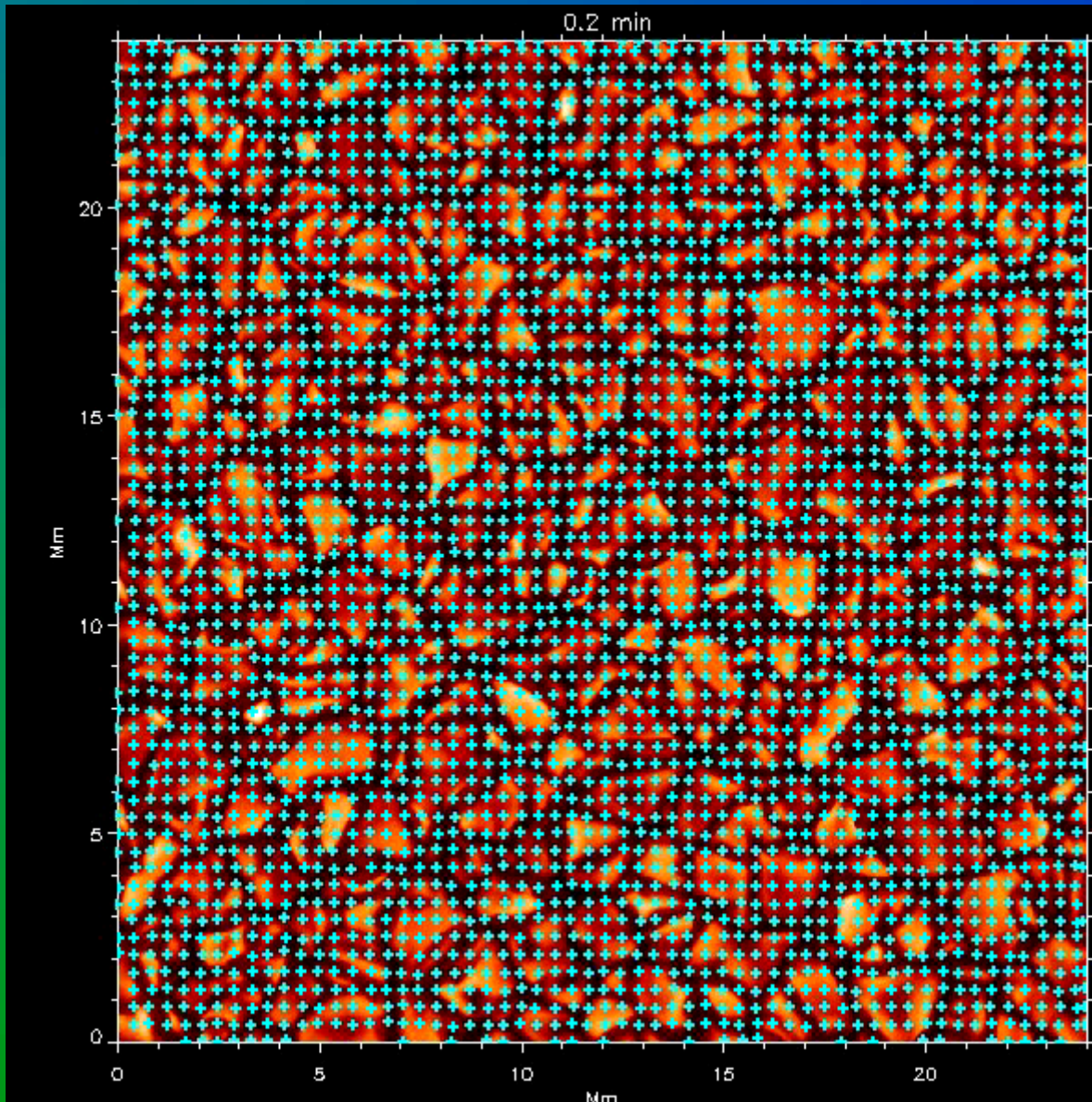
Upper photosphere



“Mesogranulation” ?

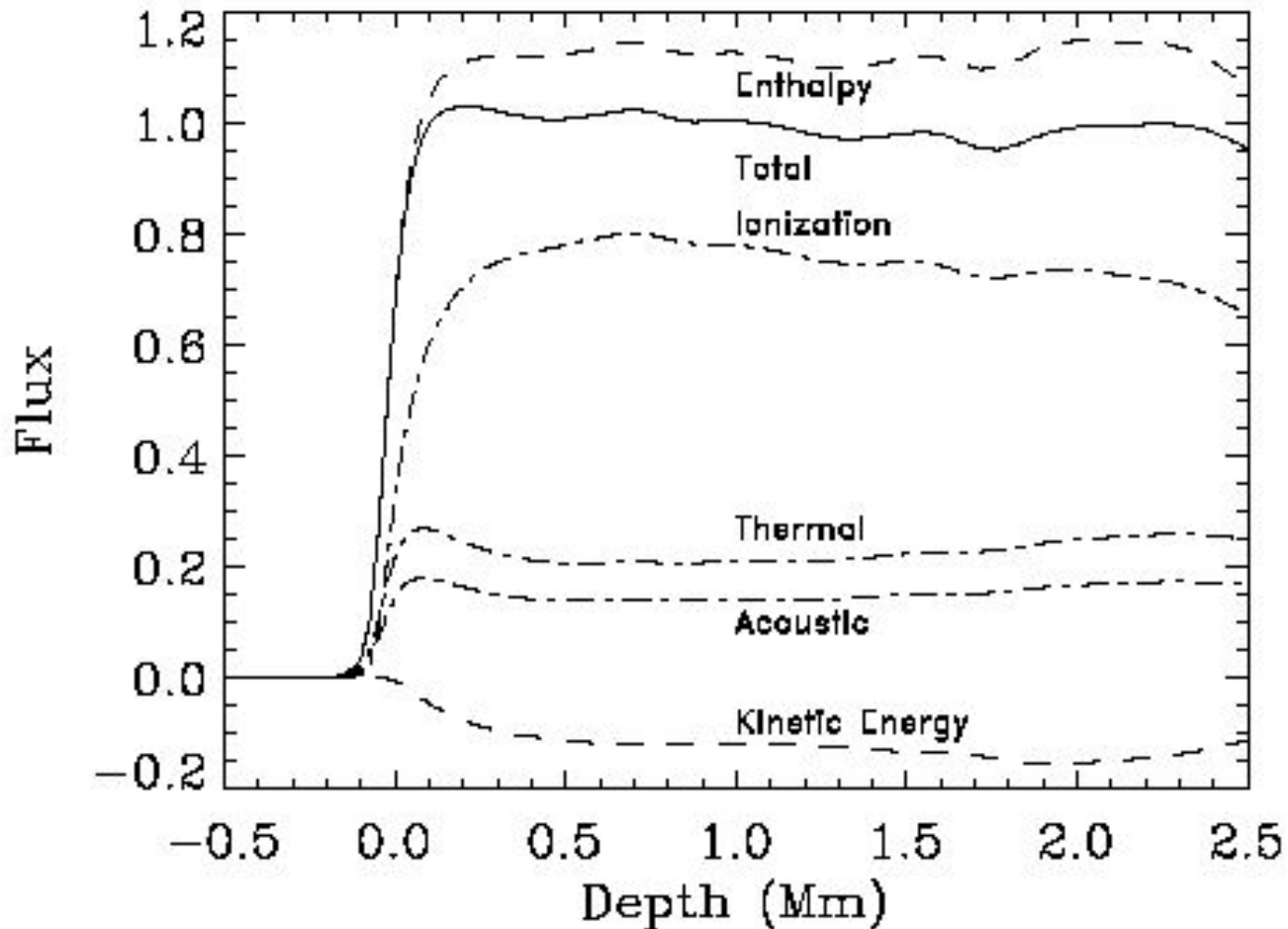


Simulated long-lived convective downflows



Virtual “corks” are carried by the horizontal flow. They accumulate in downflow regions.

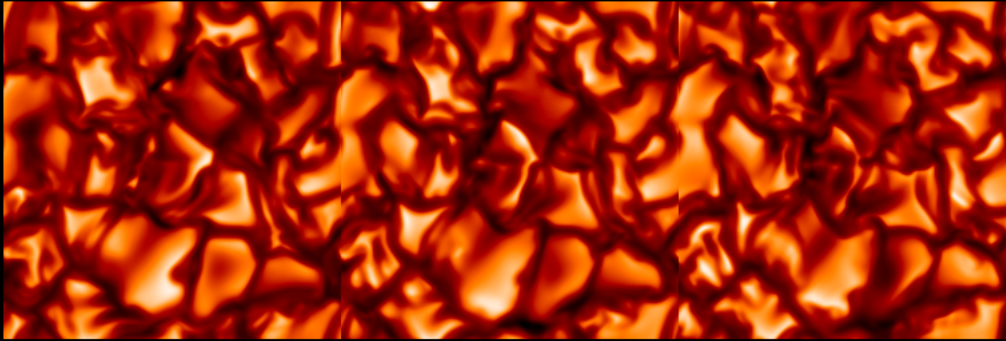
Averaged energy fluxes in a simulation of solar convection



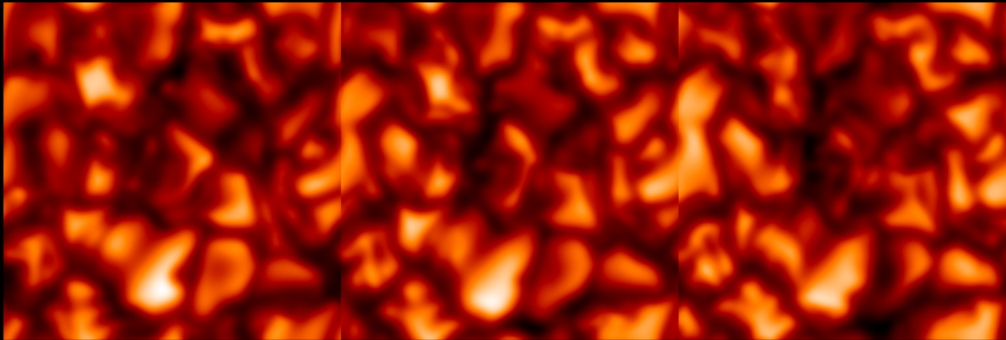
Stein & Nordlund, 2000

Simulation and observation

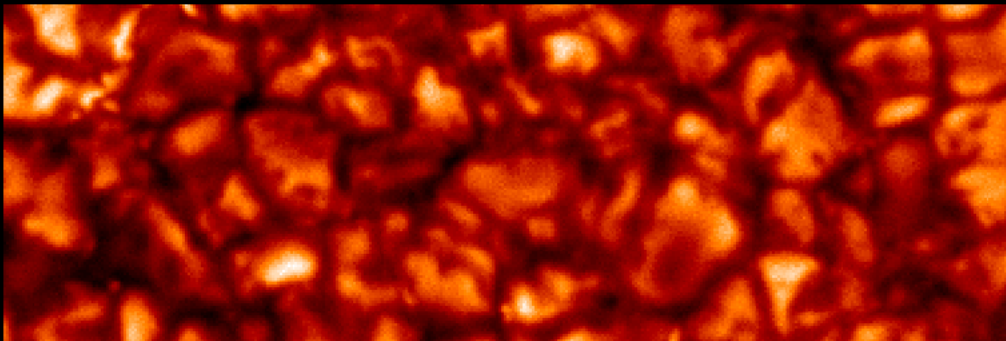
Simulation



Simulation+MTF



Observed

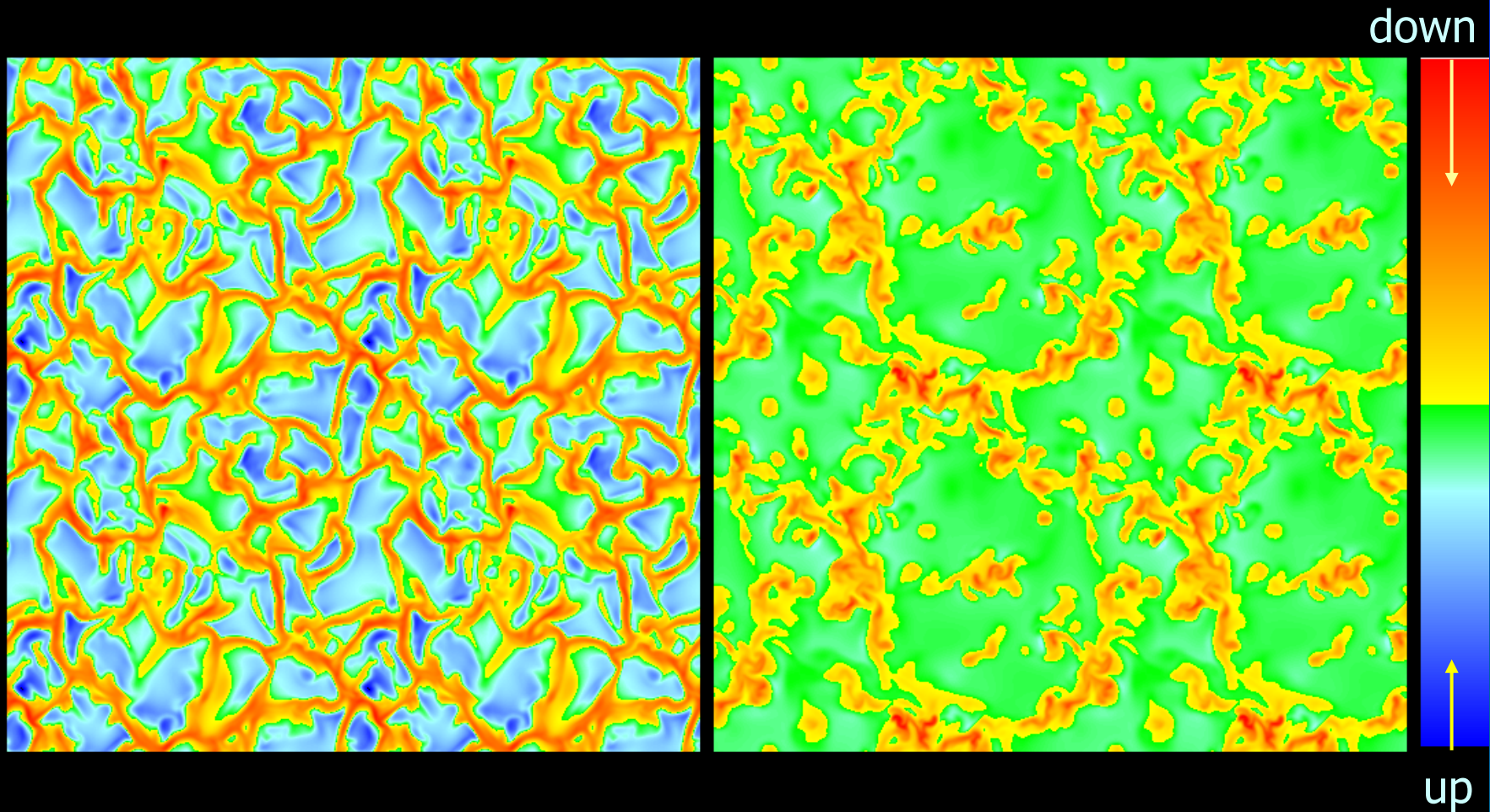


Simulation
(original)

Simulation
(smoothed)

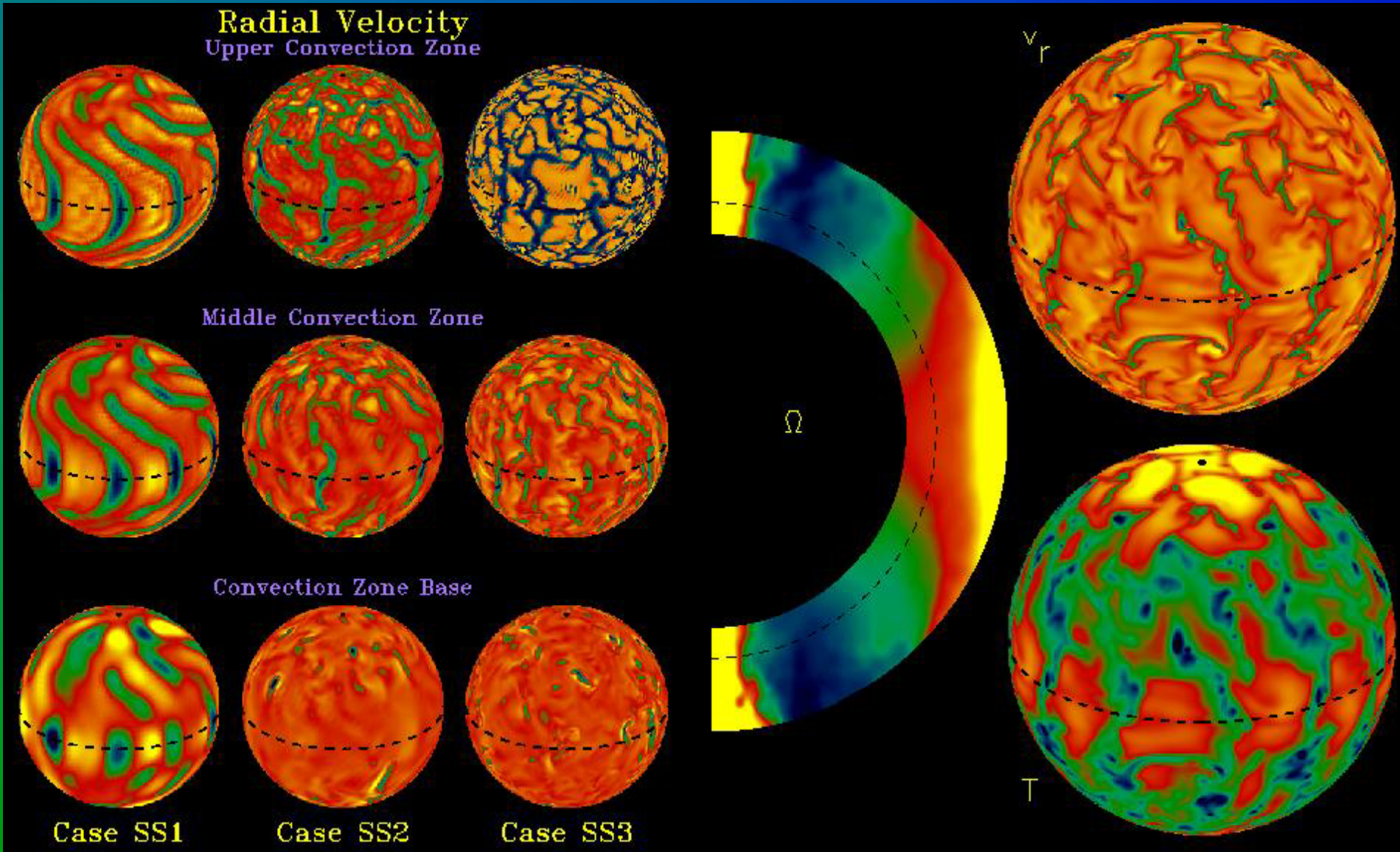
Observation

Change of downflow topology

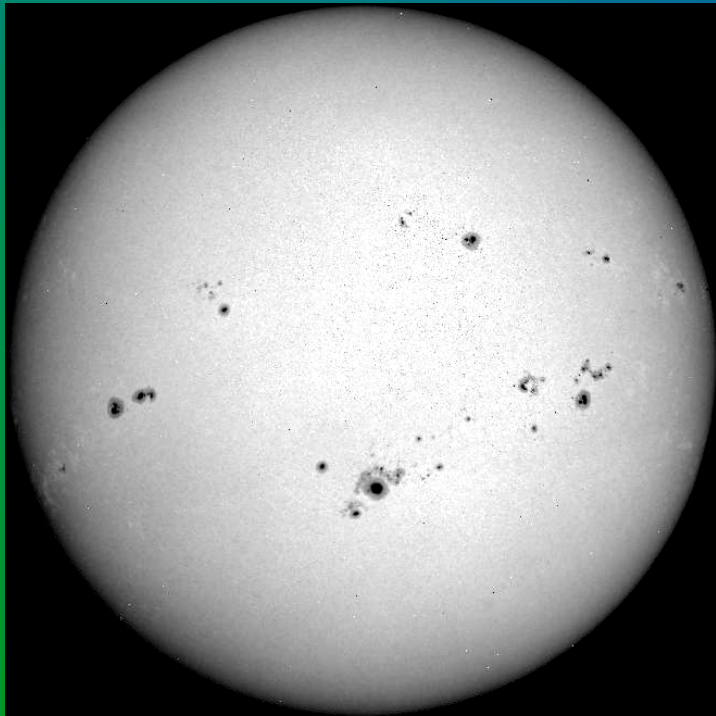


Simulated convection in a solar-like spherical shell

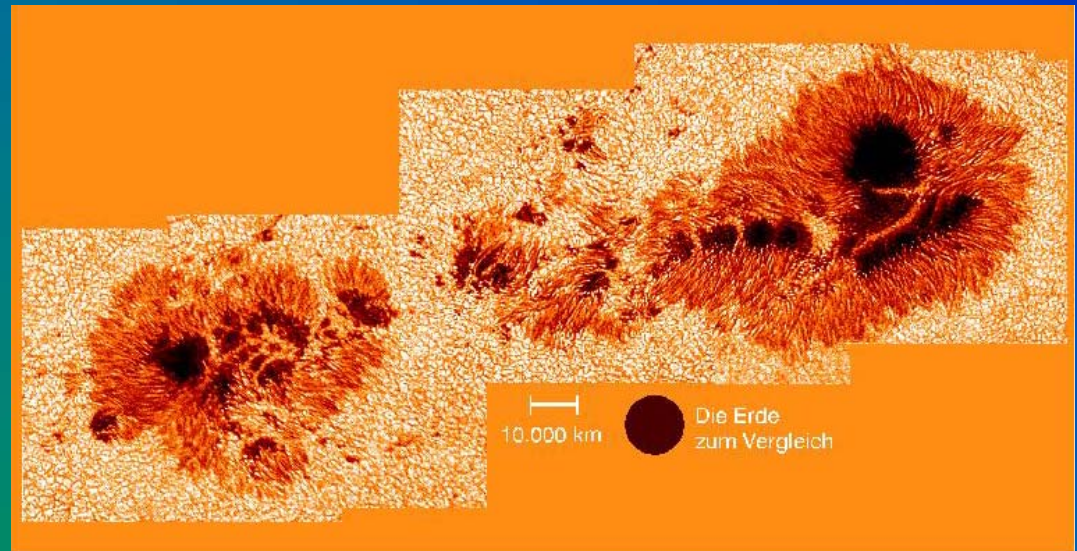
Miesch 1998



Magnetic fields on the Sun

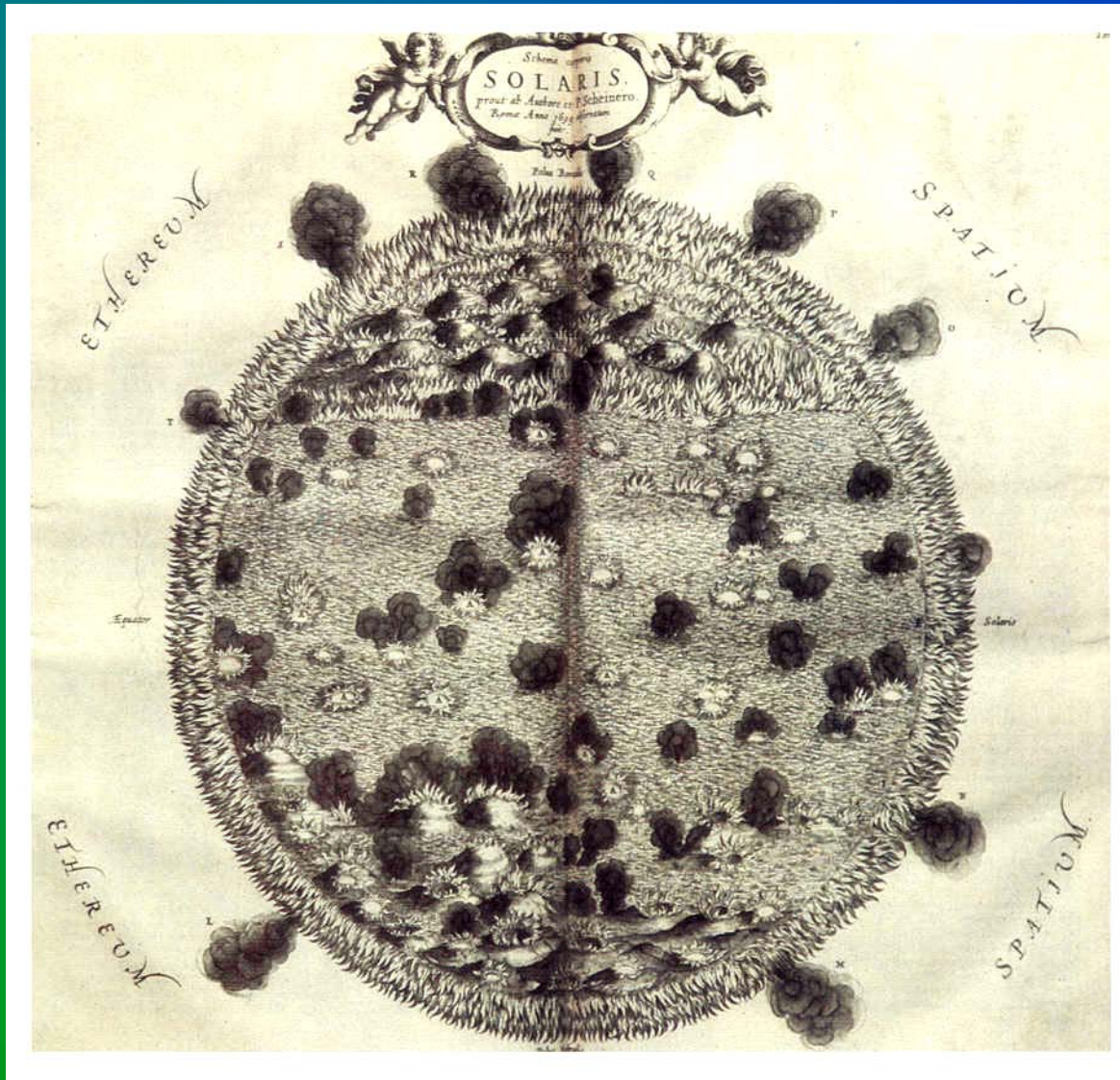


Sunspots



A large sunspot group

What is the nature of sunspots?

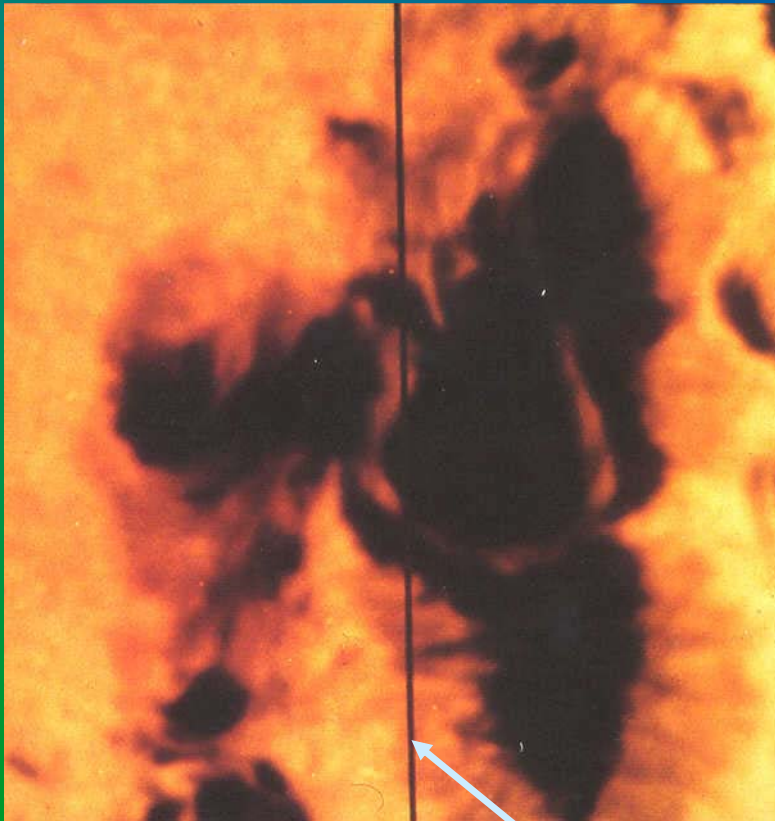
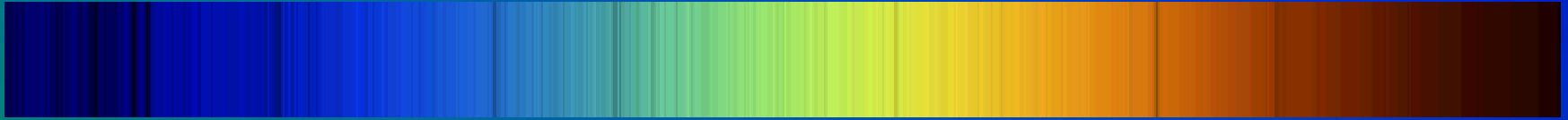


Smoke clouds ?

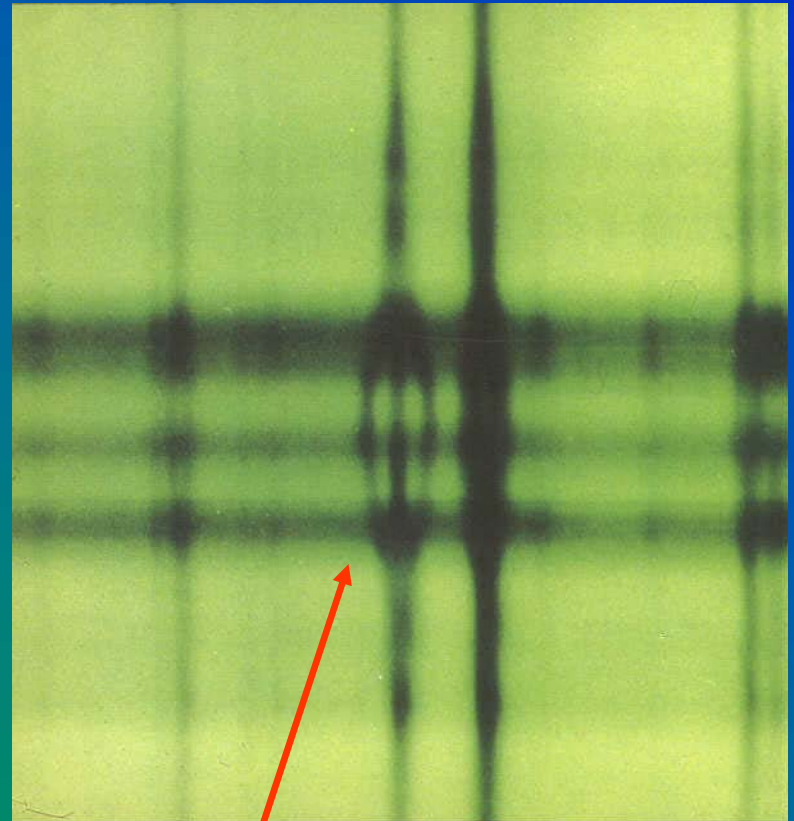
Holes ?

Tornadoes ?

The magnetic nature of sunspots

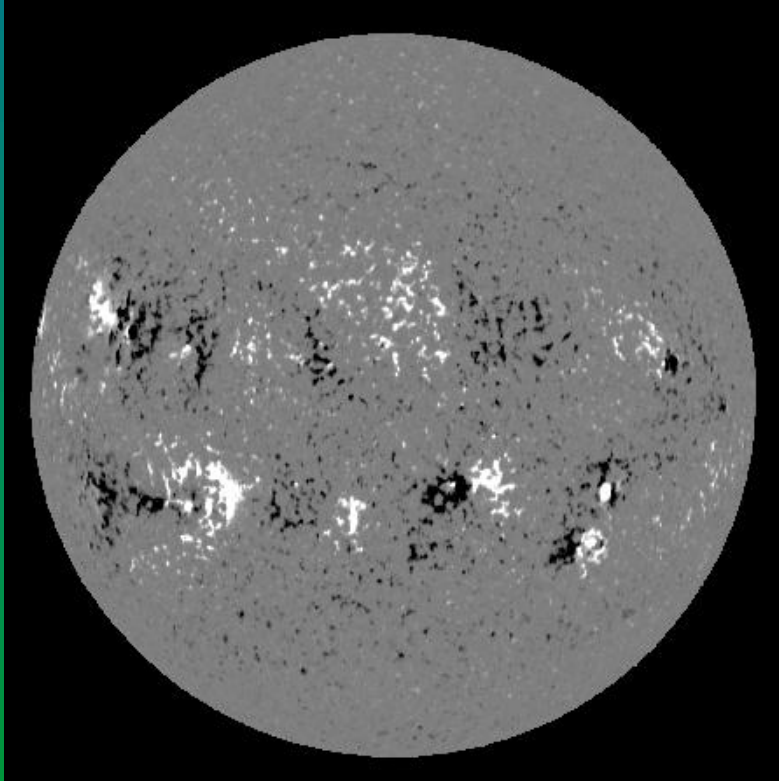


Sunspot with spectrograph slit

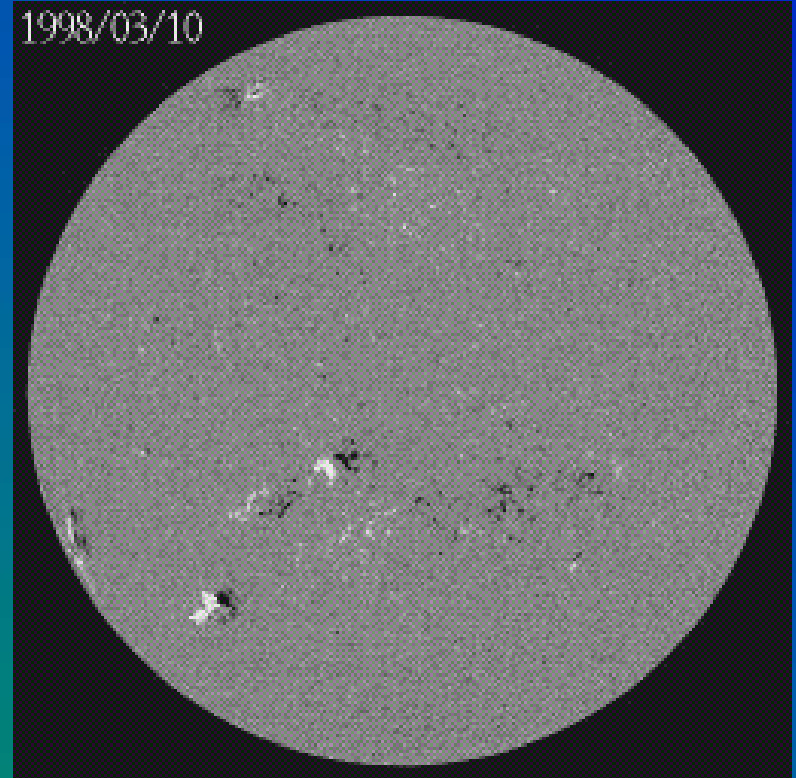


Magnetically split spectral line

Magnetic variability

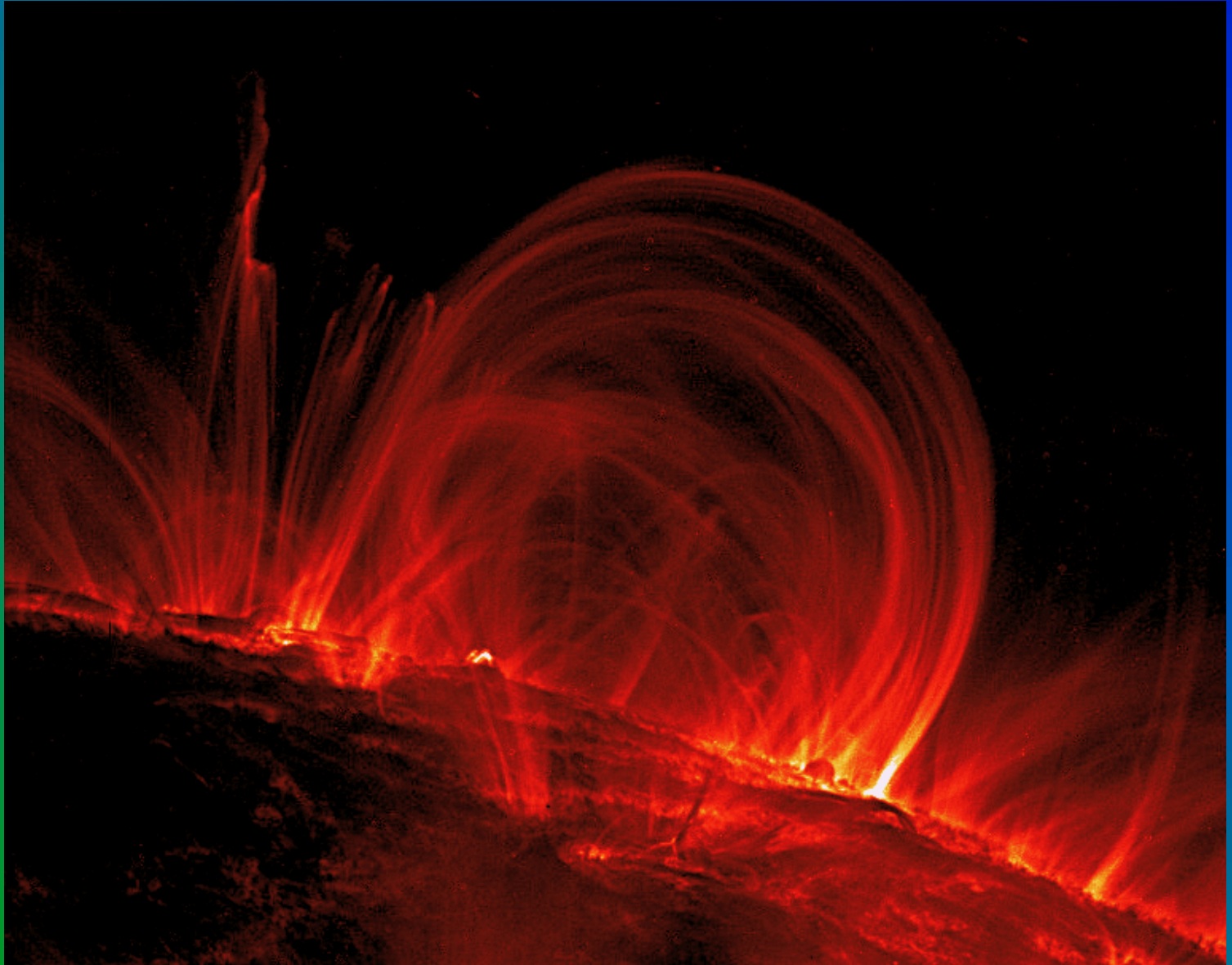


Full-disk magnetogram

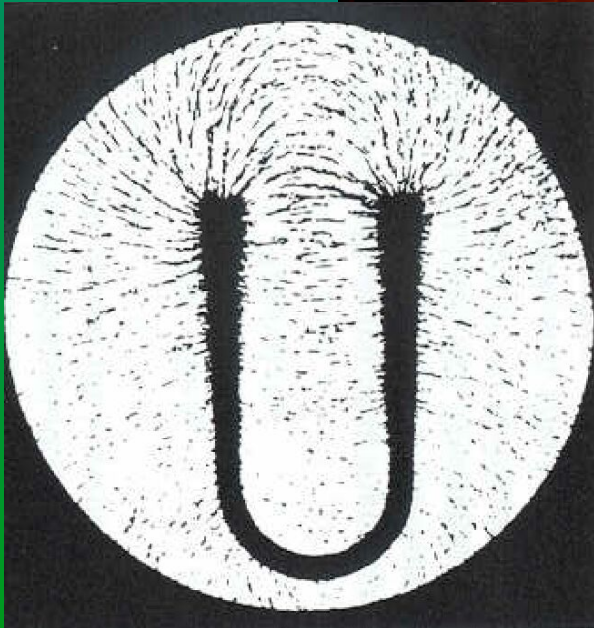
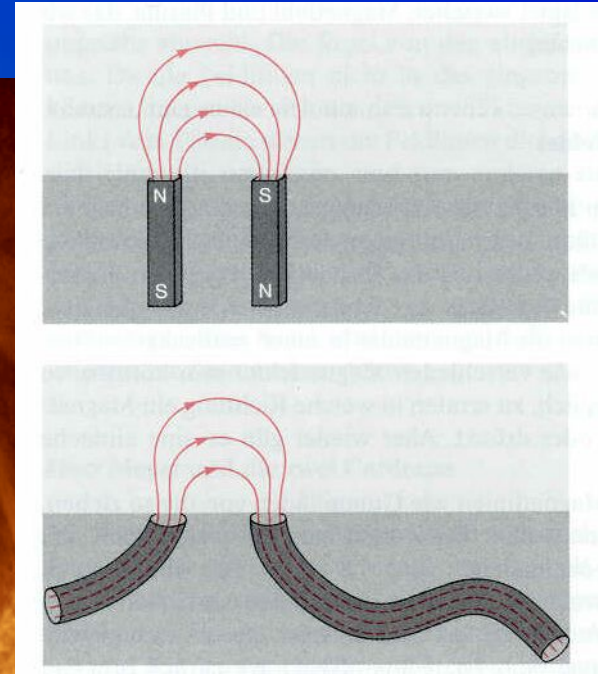
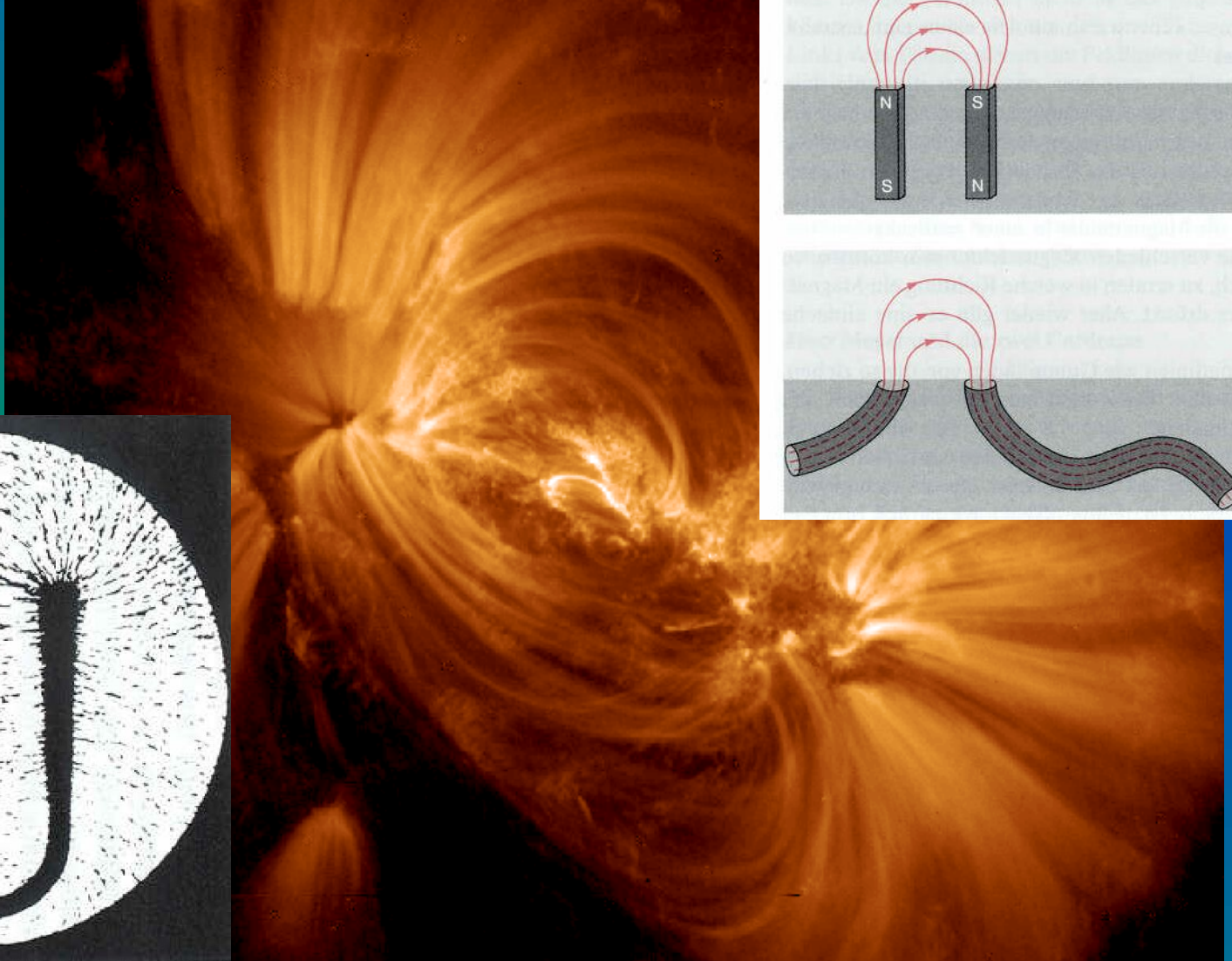


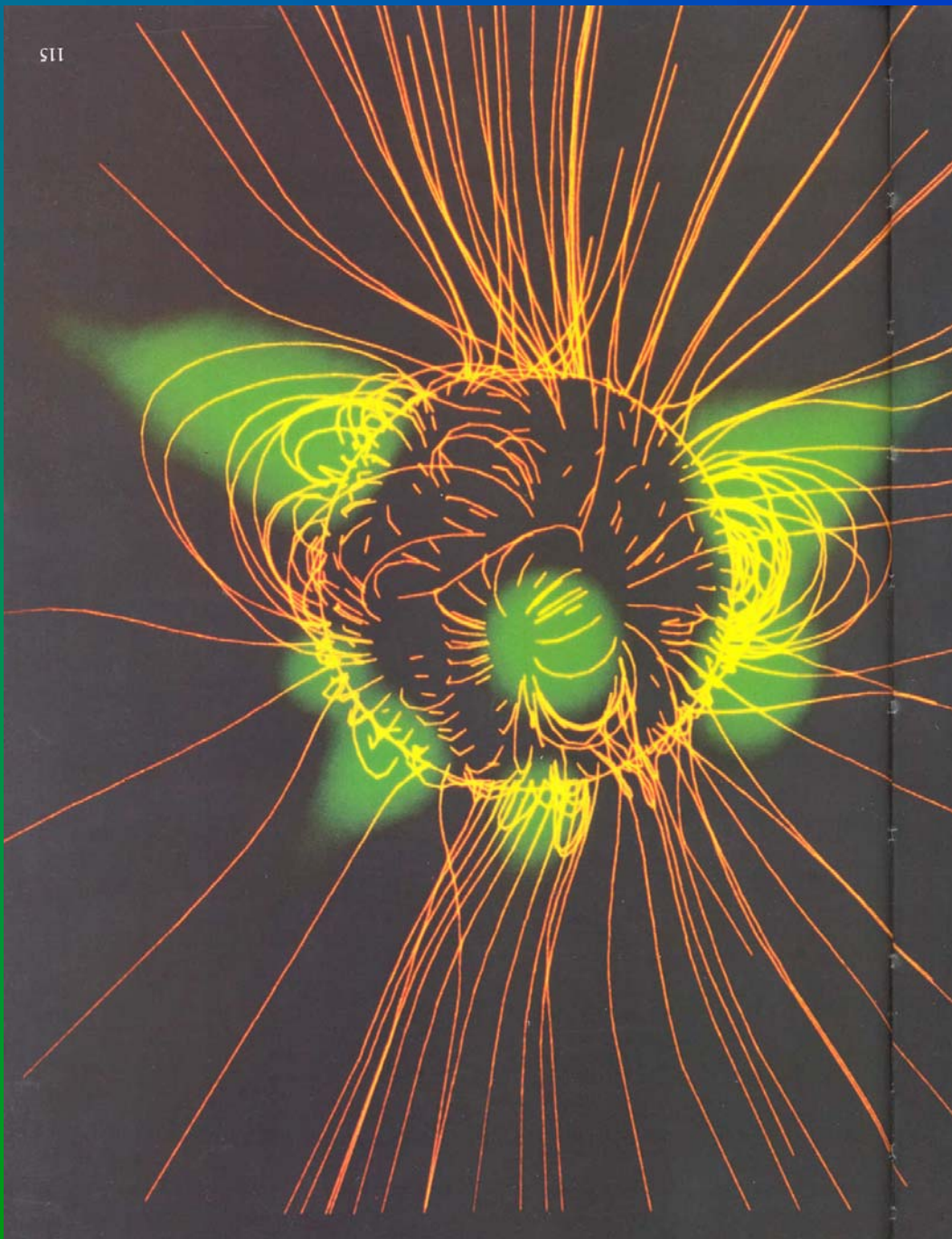
Magnetic patterns on the rotating Sun

Hot plasma draws magnetic field lines...



Hot plasma draws magnetic field lines...





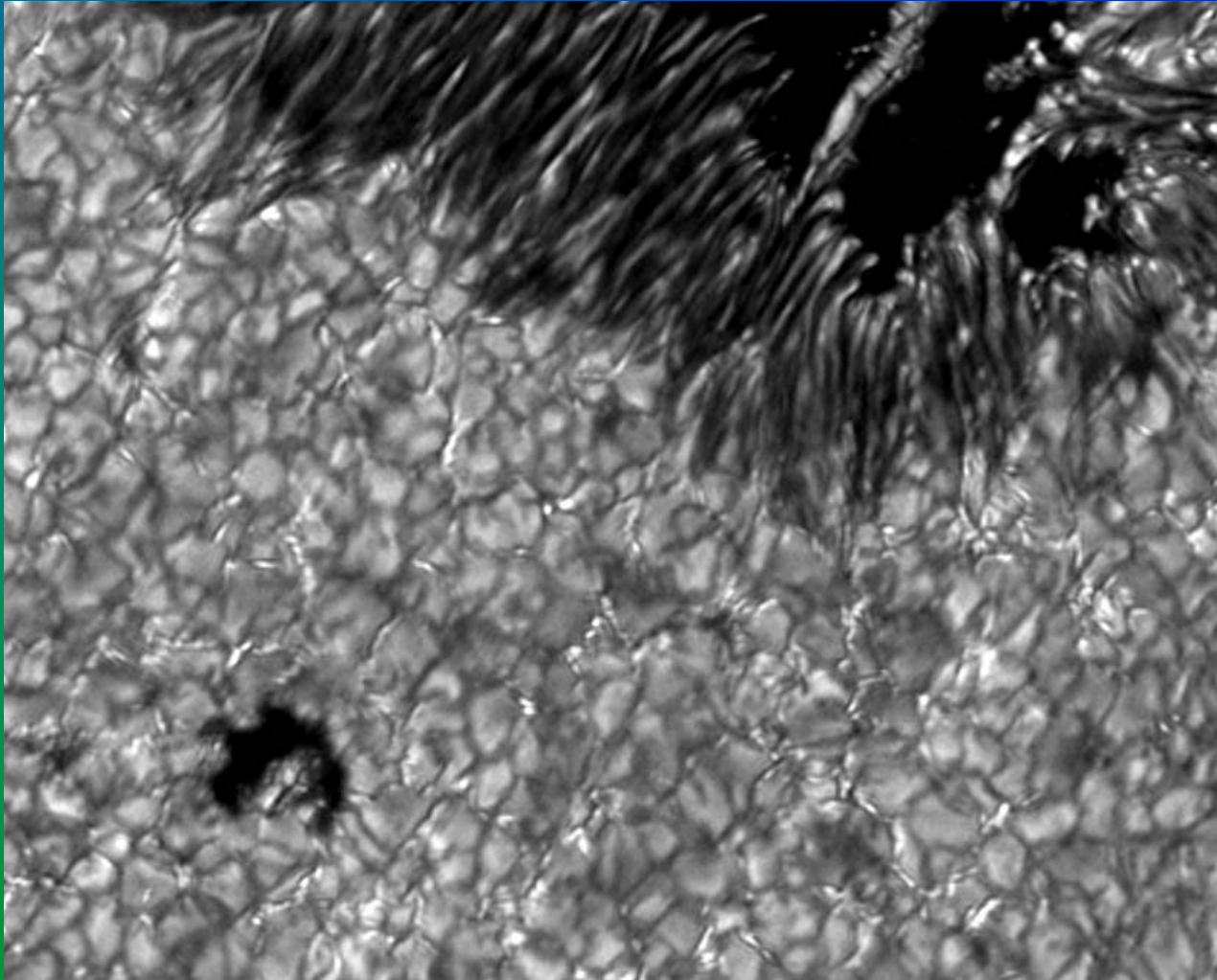
The solar magnetic field...

... continues into interplanetary space.

Its variability in the course of the 11-year cycle and its long-term modulation...

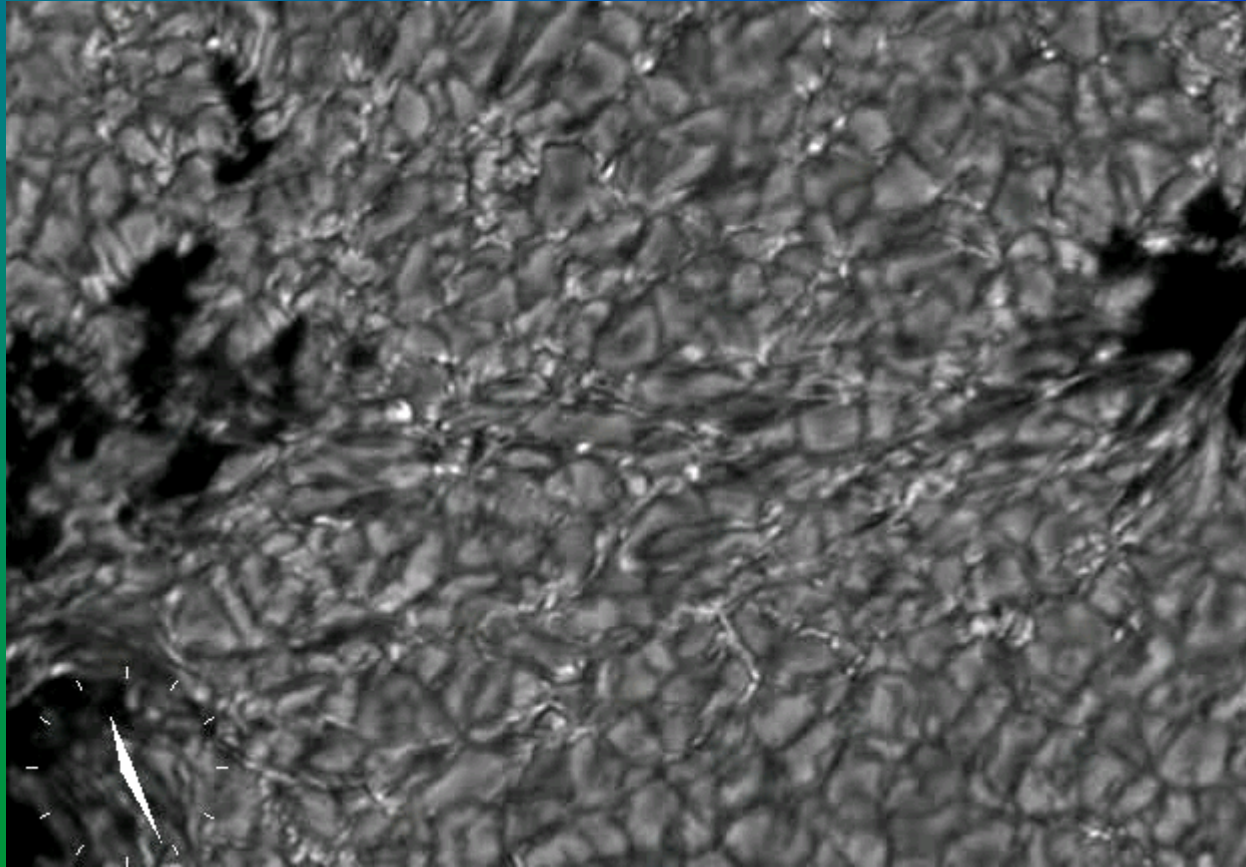
- affects cosmic rays,
- perturbs the terrestrial magnetic field.

G-band observations



KIS/VTT, Obs. del Teide, Tenerife

G-band observations



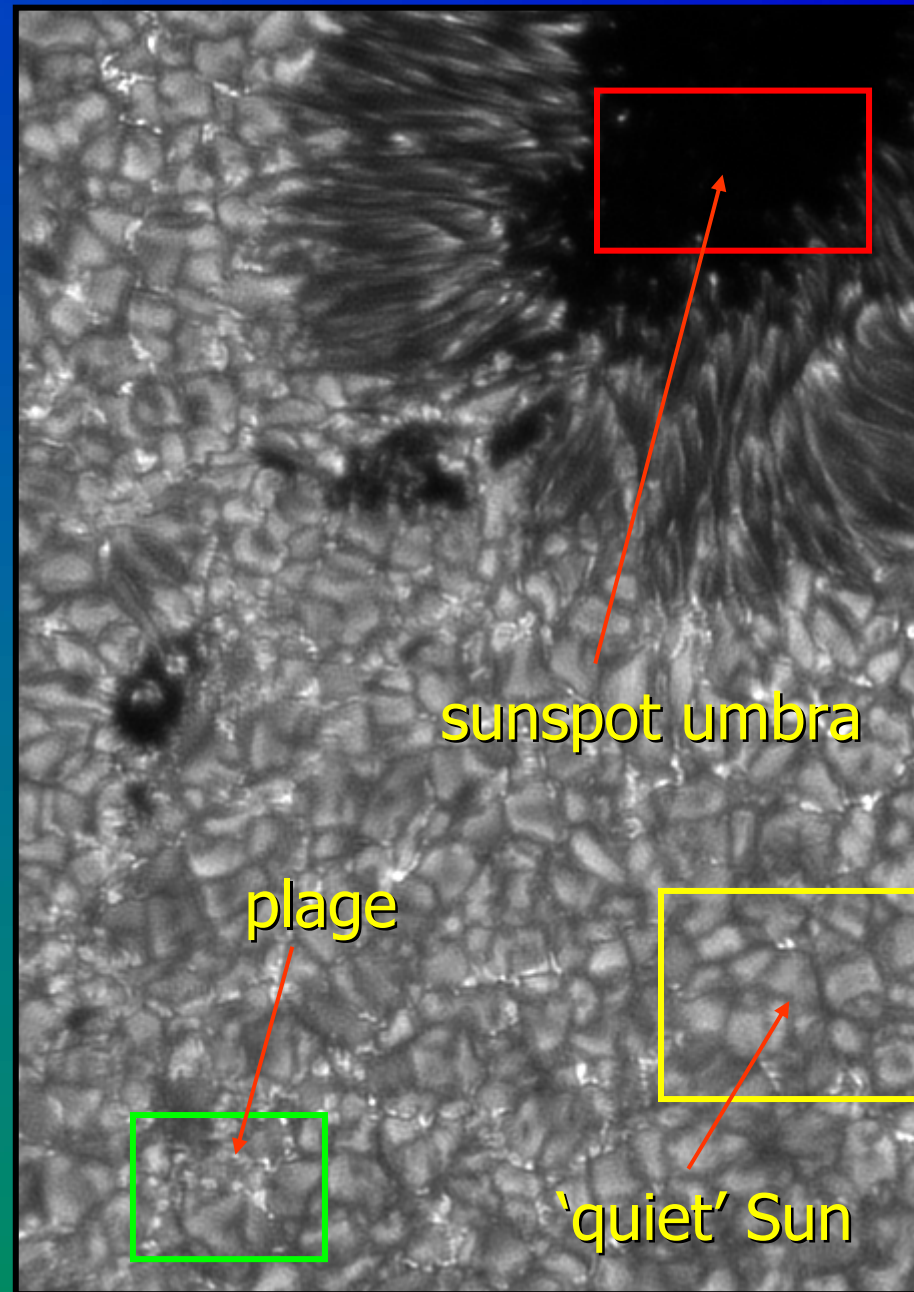
Dutch Open Telescope, Obs. del Roque de los Muchachos,
courtesy P. Sütterlin

What is magneto-convection?

- Interaction between convective flows and magnetic field in an electrically well-conducting fluid
- High Reynolds numbers: nonlinear dynamics, structure and pattern formation
- Interference with convective energy transport

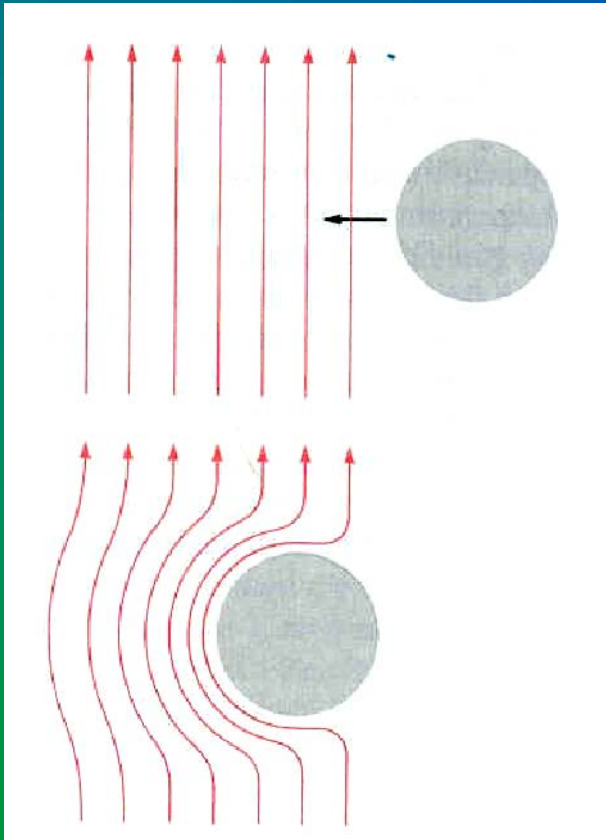
Regimes of solar magneto-convection

- $\langle B \rangle$ increases:
quiet Sun \rightarrow plage \rightarrow
 \rightarrow umbra
- horizontal scale of
convection decreases
- convective energy
transport decreases

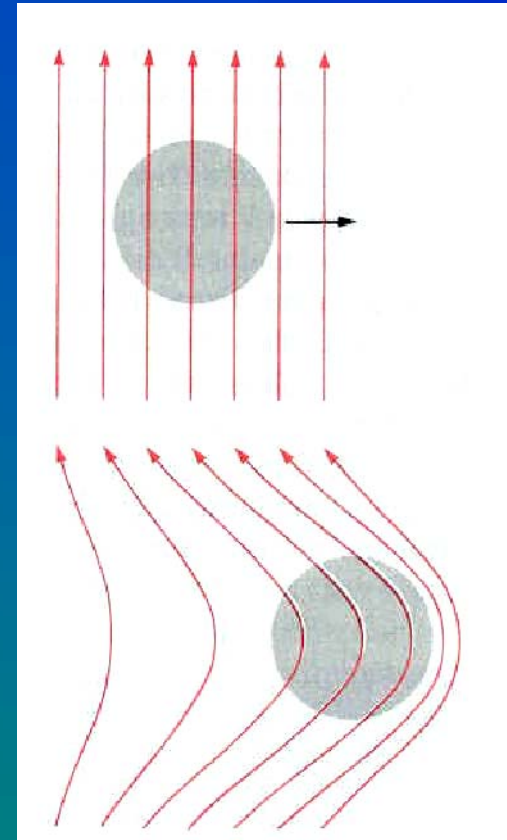


T. Berger, SVST 12 May 1998, Obs. del Roque de los Muchachos
Adapted from a figure by Thierry Emonet, Univ. Chicago

Good electrical conductors : "frozen field"

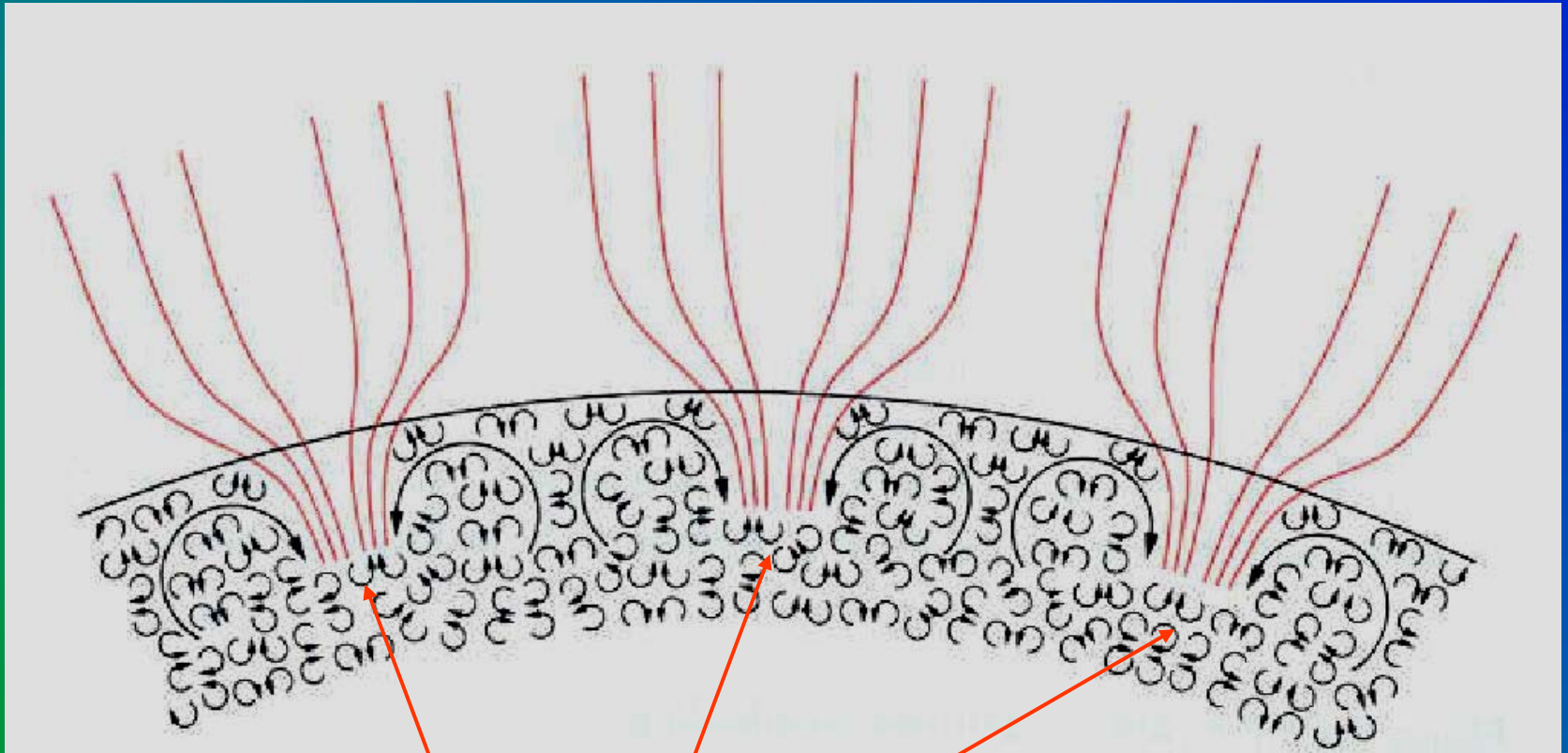


Initially field-free volumes remain field-free



Magnetic flux through a given volume remains constant

“Frozen field” in the Sun

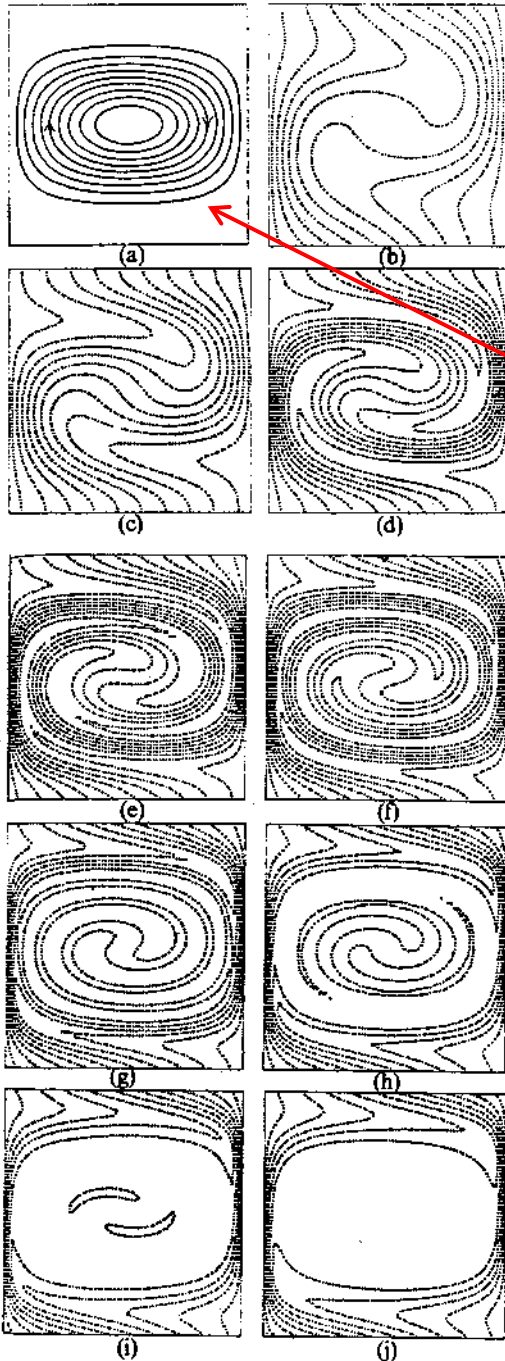


Magnetic flux is transported to the downflow regions of the convective flow patterns

“magnetic network”

Simulation of flux expulsion

(Weiss, 1966)



b: final state for $Re_m = 40$

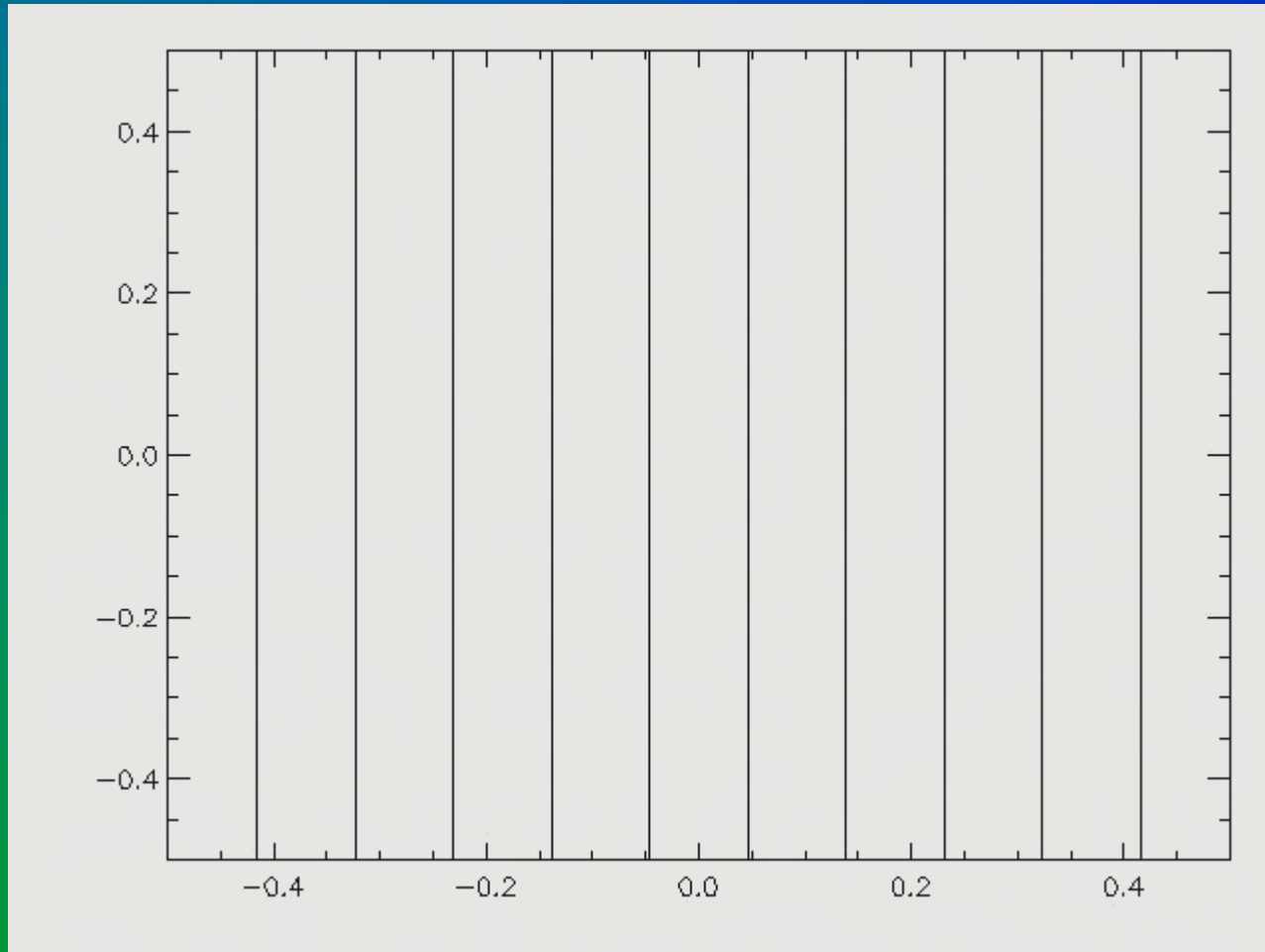
a: streamlines of the fixed velocity field

c-j: time evolution for $Re_m = 1000$

- evolution of an initially vertical magnetic field under the influence of a fixed flow field
- kinematic, 2D
- the magnetic flux is expelled from the area of closed streamlines and concentrated in narrow sheets

Flux expulsion and intermittency

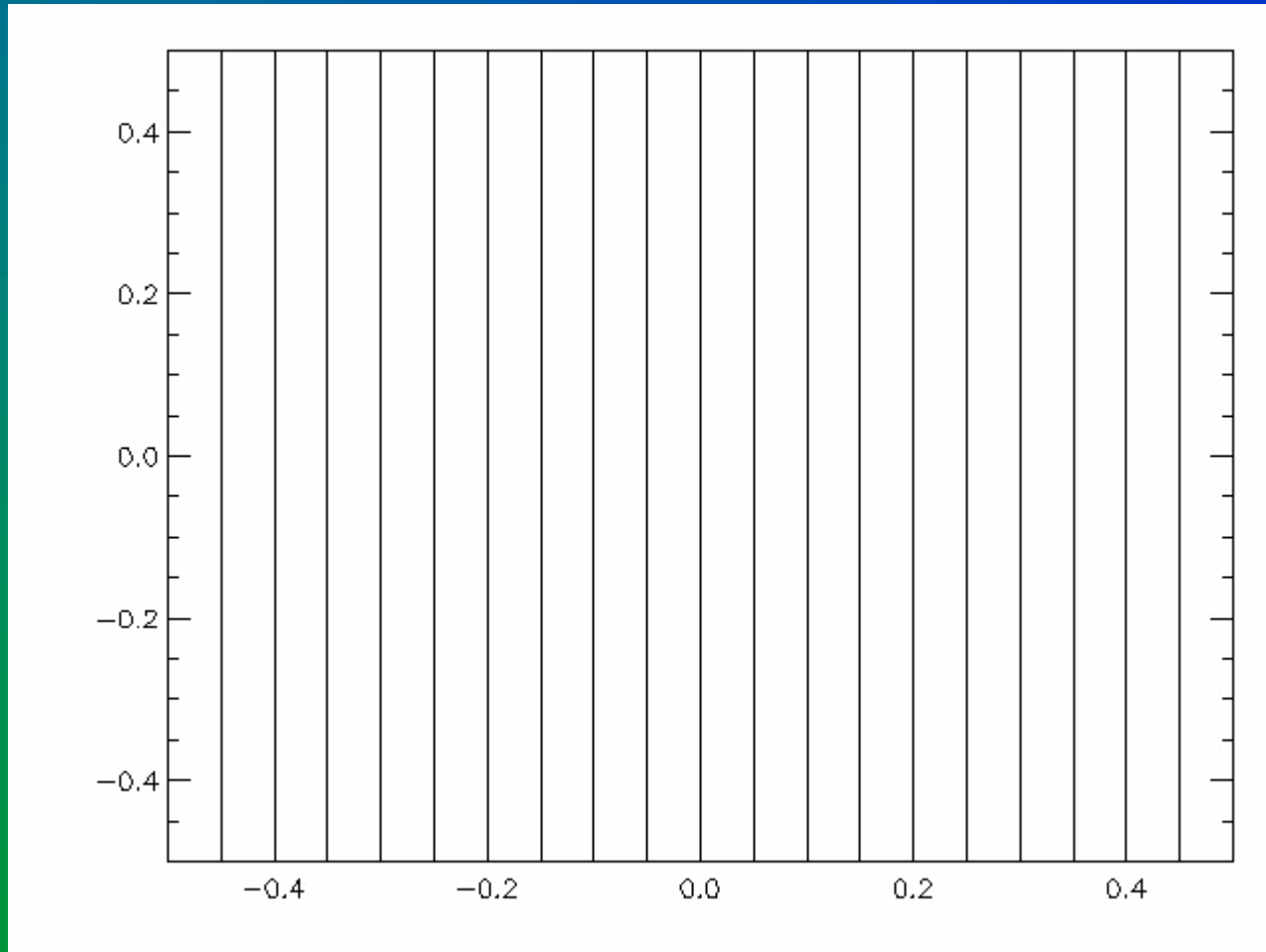
- N.O. Weiss (1964): *first simulations*



(Hupfer, KIS Freiburg, 2001)

Flux expulsion and intermittency

- N.O. Weiss (1964): *first simulations*



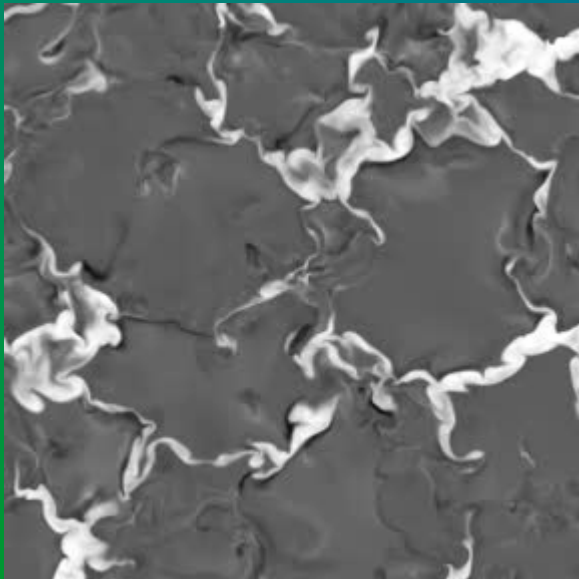
(Hupfer, KIS Freiburg, 2001)

$B_0 = 200 \text{ G}$ (plage): time evolution

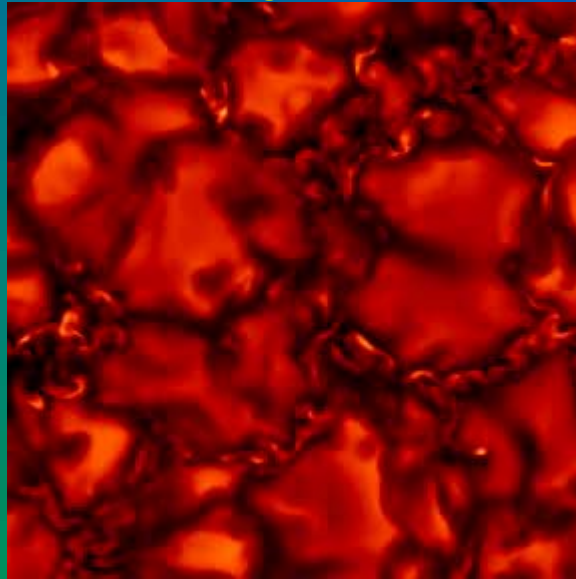
horizontal cuts near $\tau=1$
6000 km \times 6000 km \times 1400 km
576 \times 576 \times 100 grid points

up   down

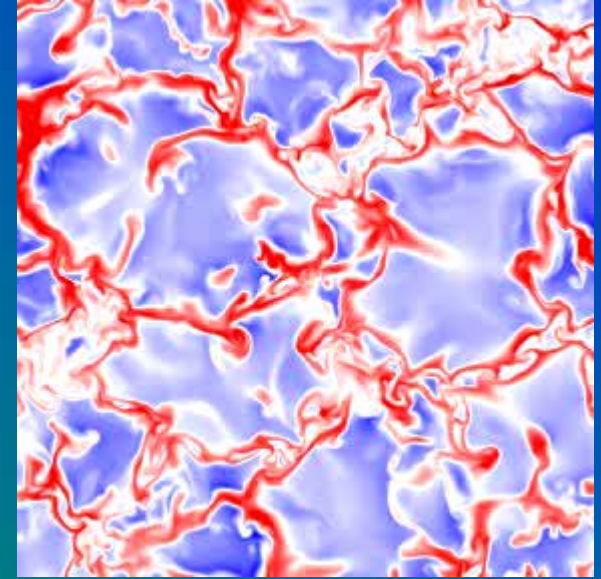
vertical magnetic field



brightness



vertical velocity



-1.3 0 +2.2

[kG]



0.6 1.0 1.5

$I/\langle I \rangle$

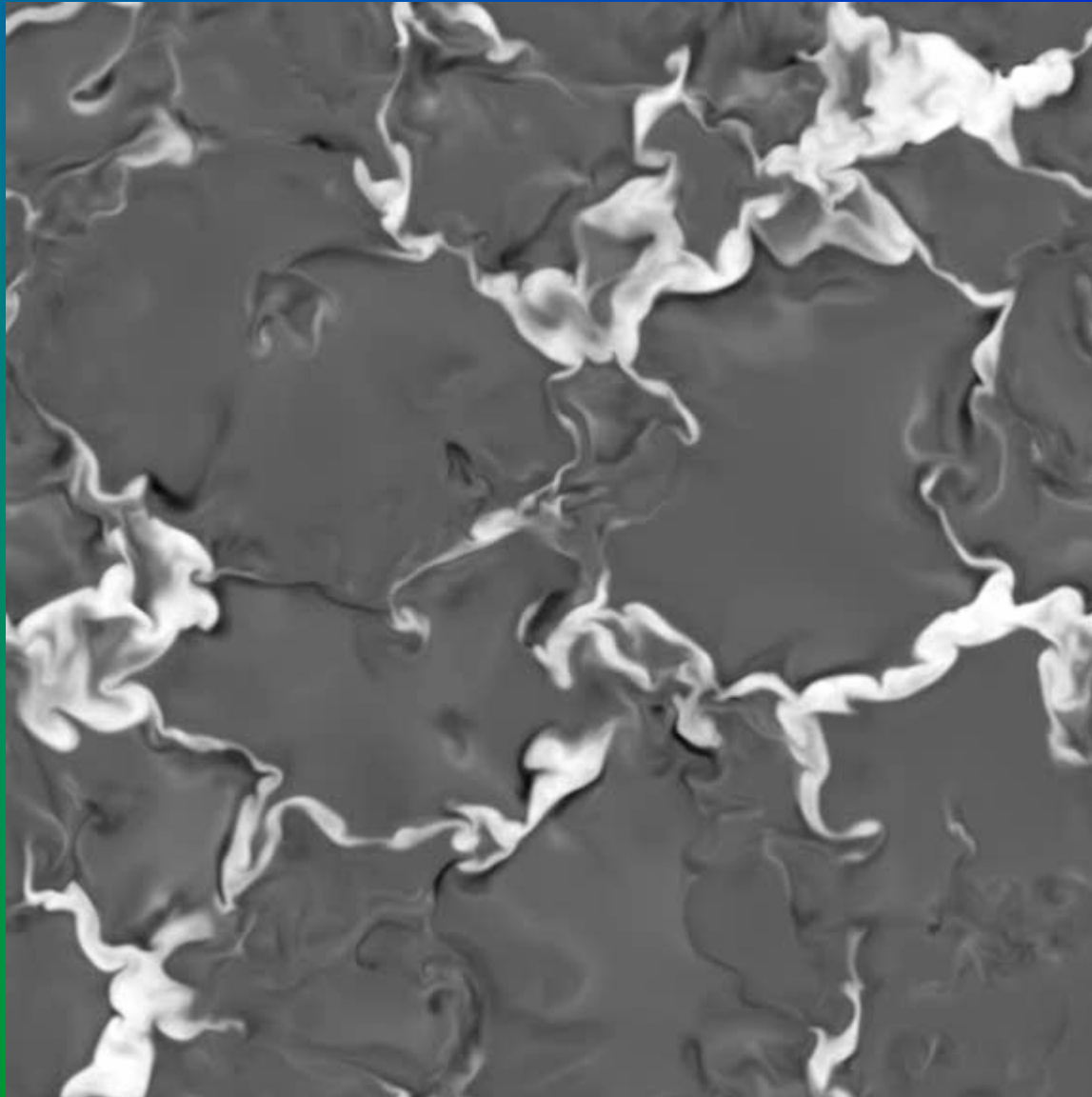


-10 0 +10

[km/s]

$B_0 = 200 \text{ G}$ (plage): time evolution

+2 kG
↑
↓
-2 kG

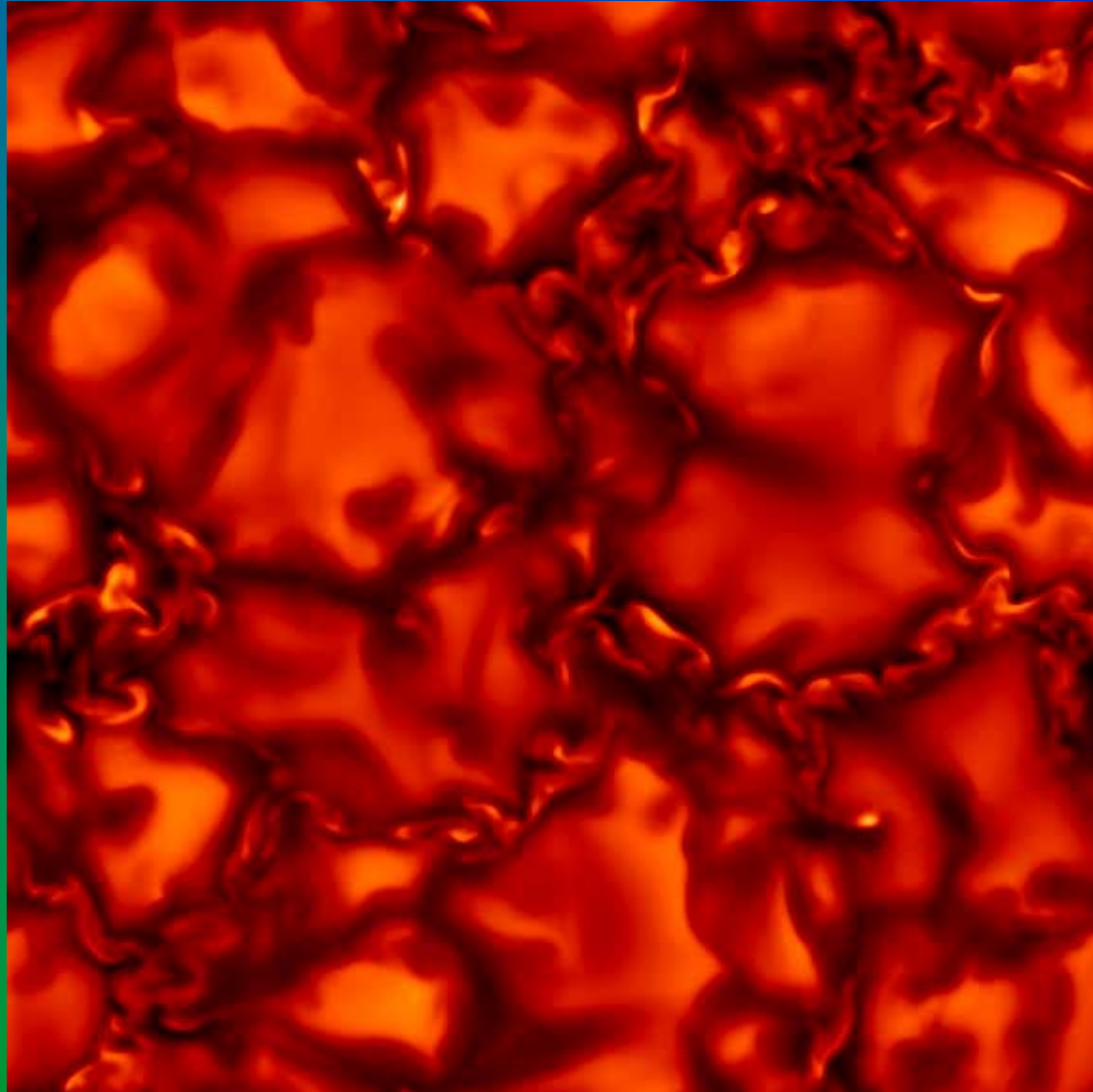


6 Mm

Vertical magnetic
field component

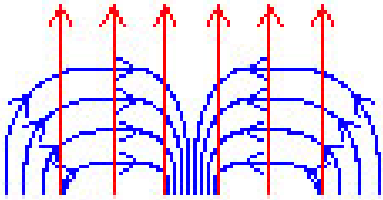
$B_0 = 200$ G (plage): time evolution

Brightness

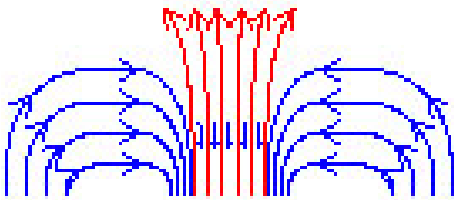


6 Mm

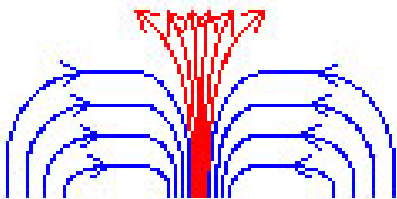
Convective intensification



- Flux advection by horizontal flow (flux expulsion)

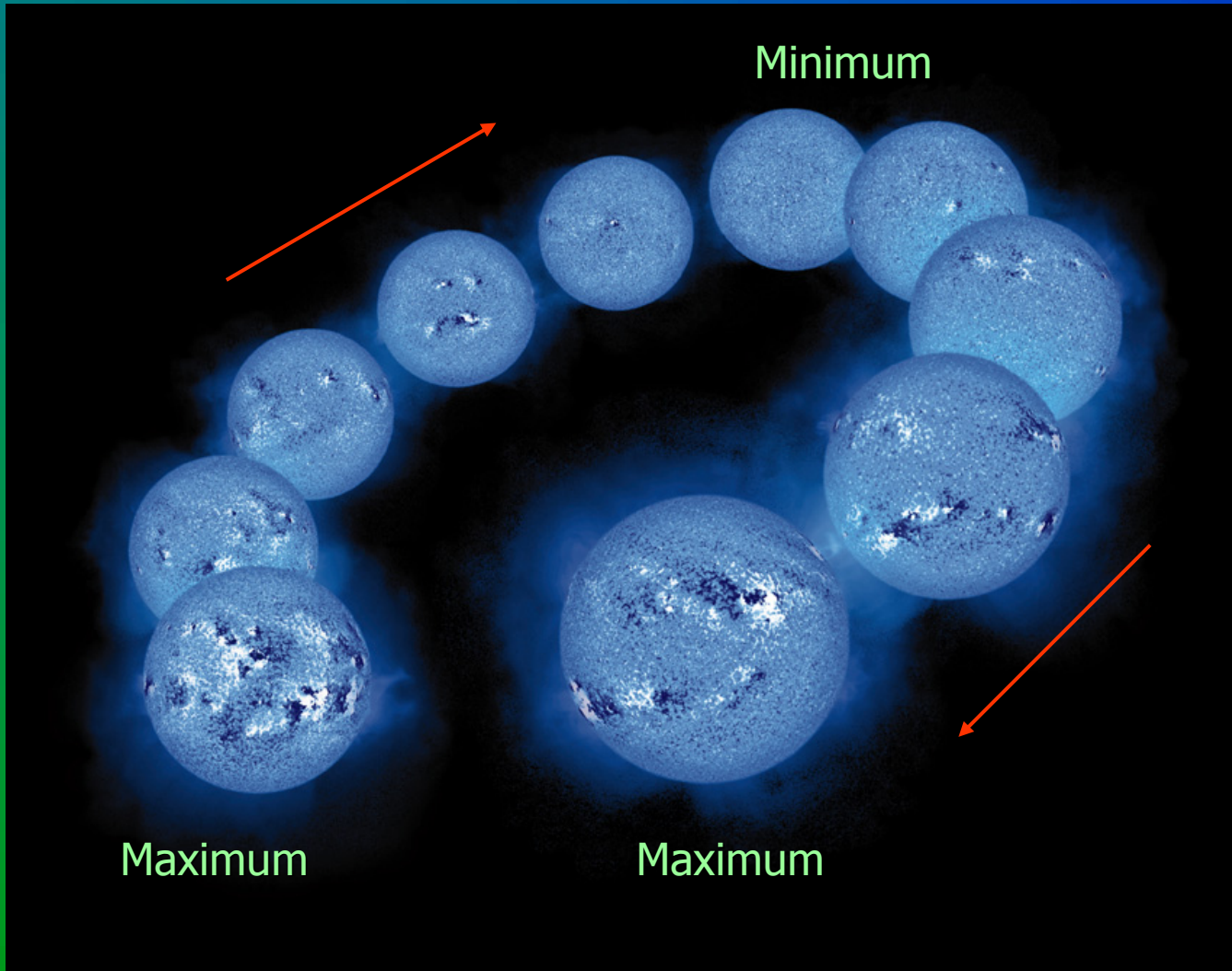


- Suppression of convection, cooling and downflow



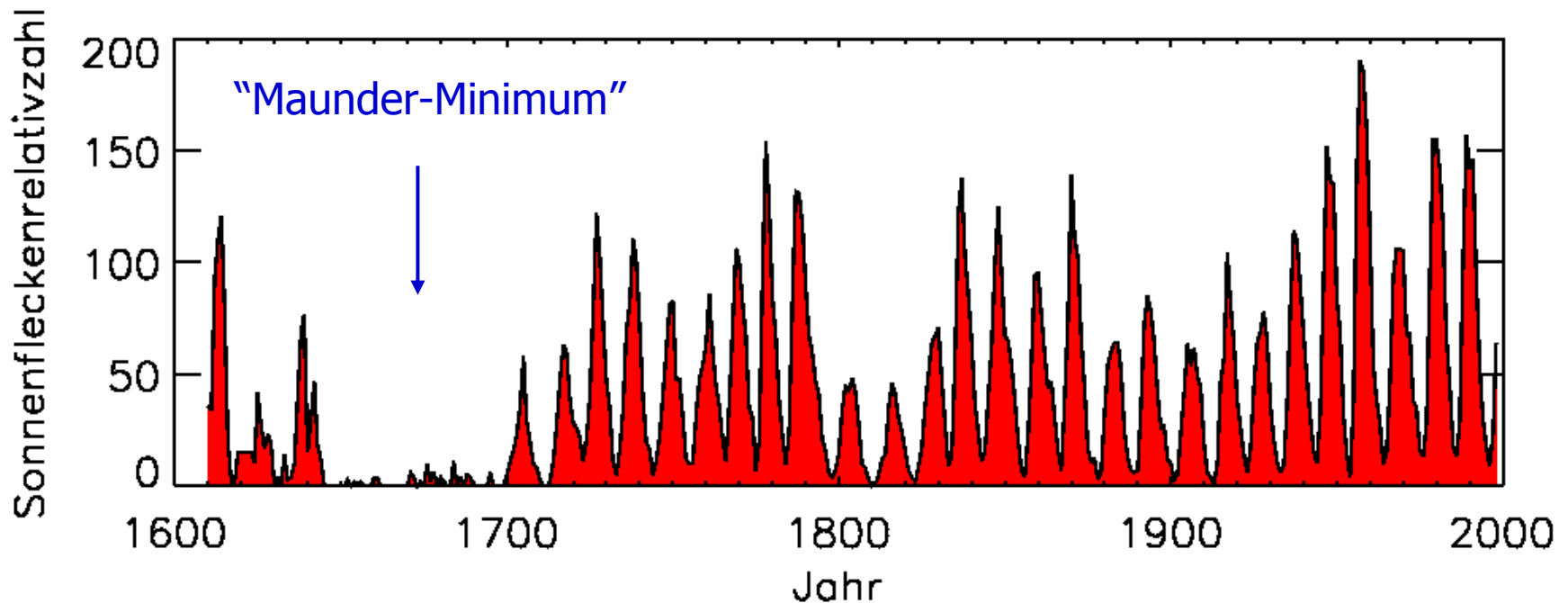
- Evacuation, field intensification

The magnetically variable Sun



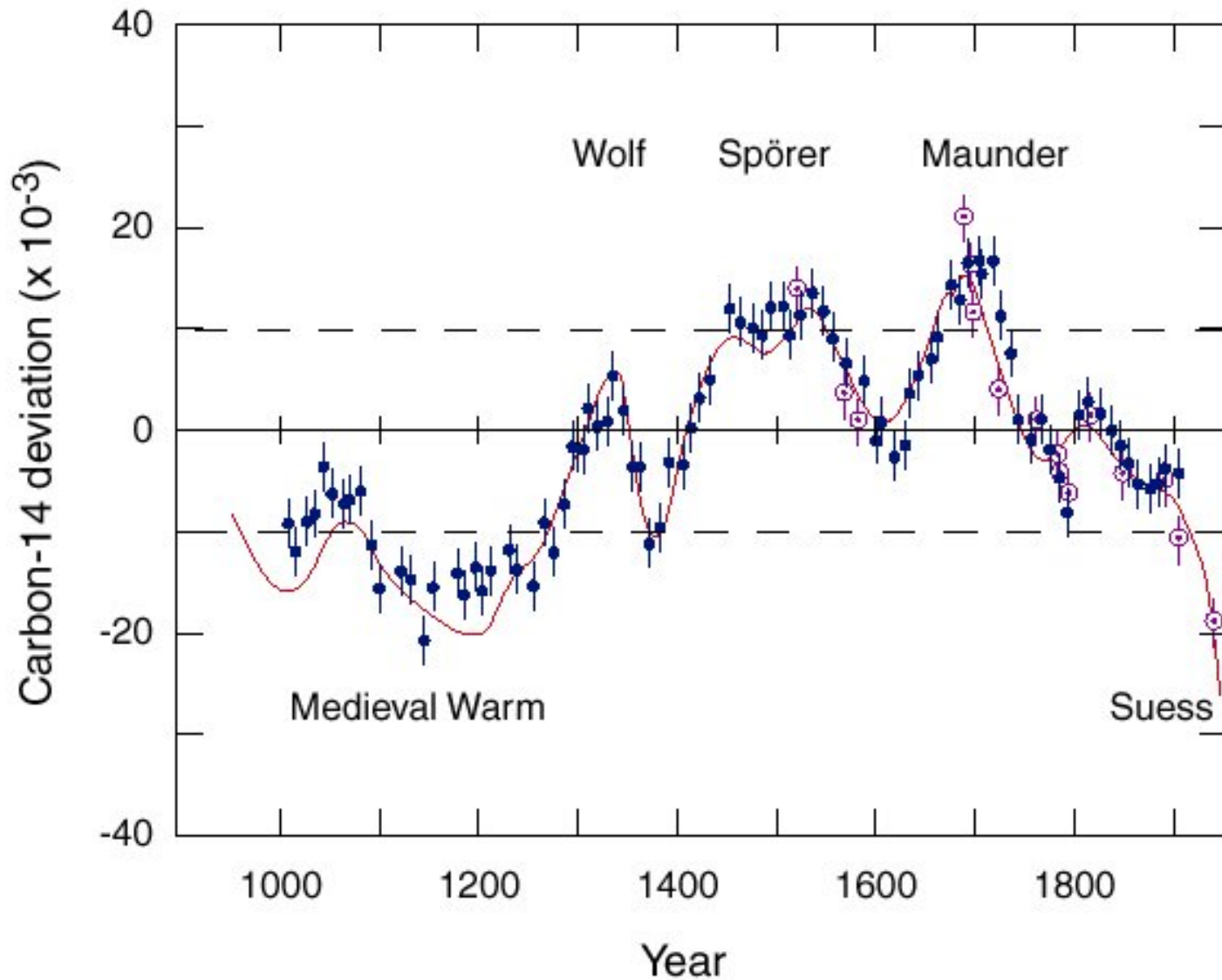
11-year cycle of magnetic activity and surface flux

The 11-year solar cycle



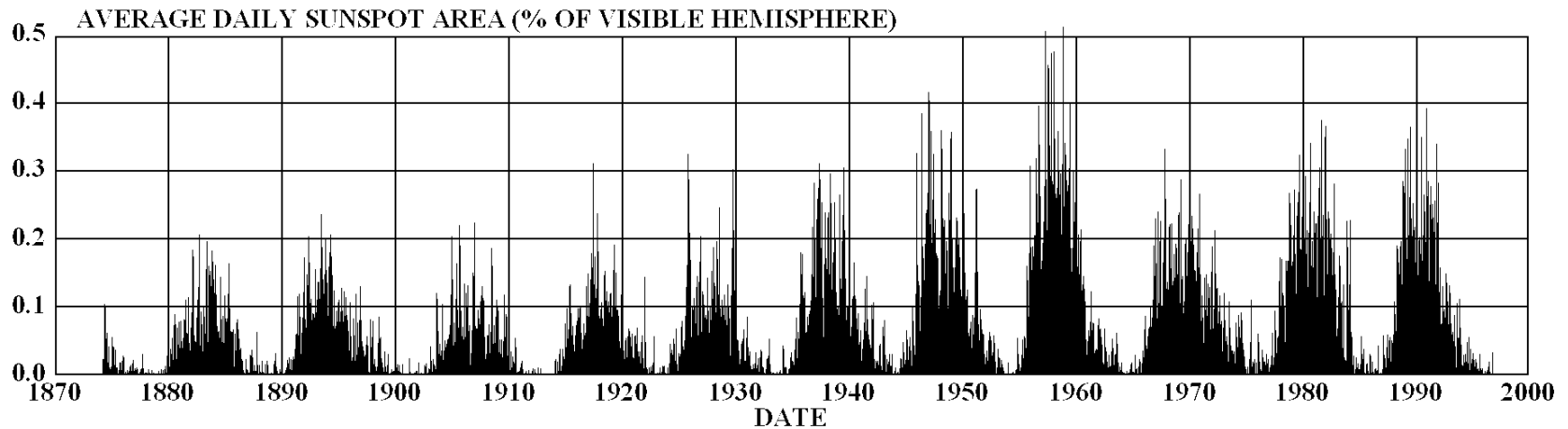
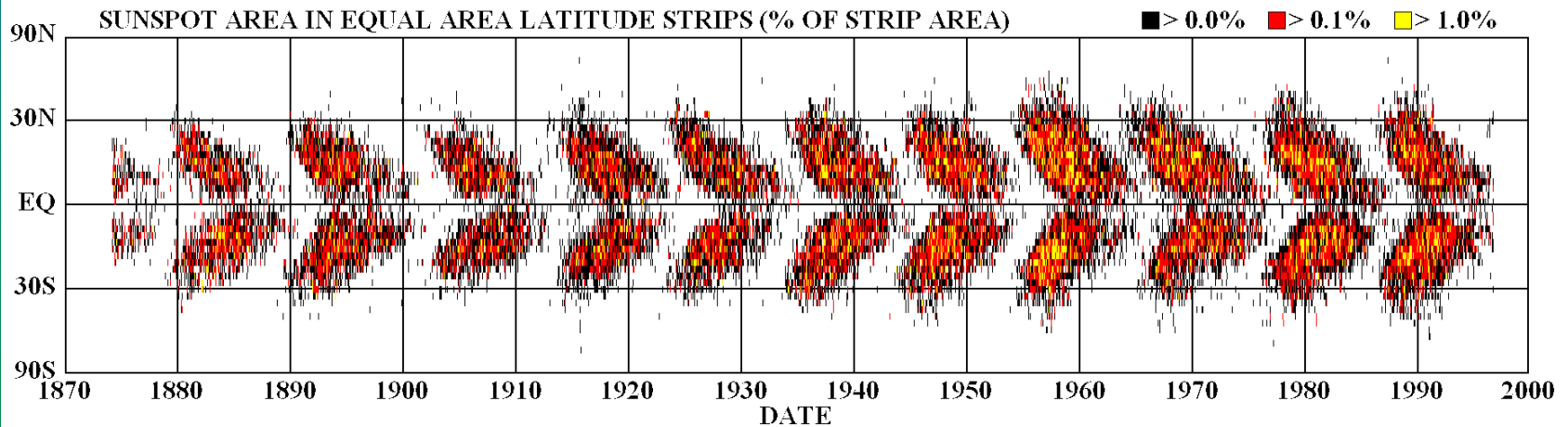
Solar magnetic activity varies with a period of roughly 11 years. Long-term variations are superposed upon this cycle.

^{14}C : Solar activity back to AD 1000



Butterfly diagram

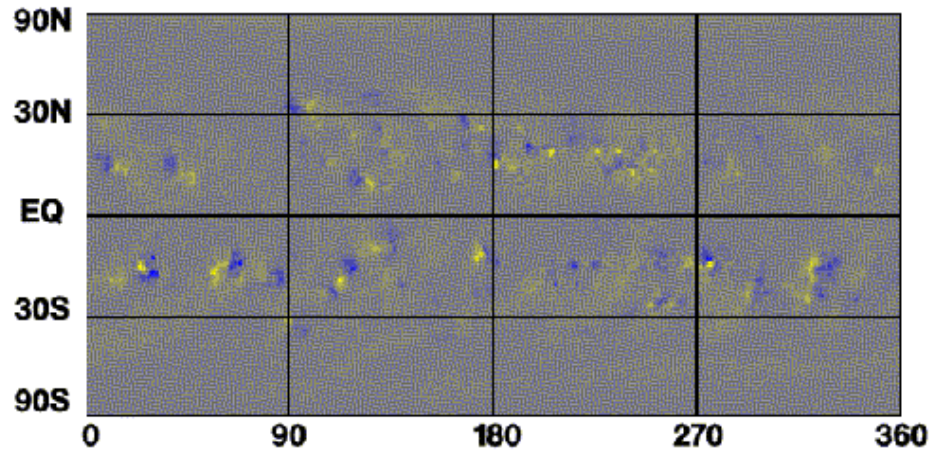
DAILY SUNSPOT AREA AVERAGED OVER INDIVIDUAL SOLAR ROTATIONS



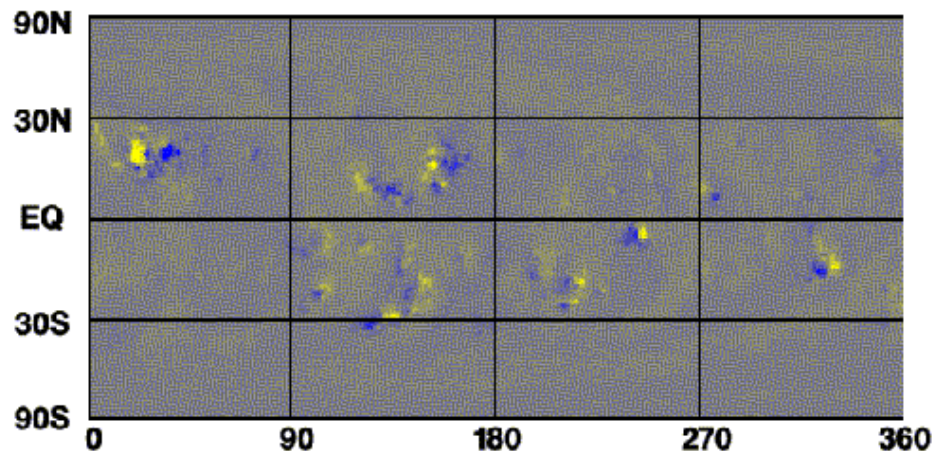
Hale's Polarity Law:

The polarity of the leading spots in one hemisphere is opposite that of the leading spots in the other hemisphere and the polarities reverse from one cycle to the next.

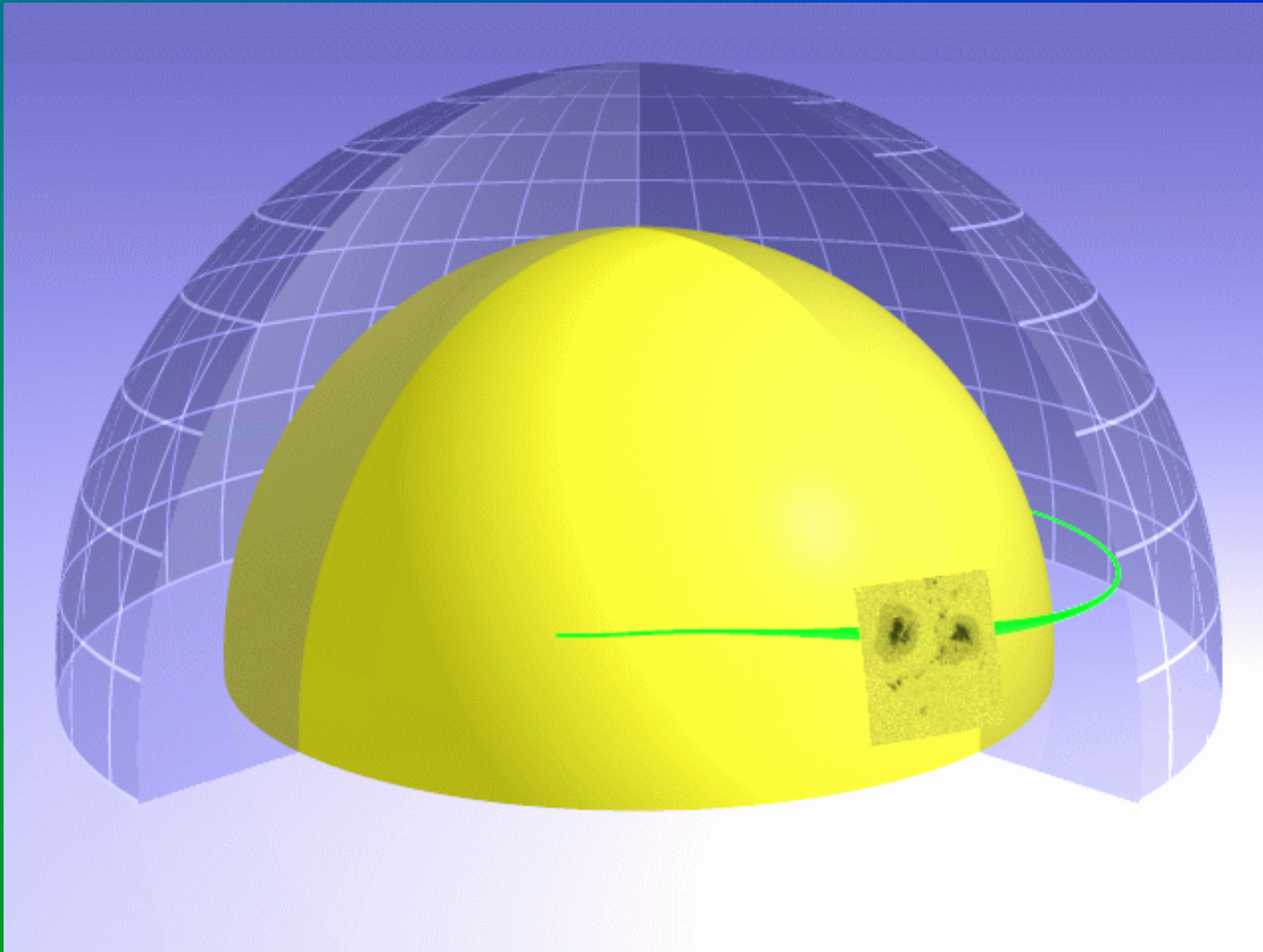
Cycle 21
Maximum



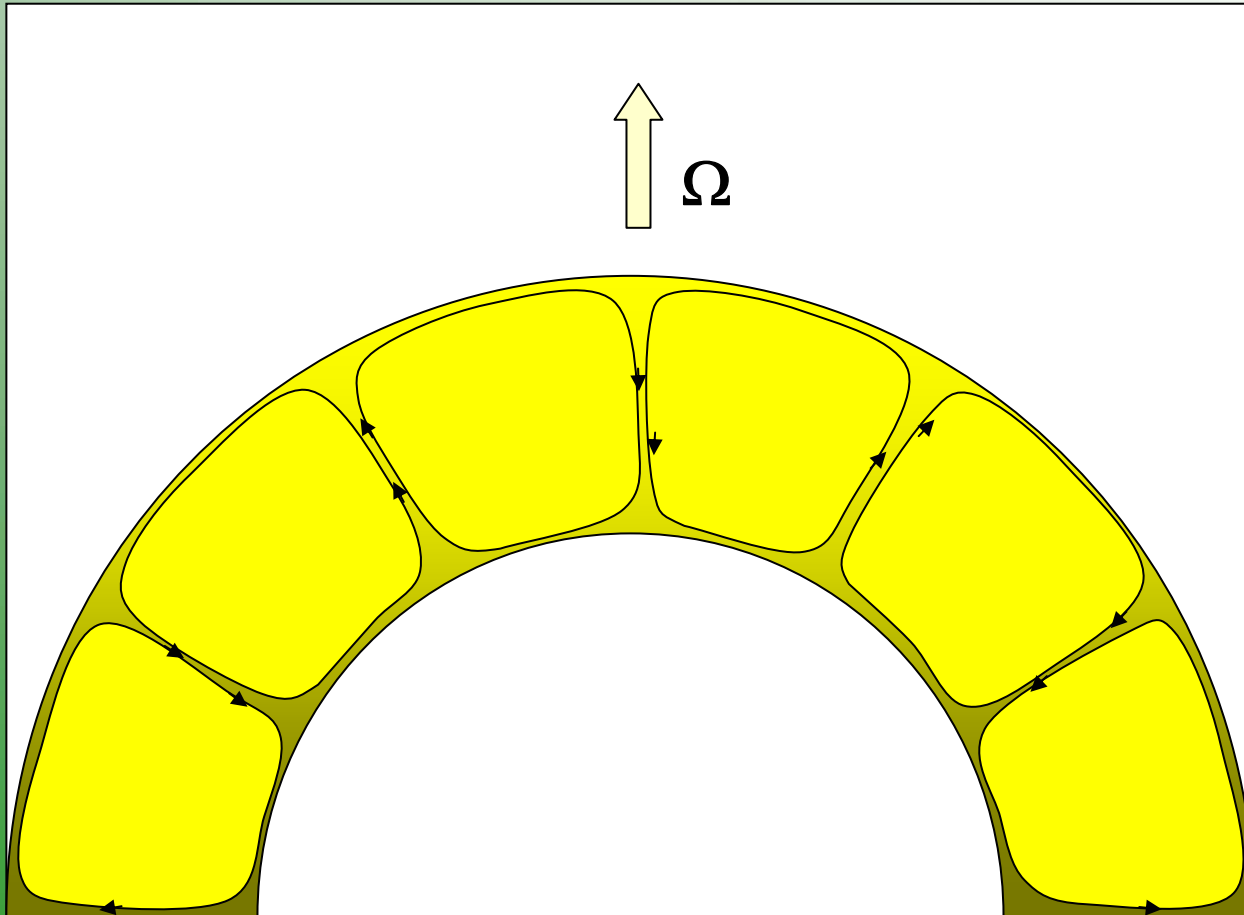
Cycle 22
Maximum



Where do the surface fields come from?



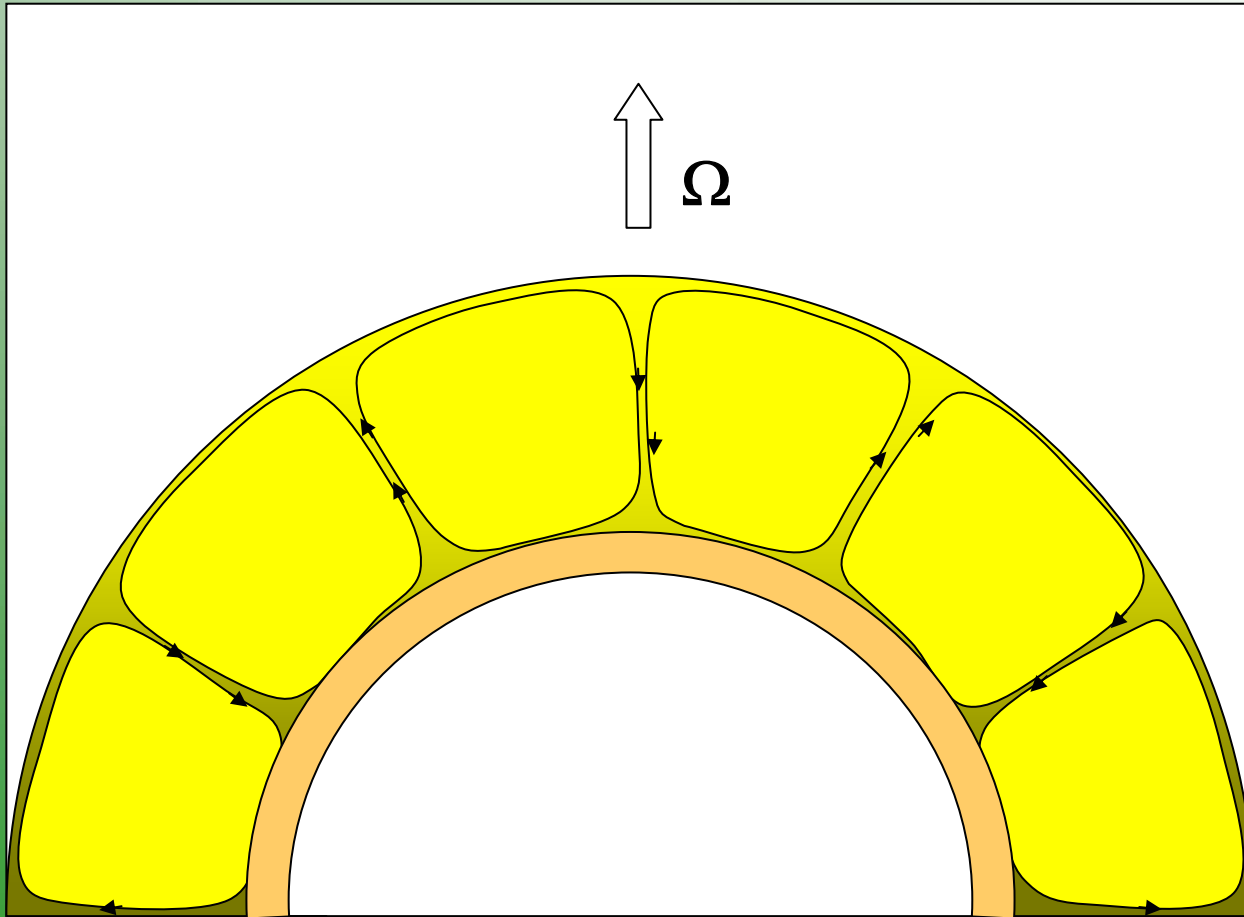
Origin of sunspots



1

•Convection zone

Origin of sunspots

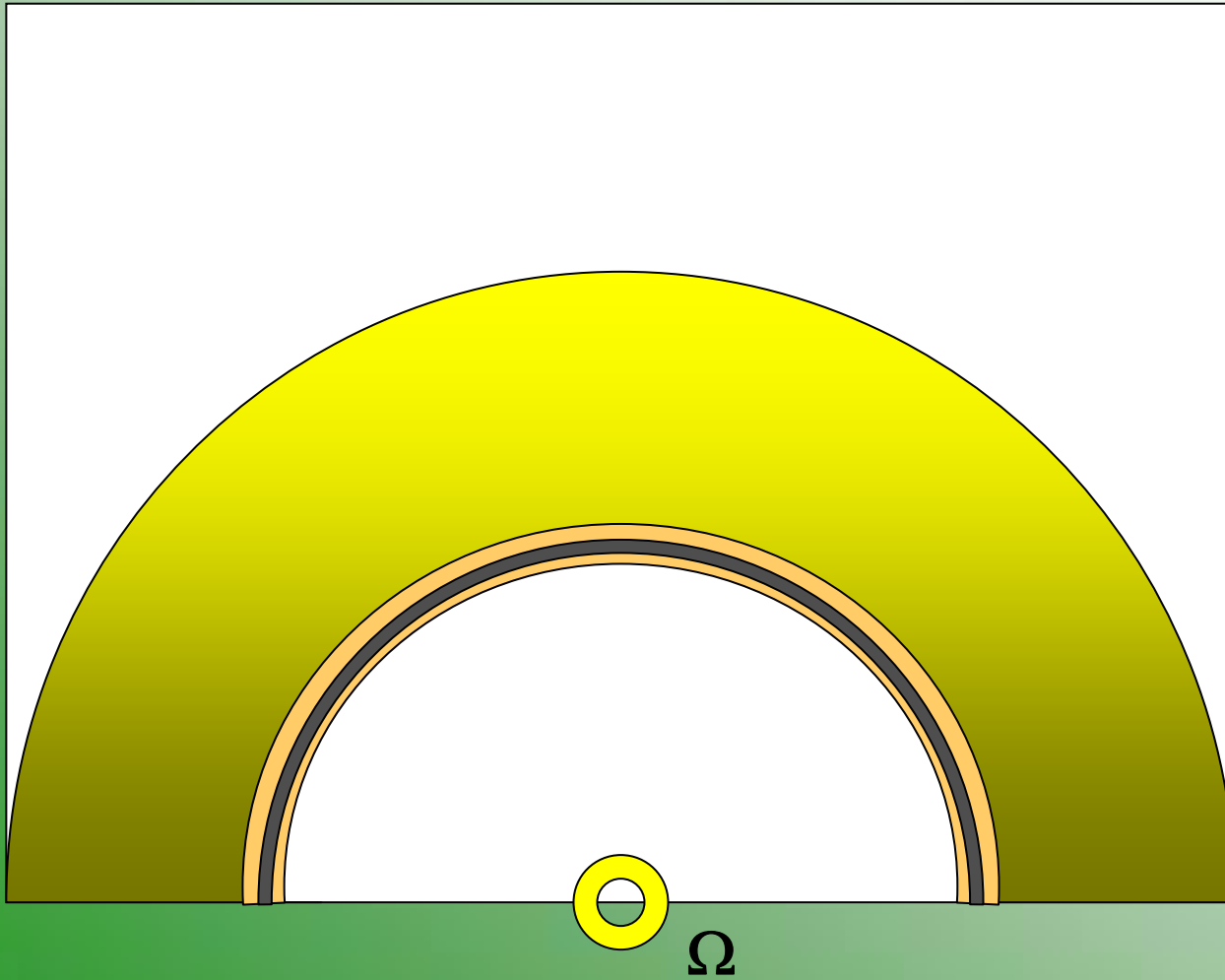


2

•Overshoot layer

Origin of sunspots

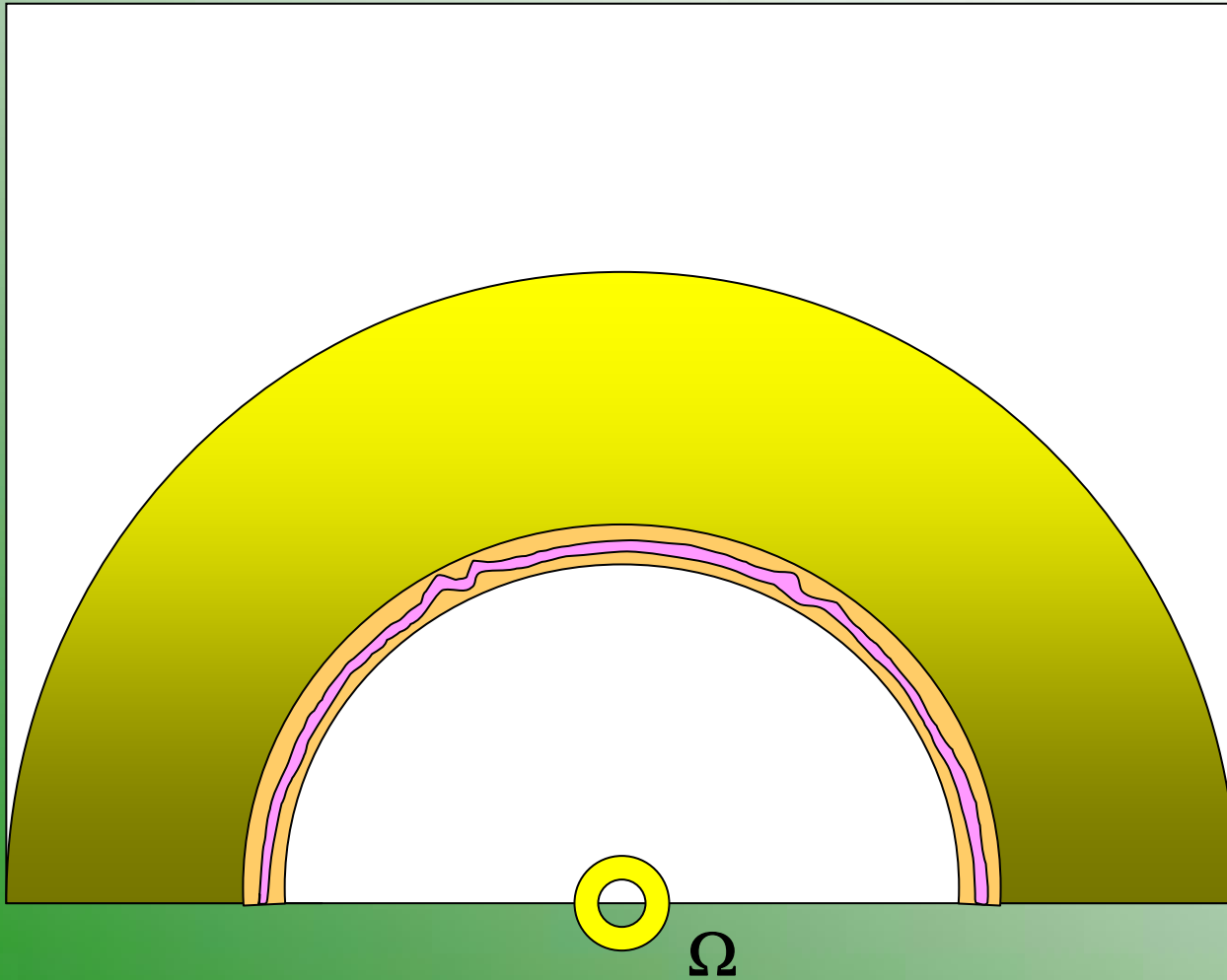
3



•Magnetic flux tube

Origin of sunspots

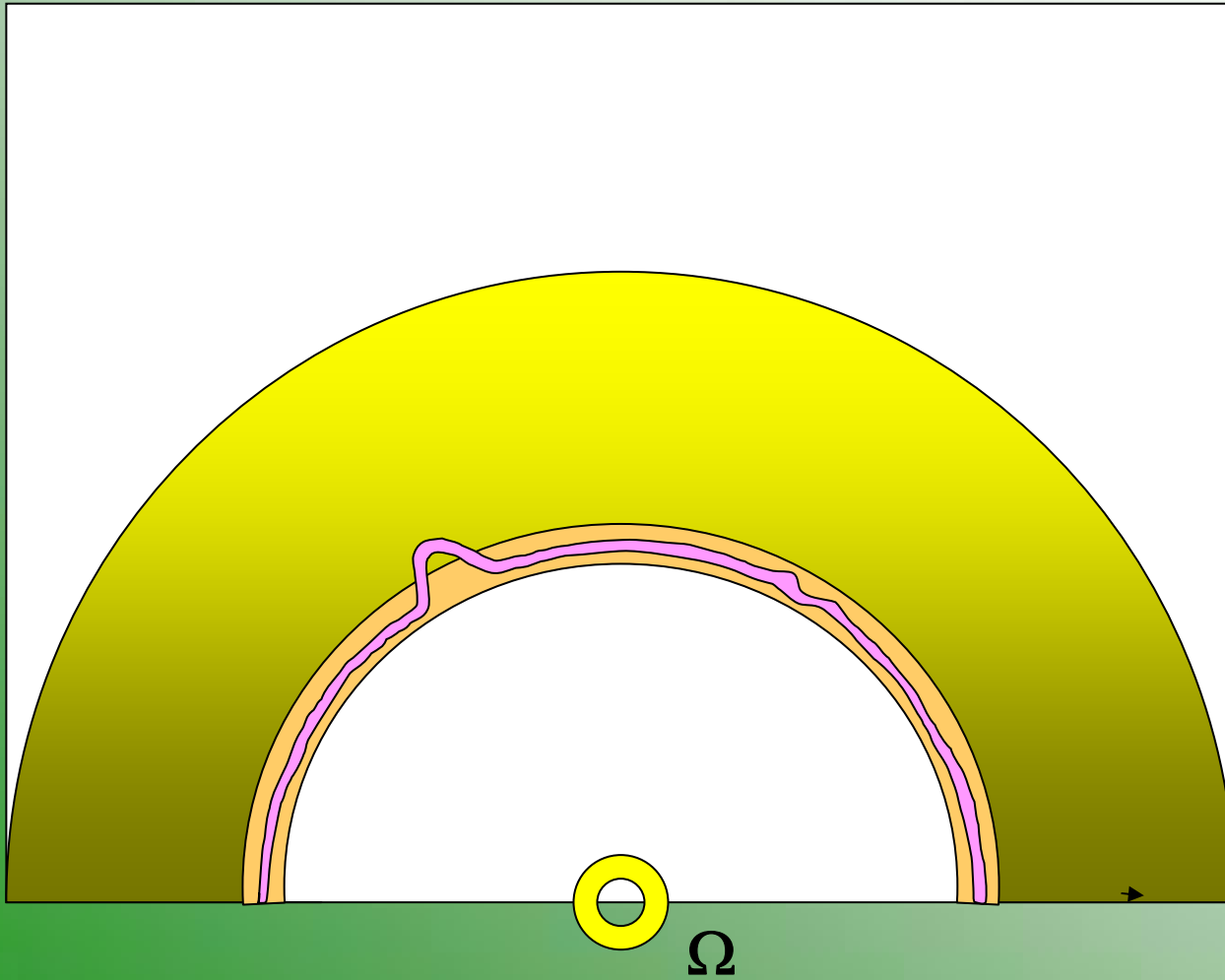
4



•Parker instability

Origin of sunspots

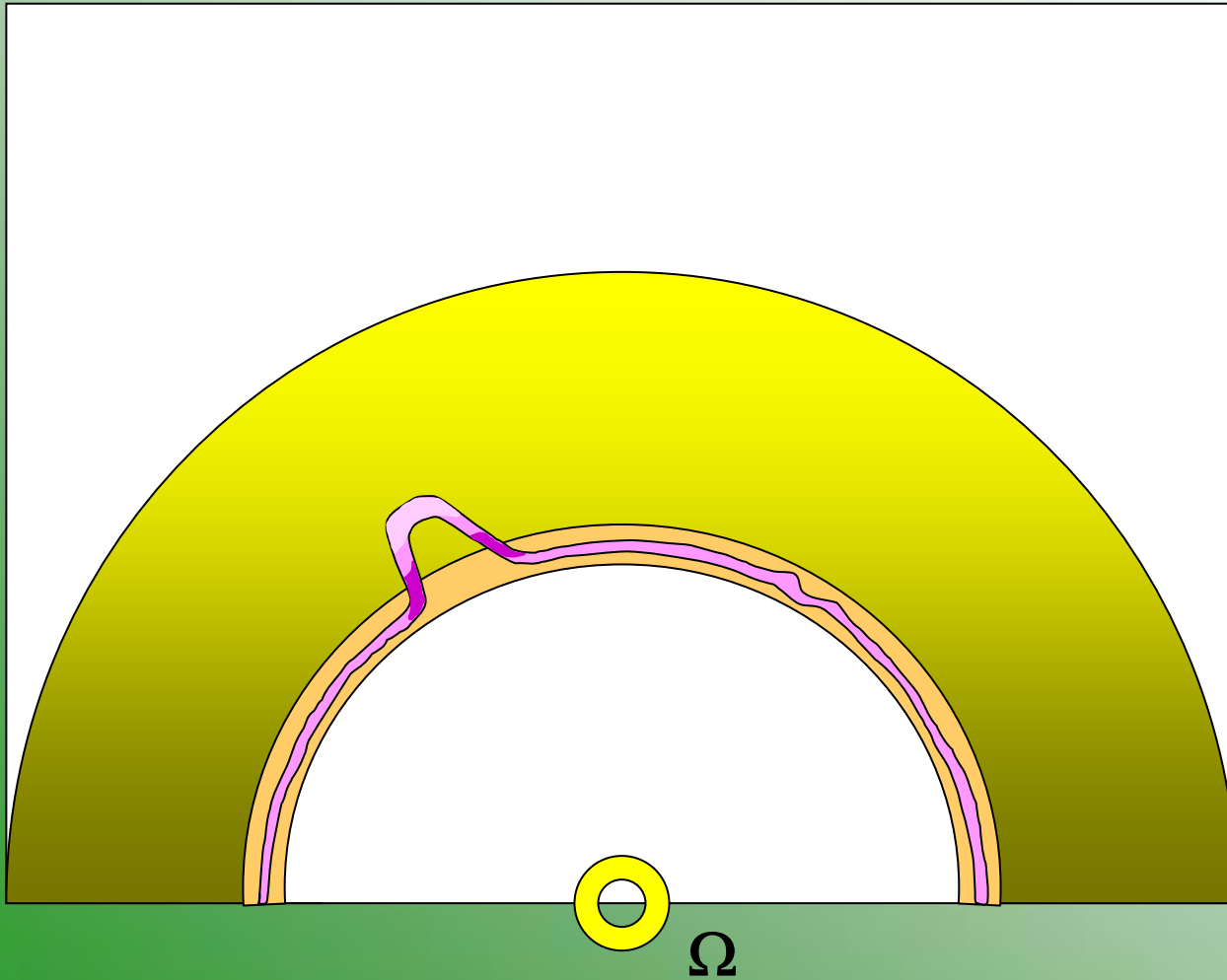
5



•Magnetic buoyancy

Origin of sunspots

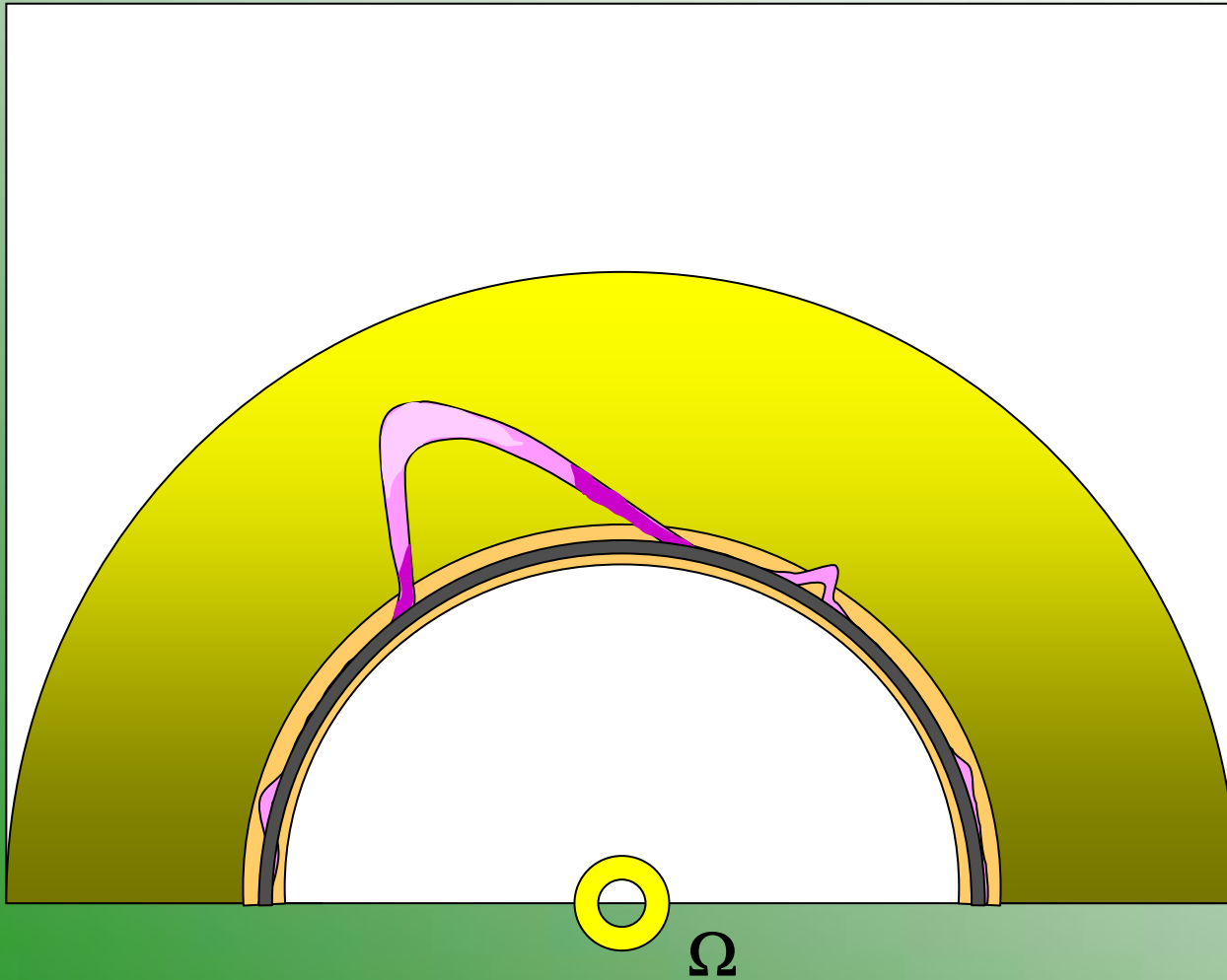
6



•Magnetic buoyancy

Origin of sunspots

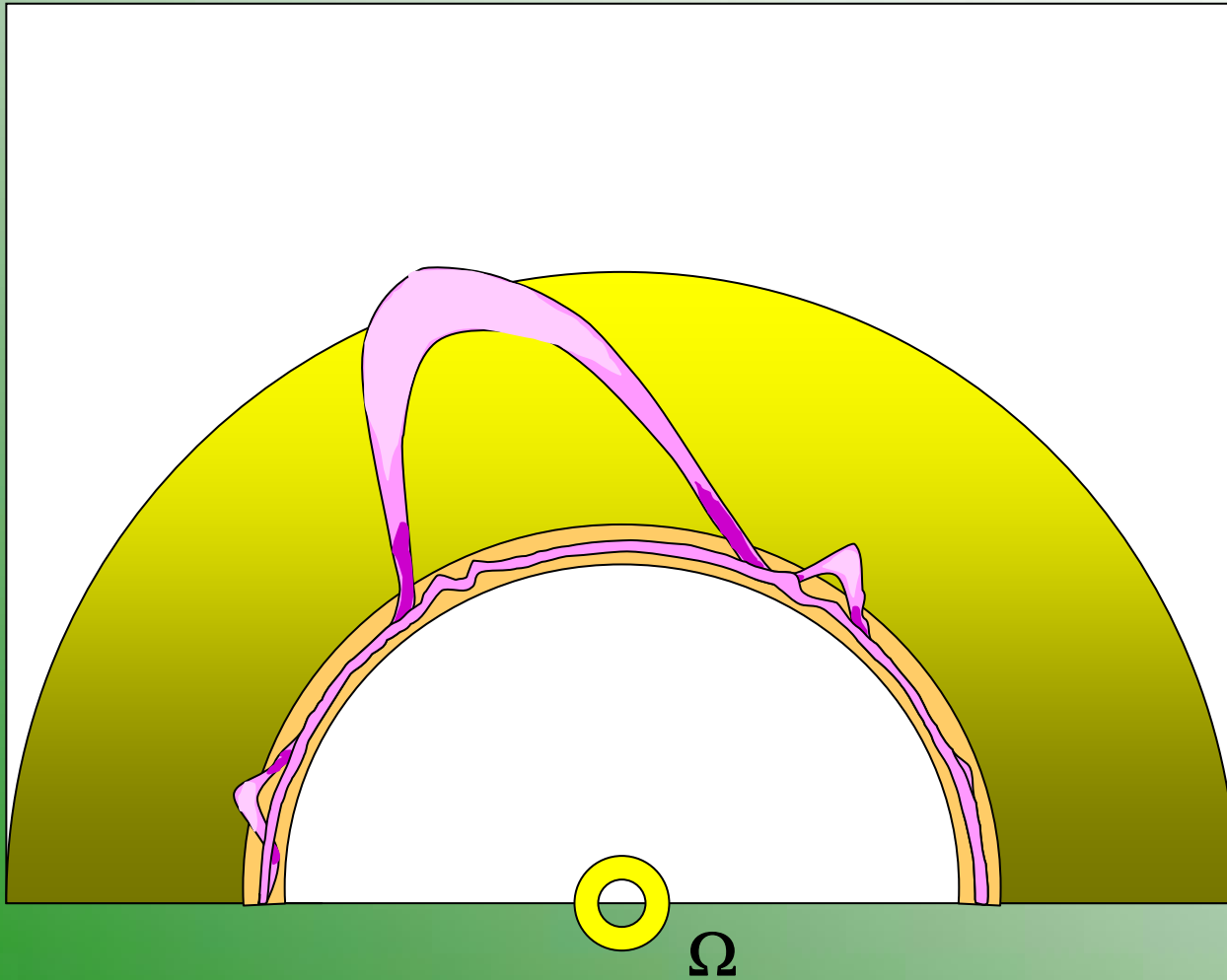
7



- Tube expansion and decreasing field strength

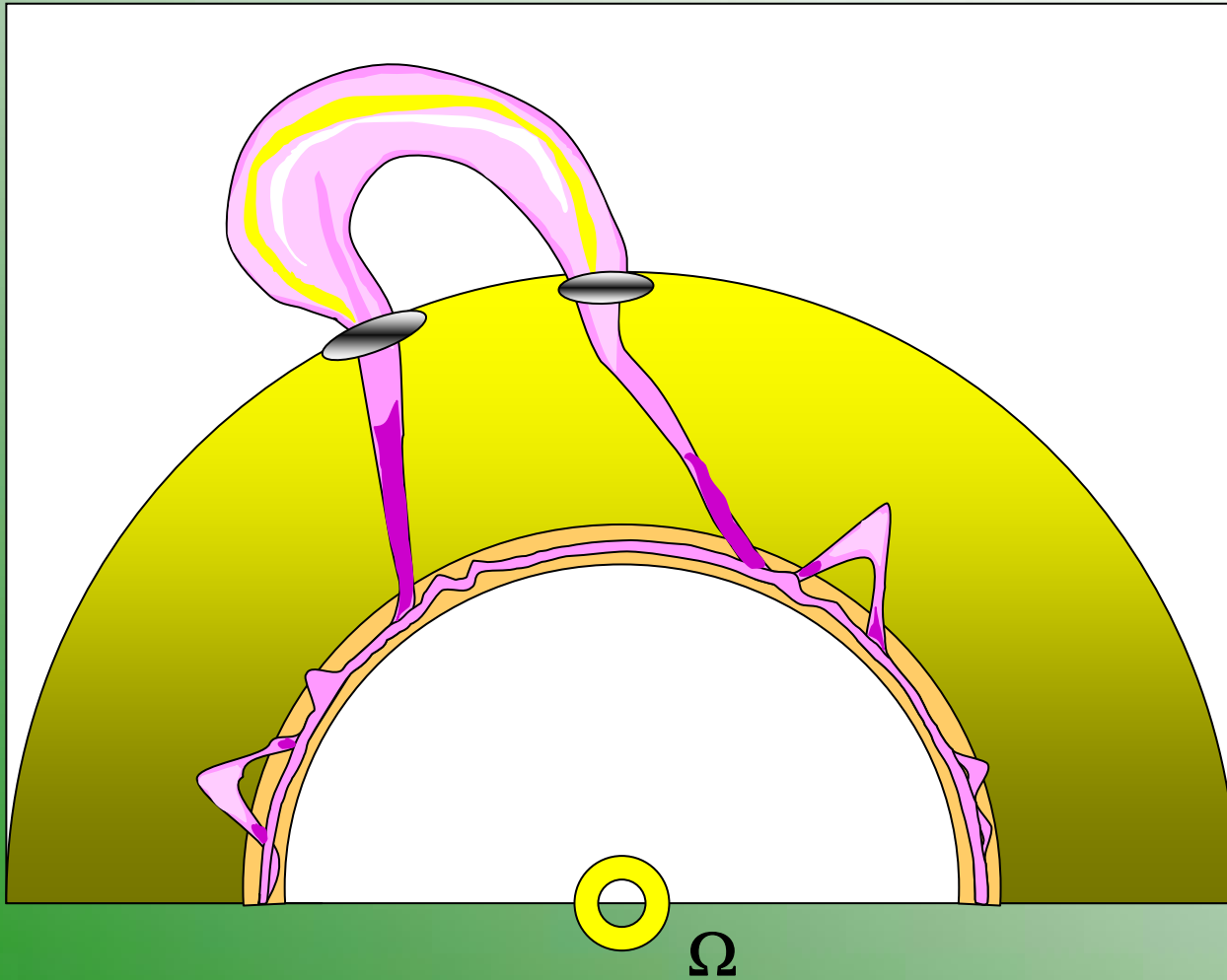
Origin of sunspots

8



•Eruption at the solar surface

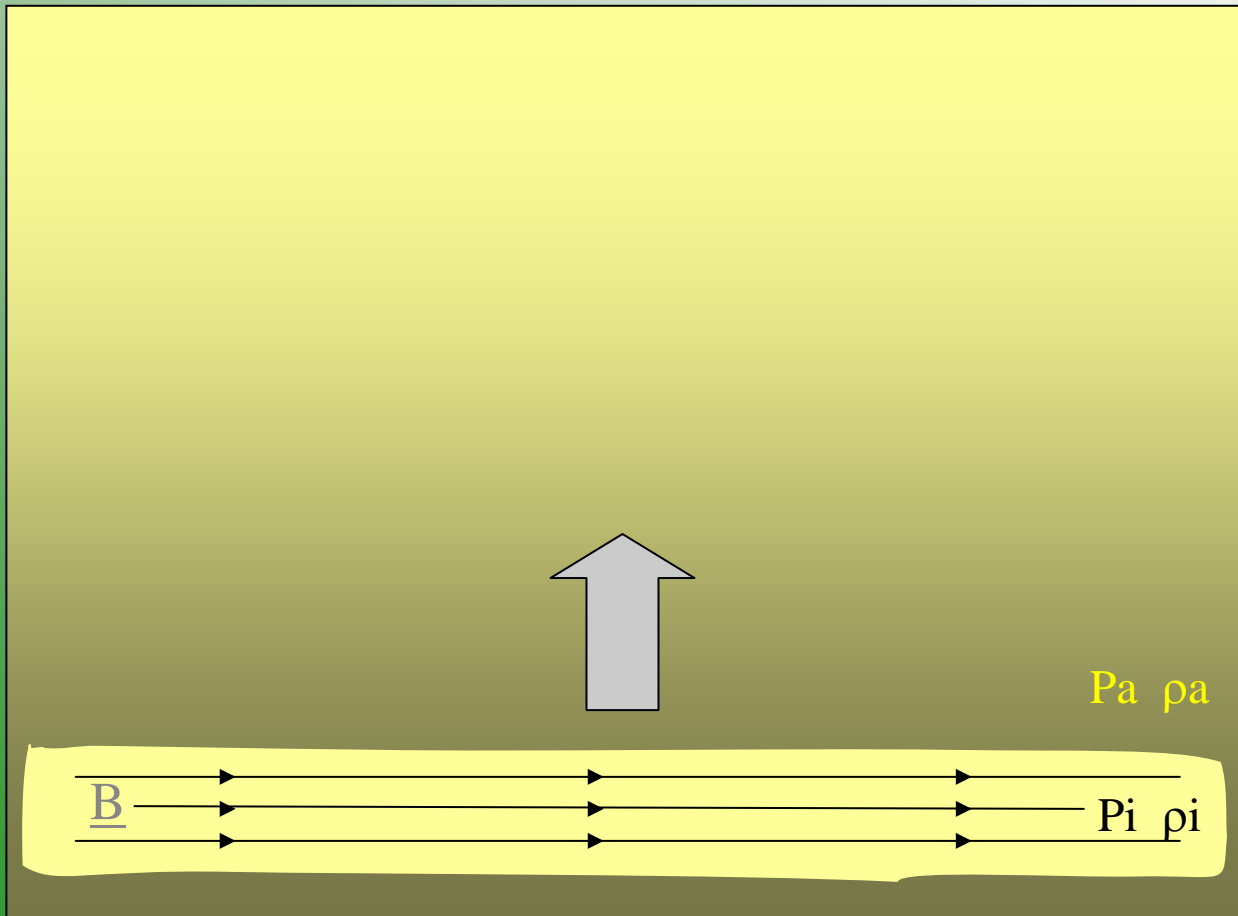
Origin of sunspots



9

•Formation of a bipolar sunspot pair/group

Magnetic buoyancy of a flux tube



Pressure equilibrium

$$P_a = P_i + B^2/8\pi$$

$$B \neq 0 \Rightarrow P_i < P_a$$

$$\Rightarrow \rho_i < \rho_a$$

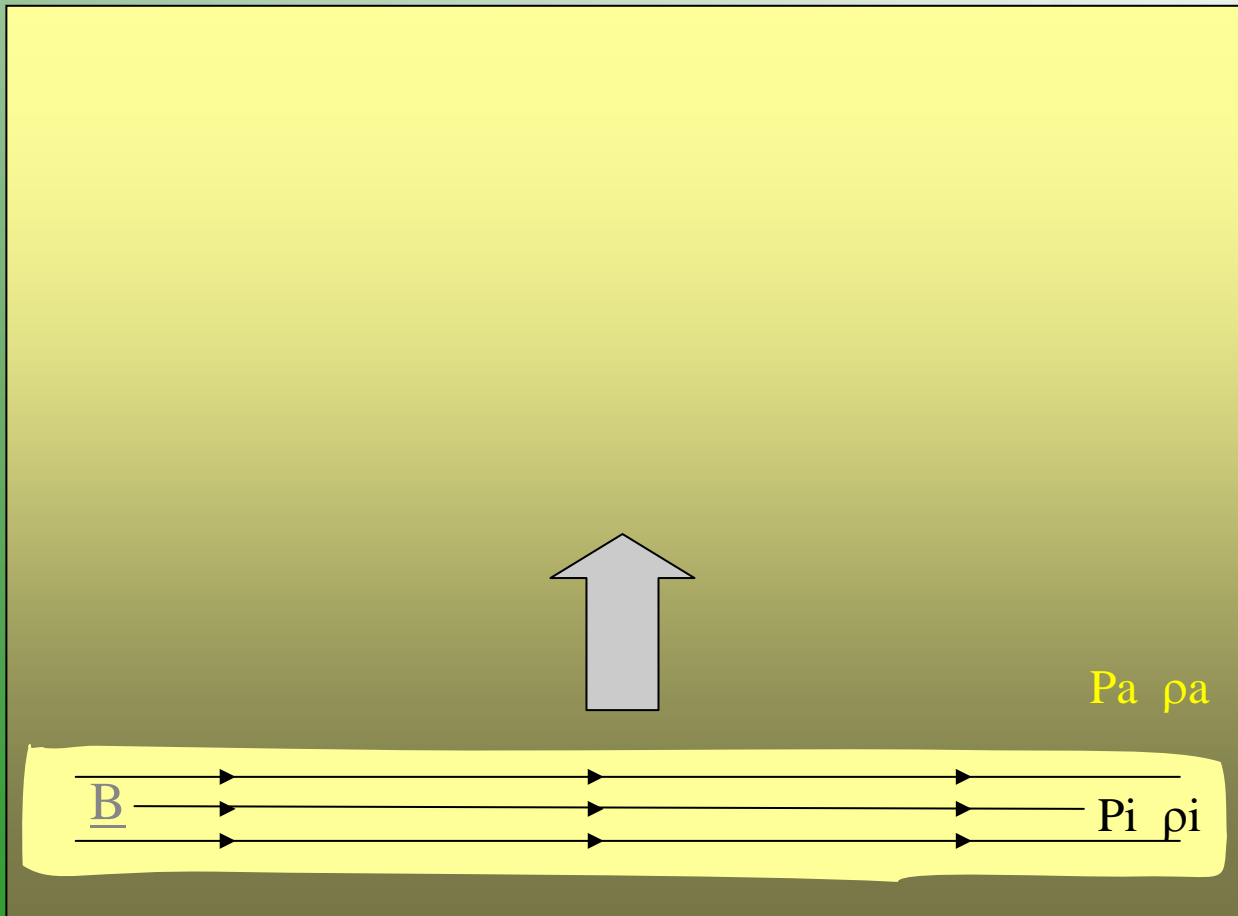
\Rightarrow buoyancy

P_a, ρ_a external pressure, density

P_i, ρ_i internal pressure, density

Parker instability

Magnetic buoyancy of a flux tube



Pressure equilibrium

$$P_a = P_i + B^2/8\pi$$

$$B \neq 0 \Rightarrow P_i < P_a$$

$$\Rightarrow \rho_i < \rho_a$$

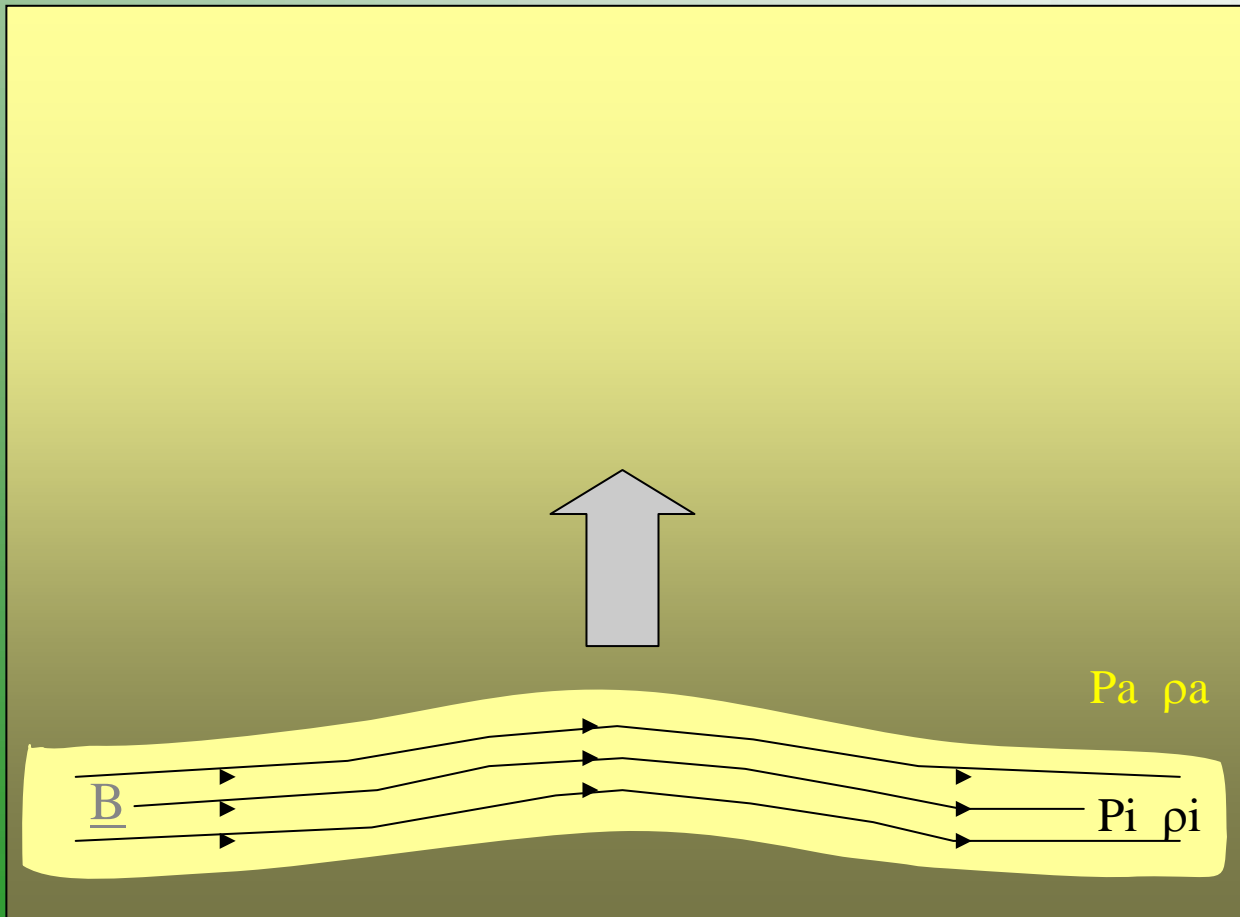
\Rightarrow buoyancy

P_a, ρ_a external pressure, density

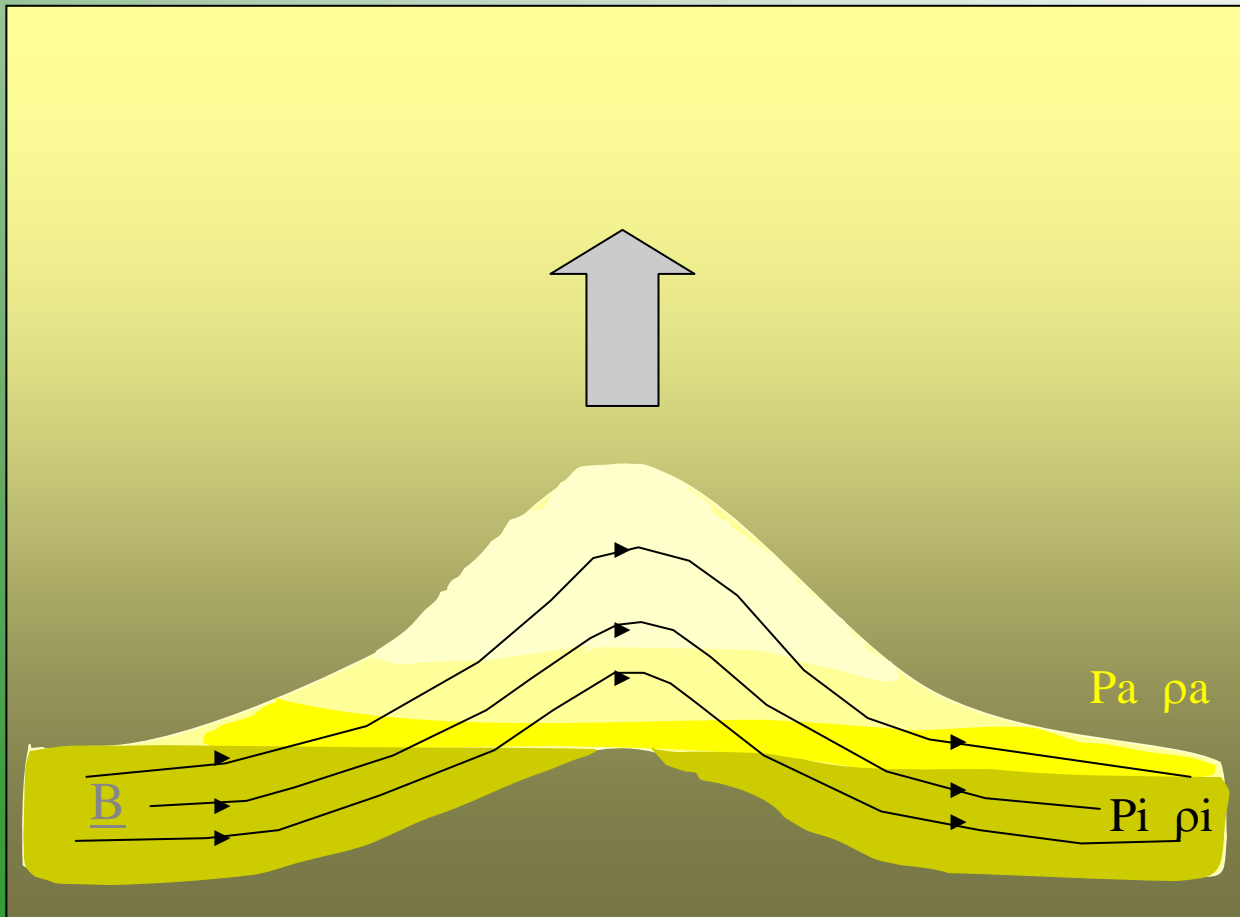
P_i, ρ_i internal pressure, density

Parker instability

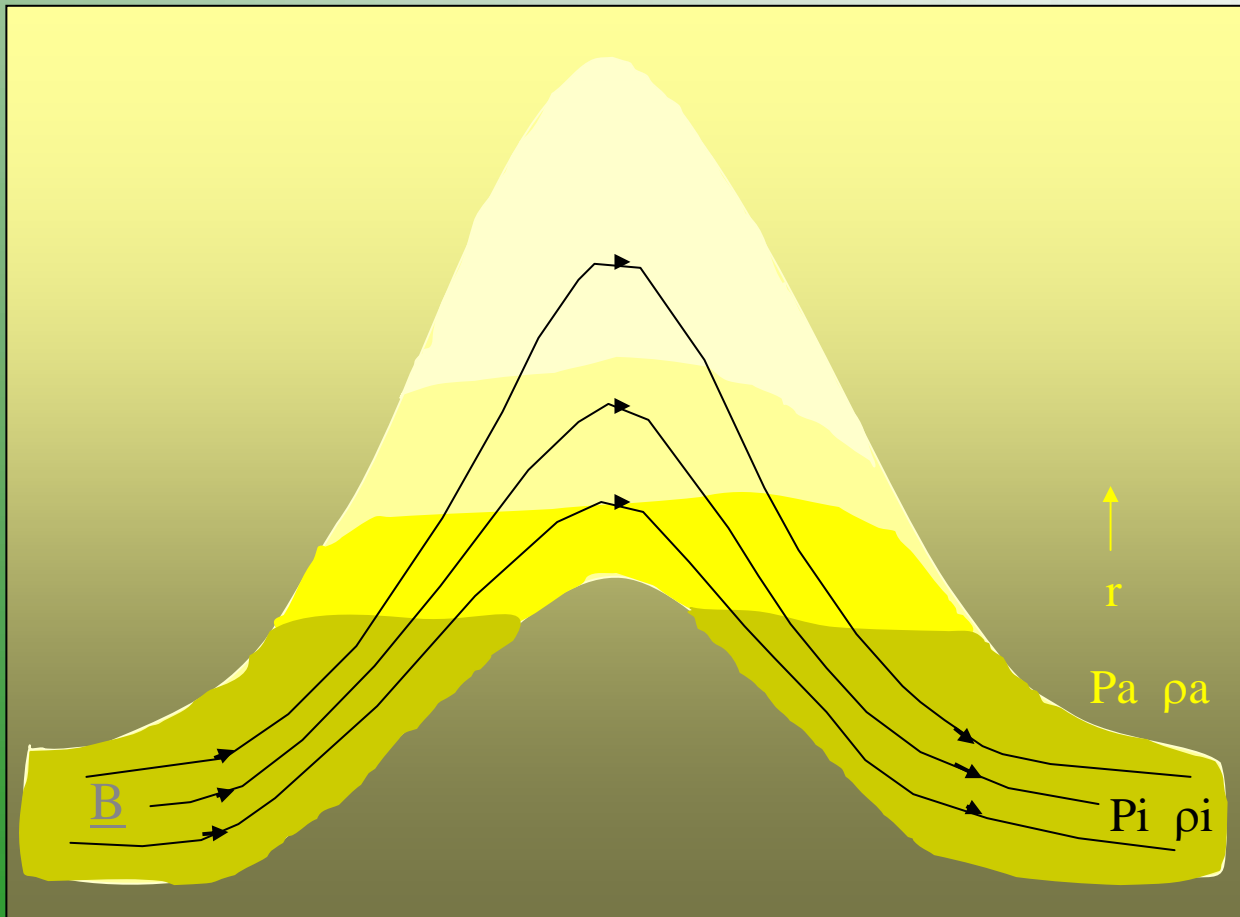
Magnetic buoyancy of a flux tube



Magnetic buoyancy of a flux tube



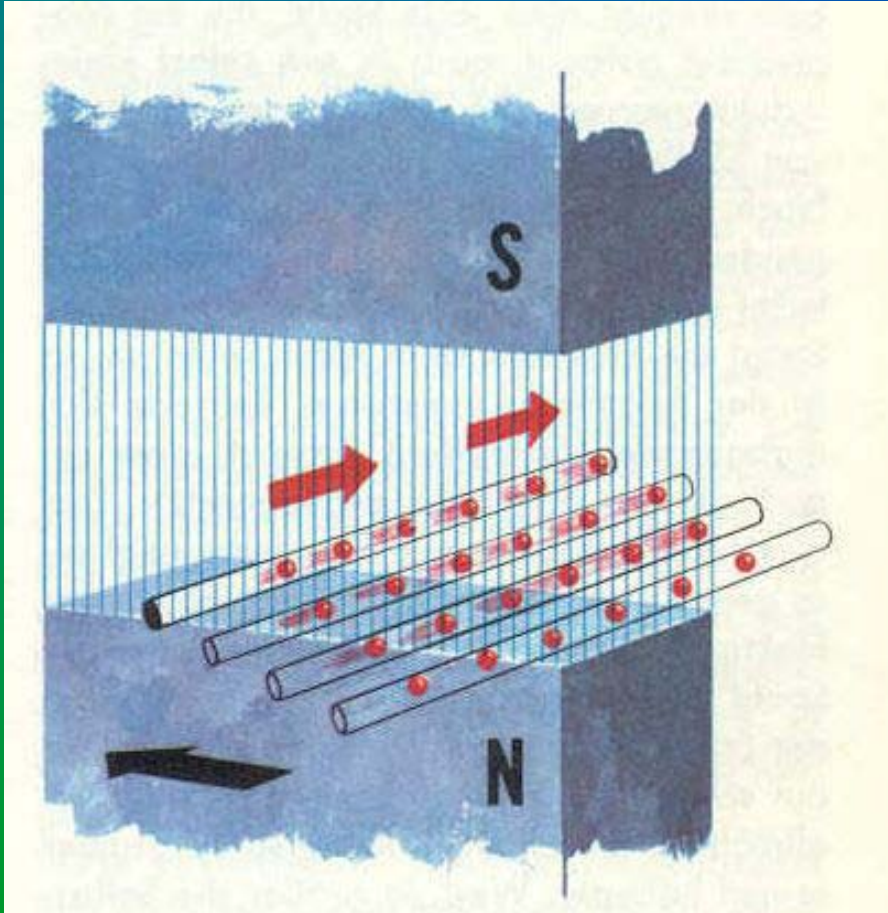
Magnetic buoyancy of a flux tube



Generation of magnetic flux...

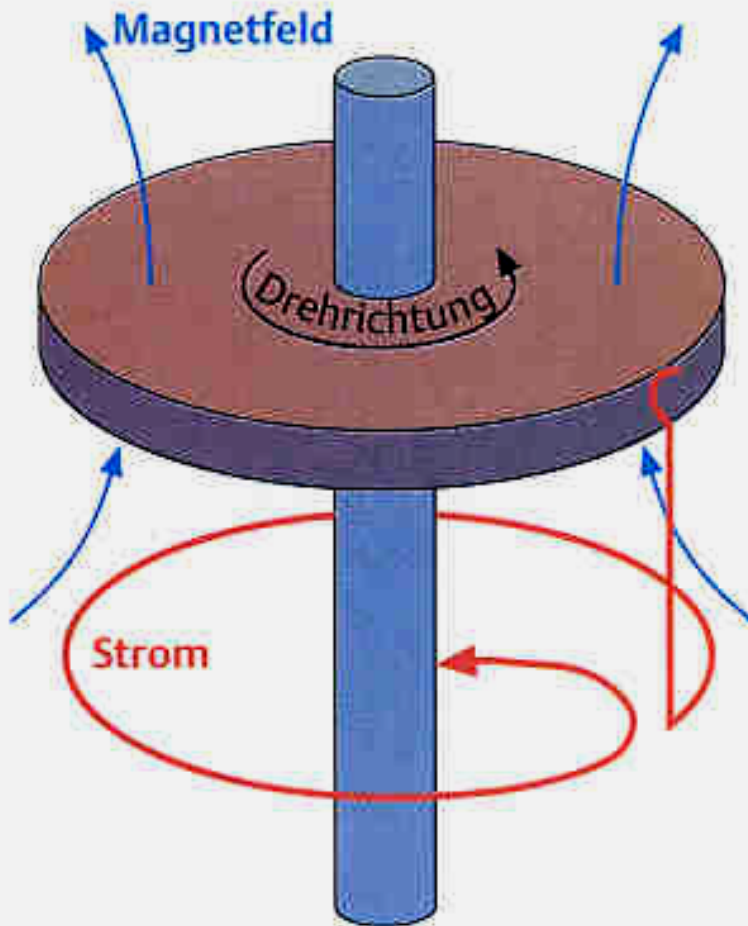
- ... requires an electrically conducting medium
 - ➔ plasma (ionized gas)
- ... requires fluid motion for induction
 - ➔ convective flows
 - ➔ (differential) rotation
- ... how is the field maintained against dissipation?
 - ➔ (self-excited) dynamo process

The induction principle



- Conductor moving in a magnetic field
- perpendicular electrical field and force
- electrical current
- new magnetic field
- **Lenz's rule!**
(no perpetuum mobile)

A simple dynamo



Initially weak "seed field"

→ Rotation induces electrical field between axis and edge

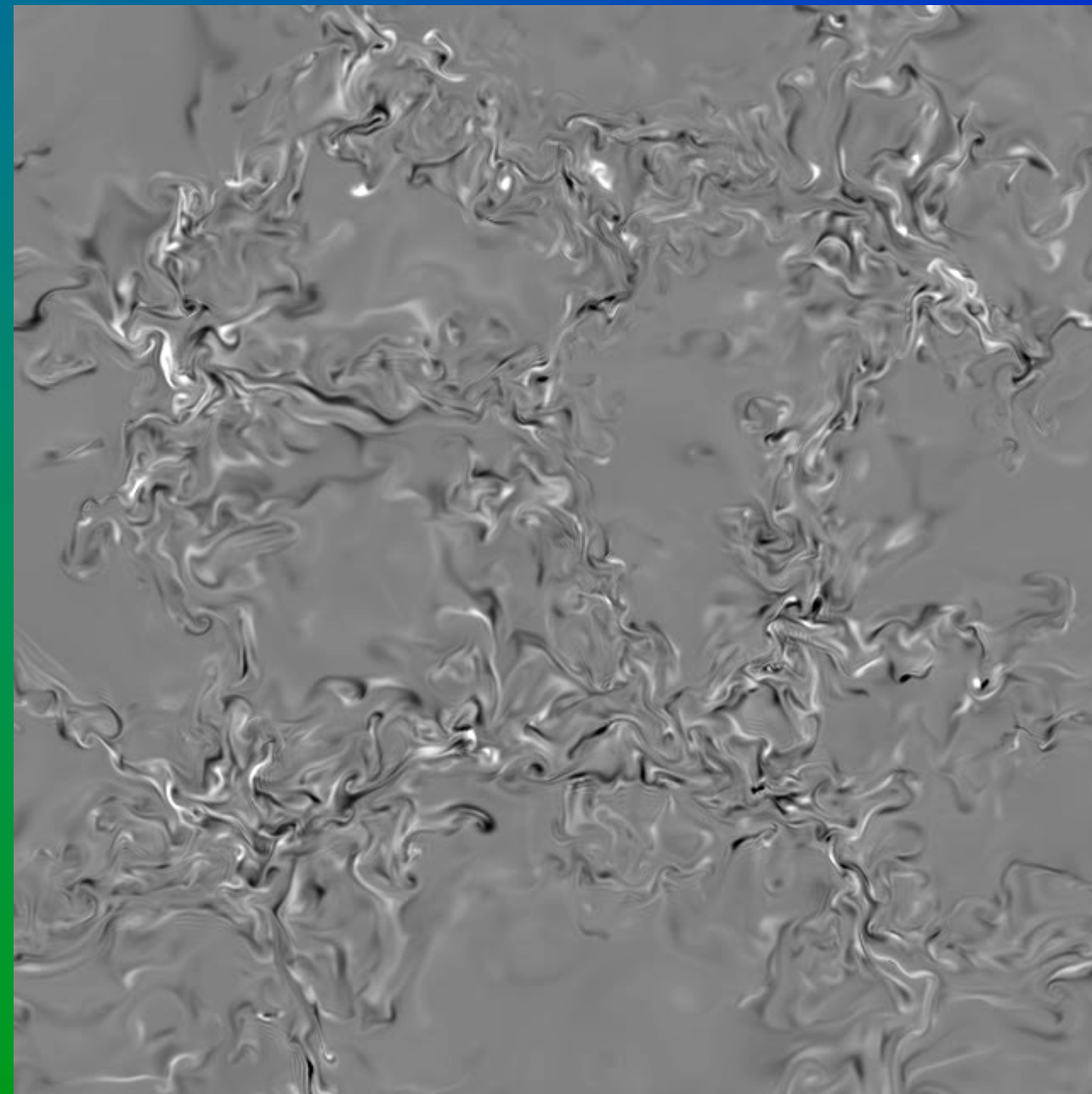
→ Current closed by wire

→ Current generates a magnetic field which amplifies the seed field

→ Sun: no isolated wires

→ "homogeneous dynamo"

Local dynamo

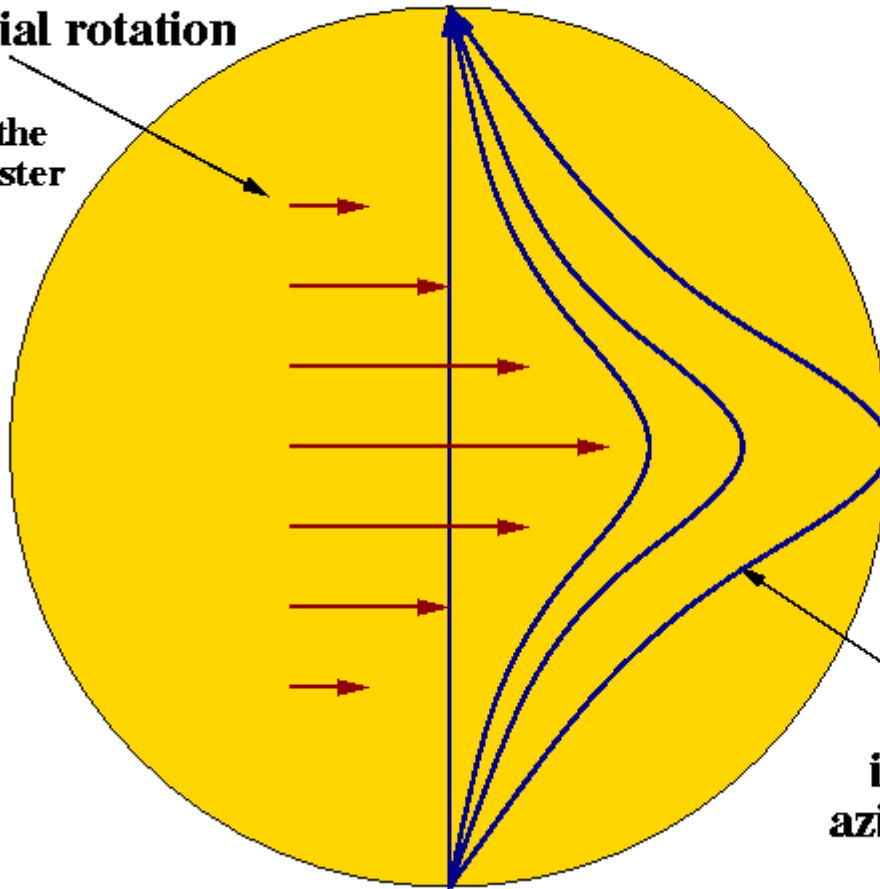


Vögler & Sch. 2007

Differential rotation generates azimuthal (toroidal) magnetic field

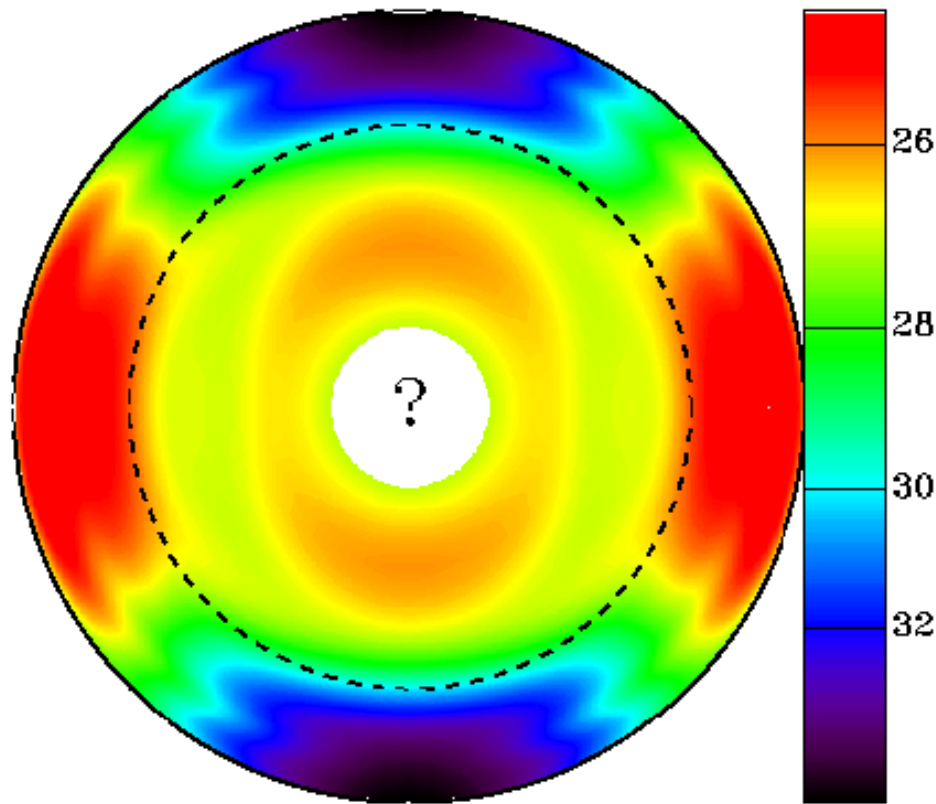
**Azimuthal flow
of differential rotation**

The longer the
arrow the faster
the flow



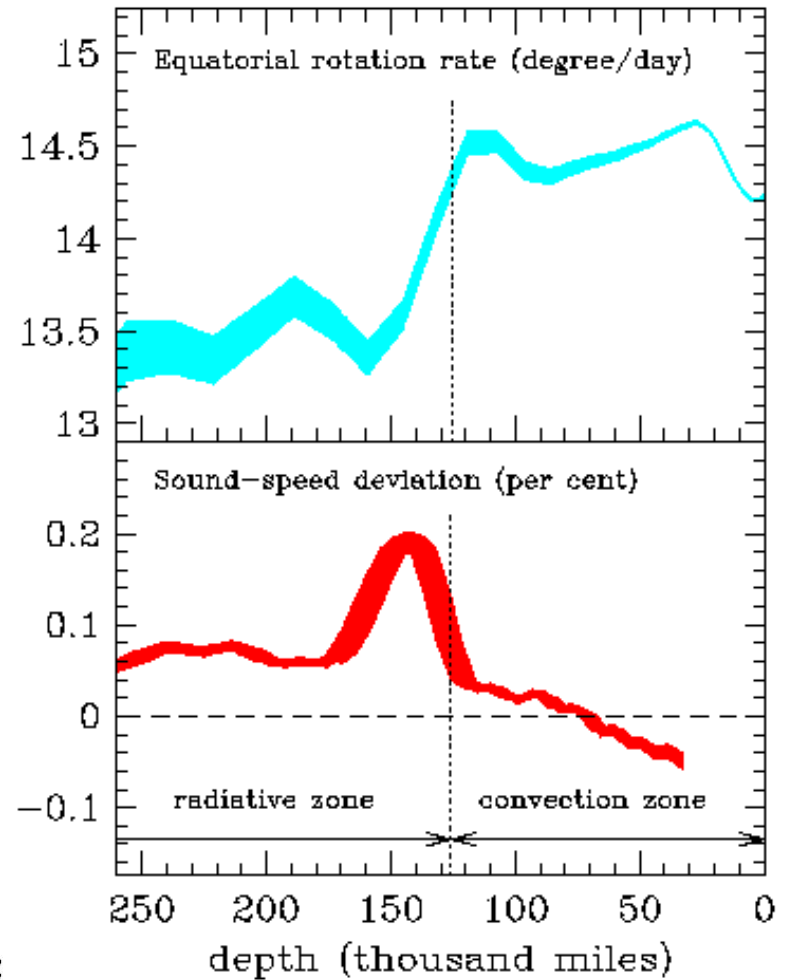
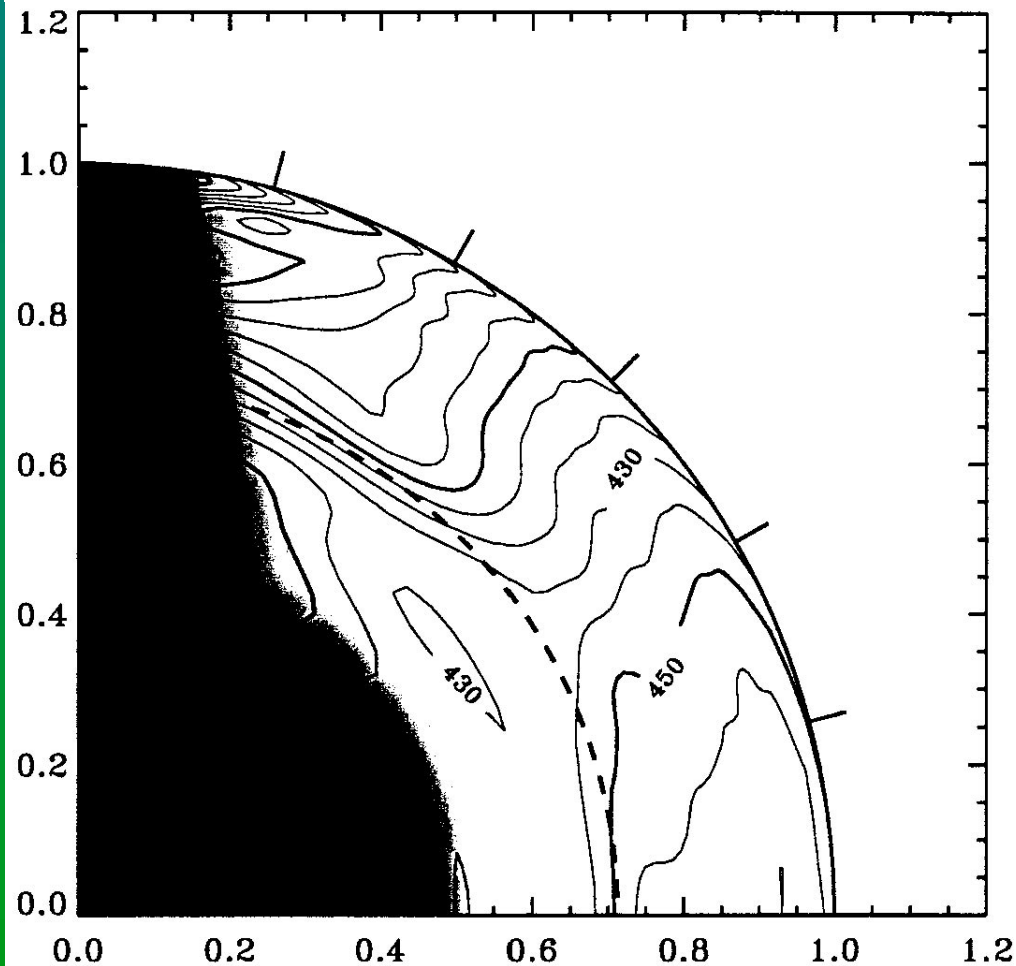
**Meridional
magnetic field
is transformed into
azimuthal magnetic field**

Internal rotation of the Sun as determined by helioseismology

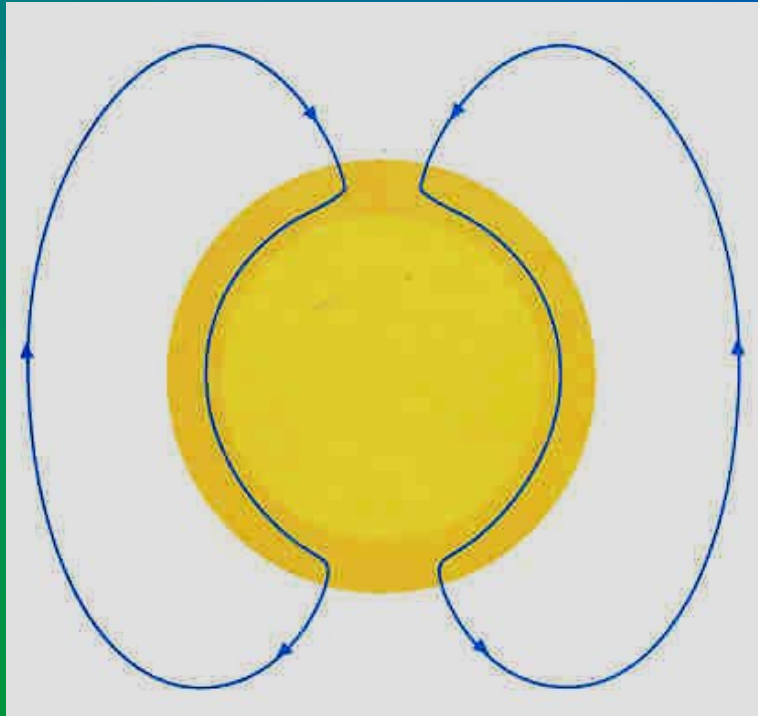


- Convection zone rotates similar to surface
- Core rotates nearly rigidly
- Steep transition at the bottom of the convection zone; width $\sim 2\% R_{\text{sun}}$
- Region of strongest shear
→ Dynamo!

Internal rotation of the Sun as determined by helioseismology

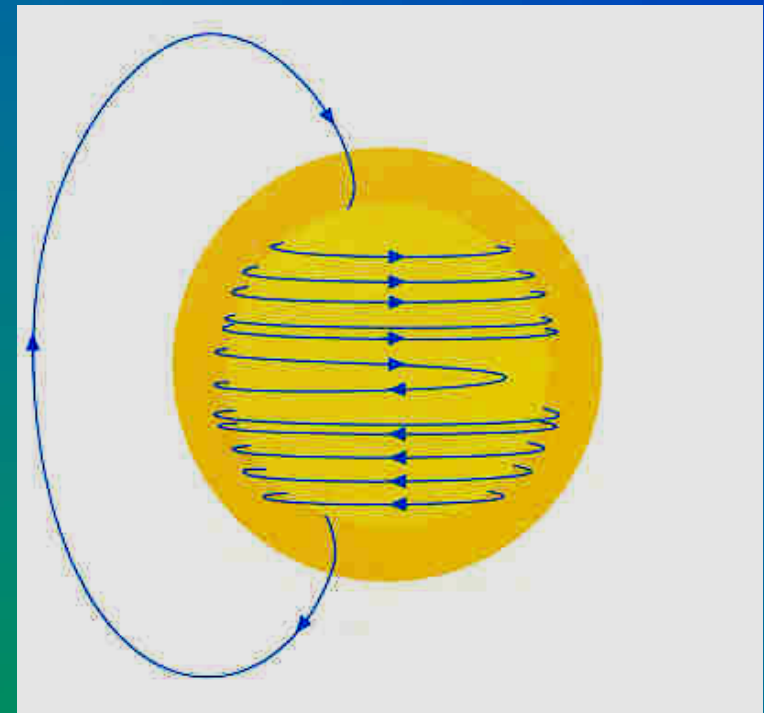


The solar dynamo (1)

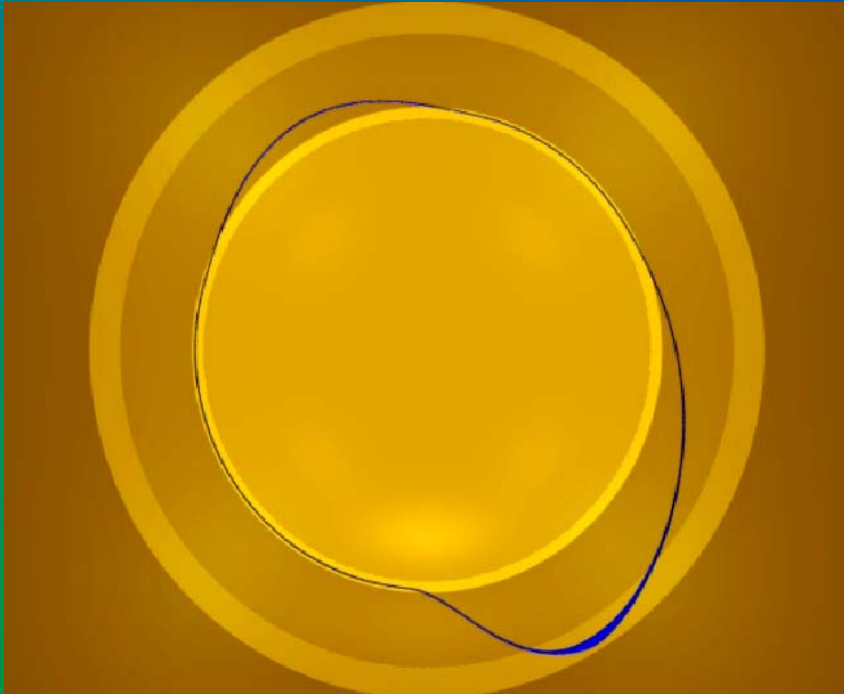


Dipole field in the convection zone

Winding up of the field by differential rotation
→ strong toroidal field

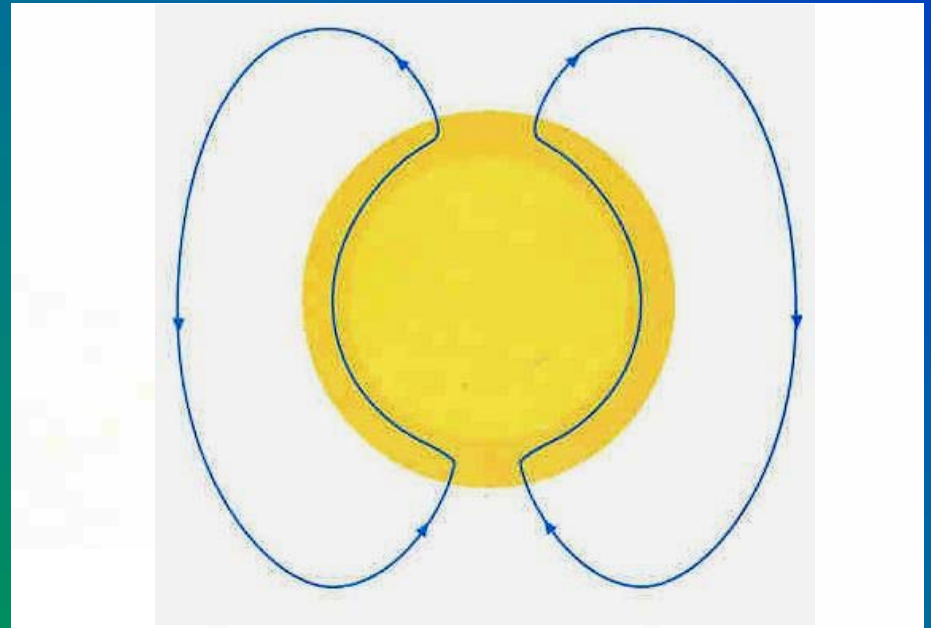


The solar dynamo (2)

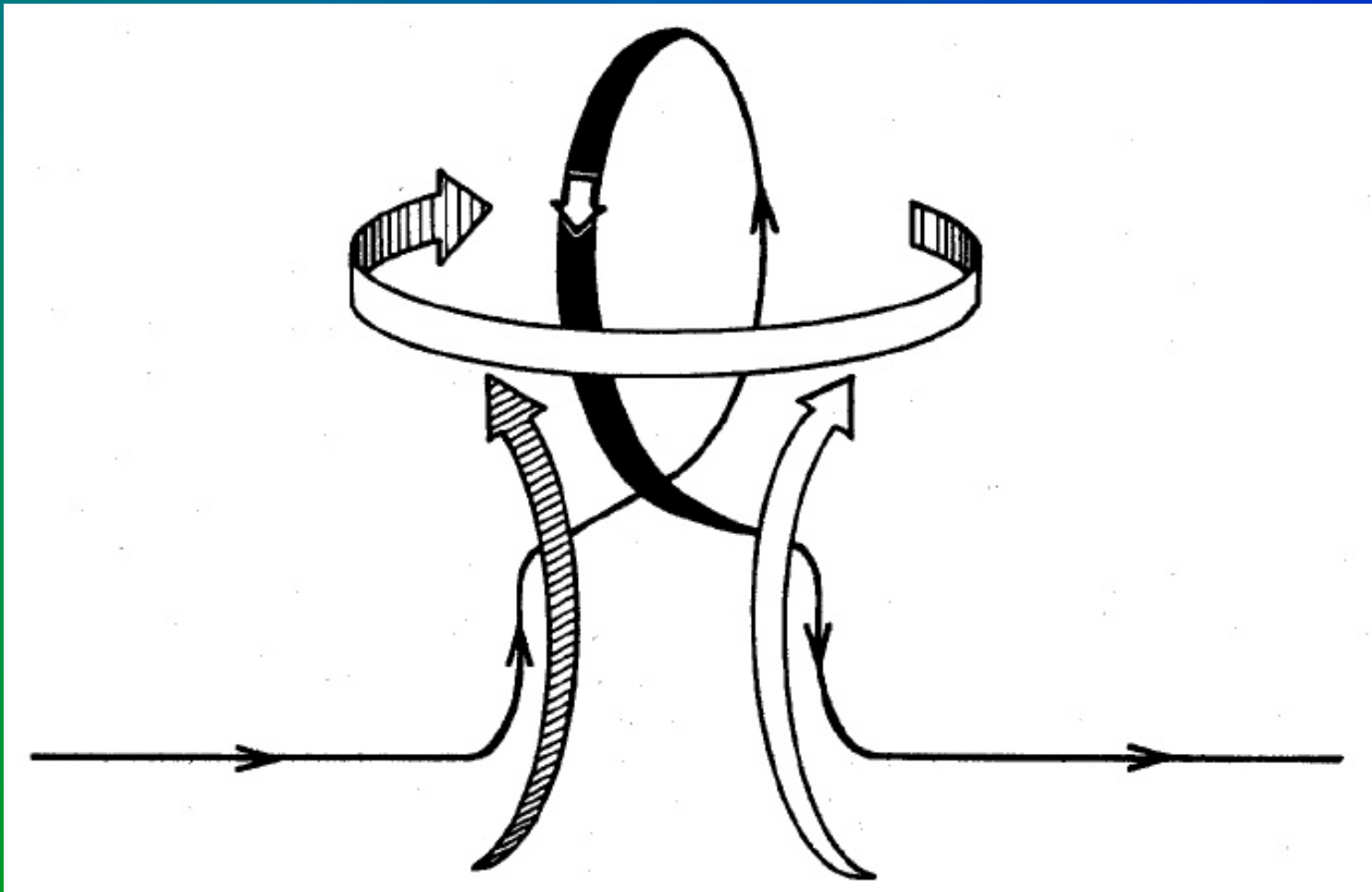


Rise and eruption of magnetic flux tubes
→ sunspots

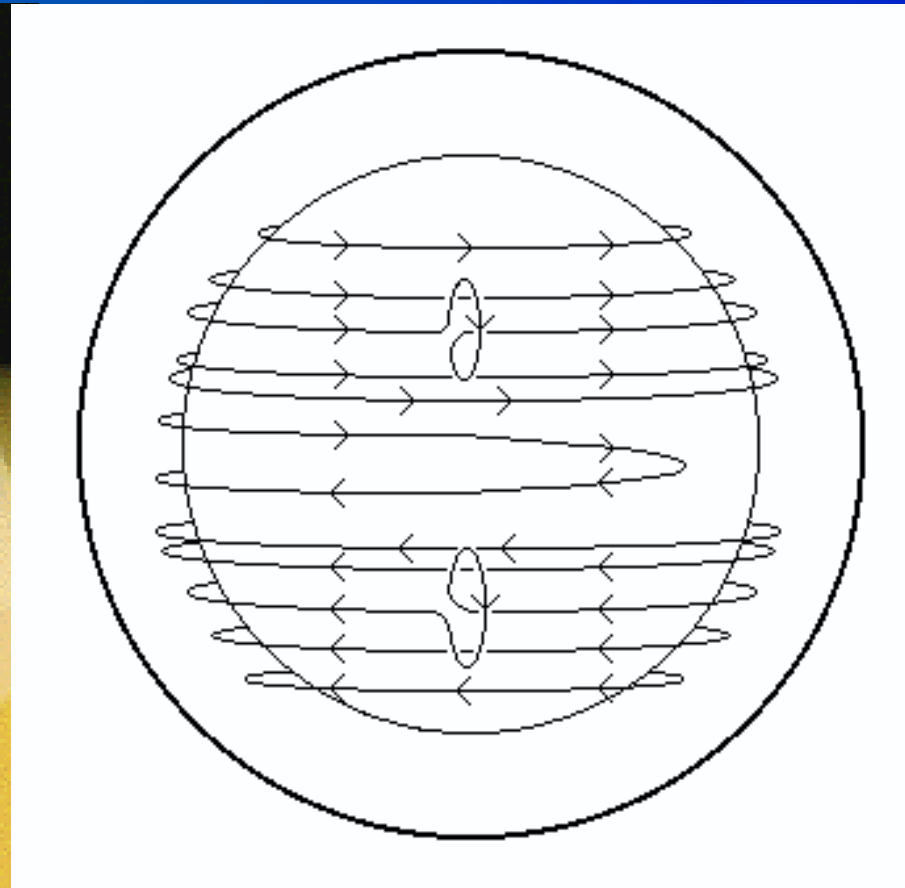
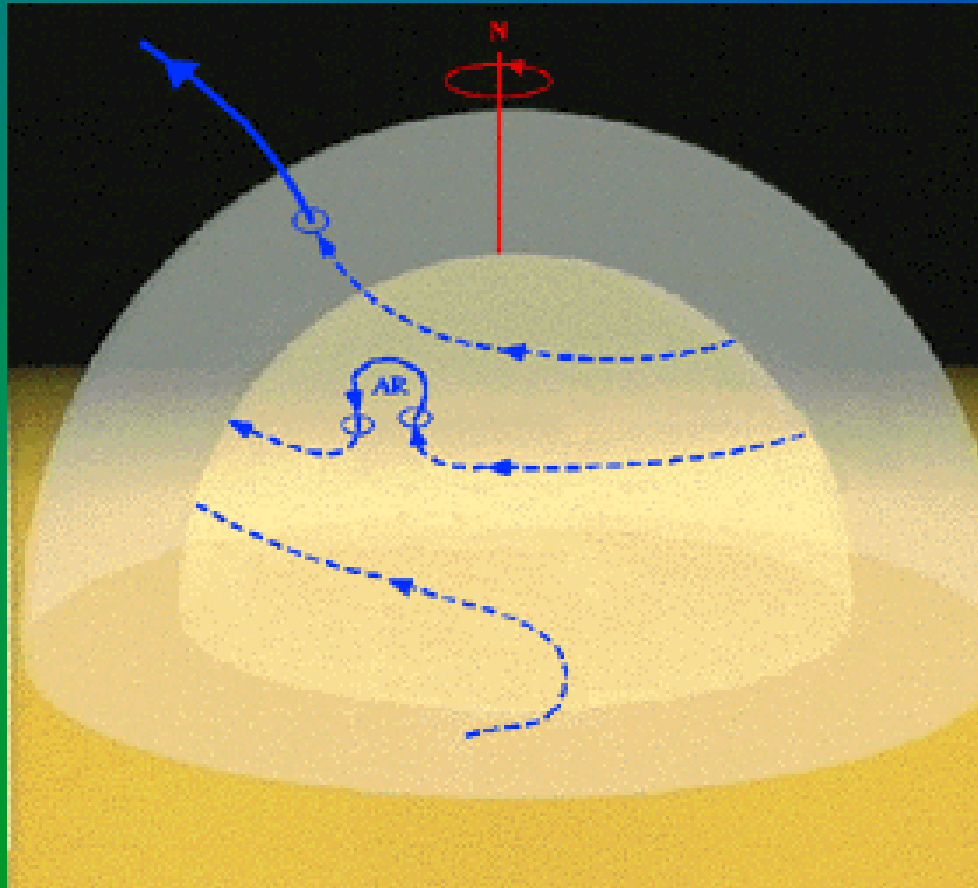
Twist of the erupting field
by Coriolis force
→ reversed dipole field



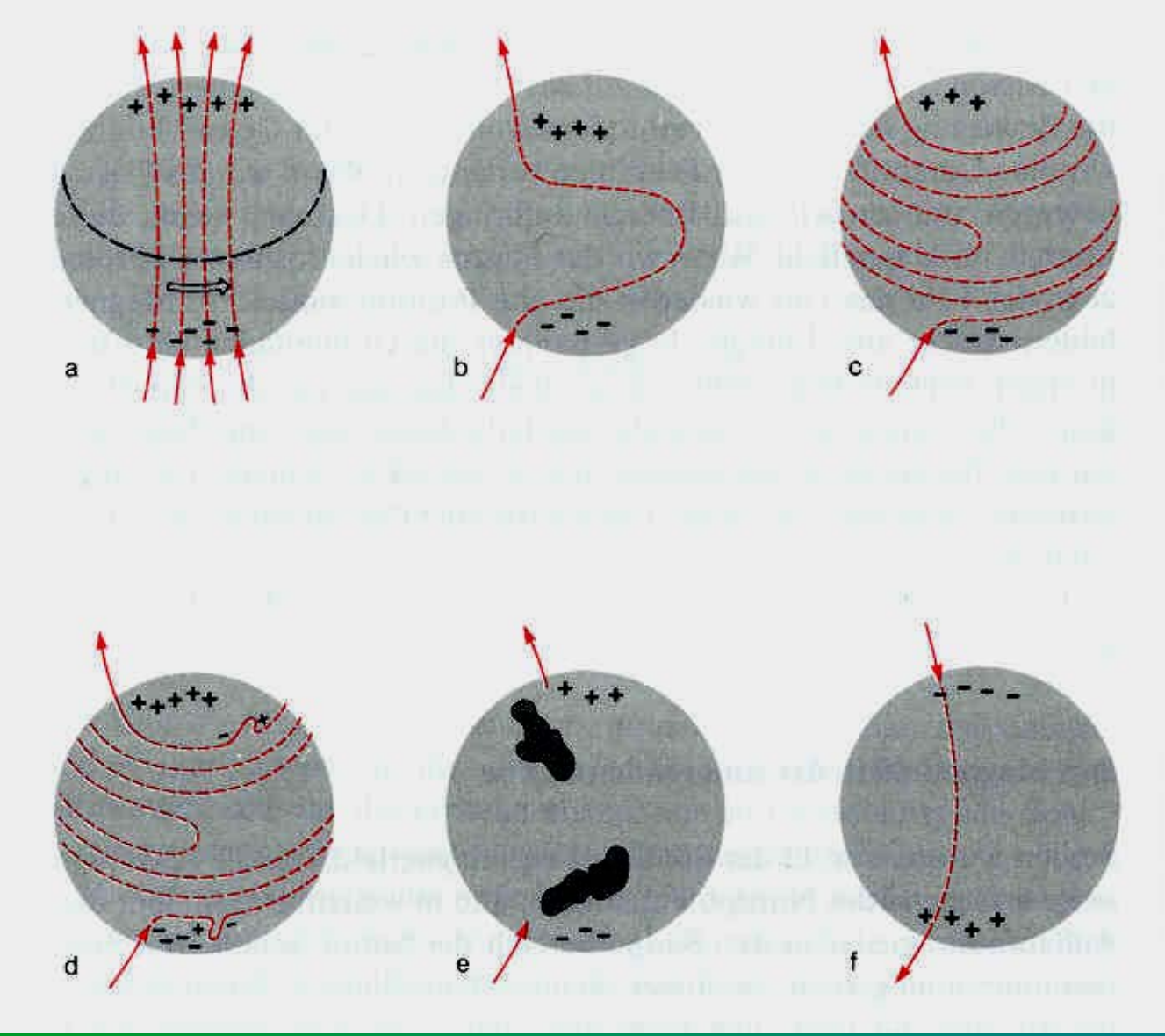
Twisting of a field line in a rising & expanding convective flow by the action of the Coriolis force (Parker, 1955)

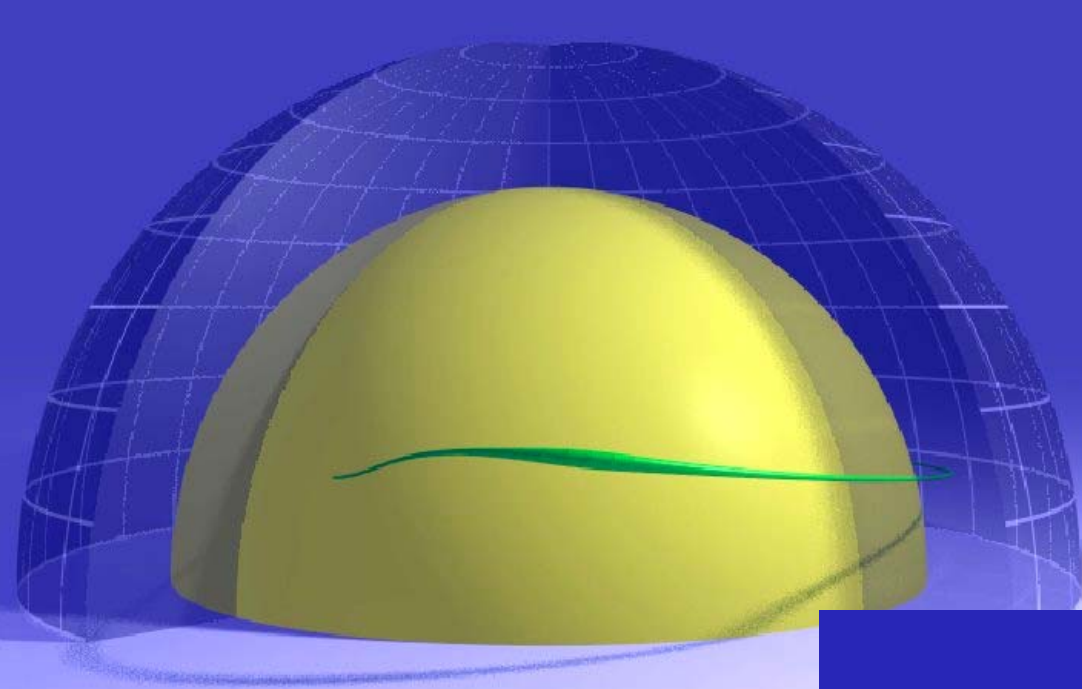


Reversal of the meridional field

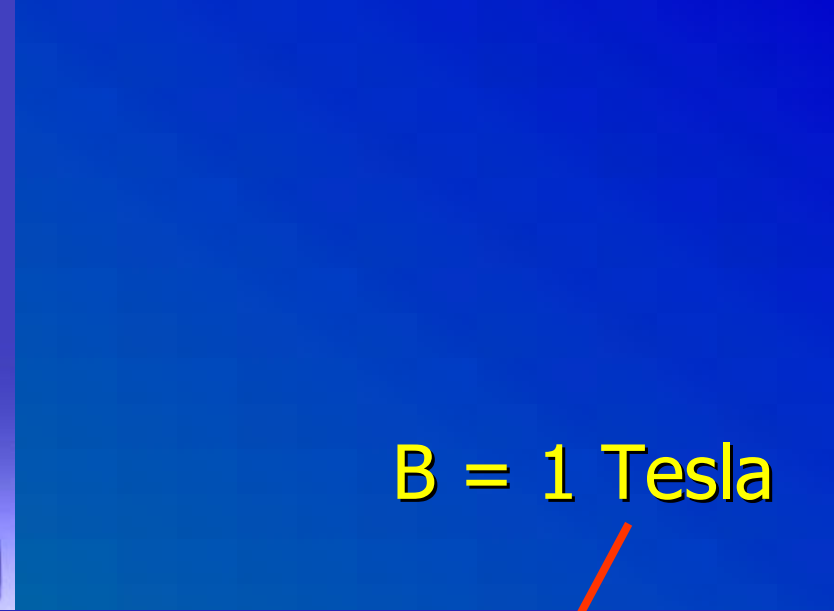


Reversal of the meridional field

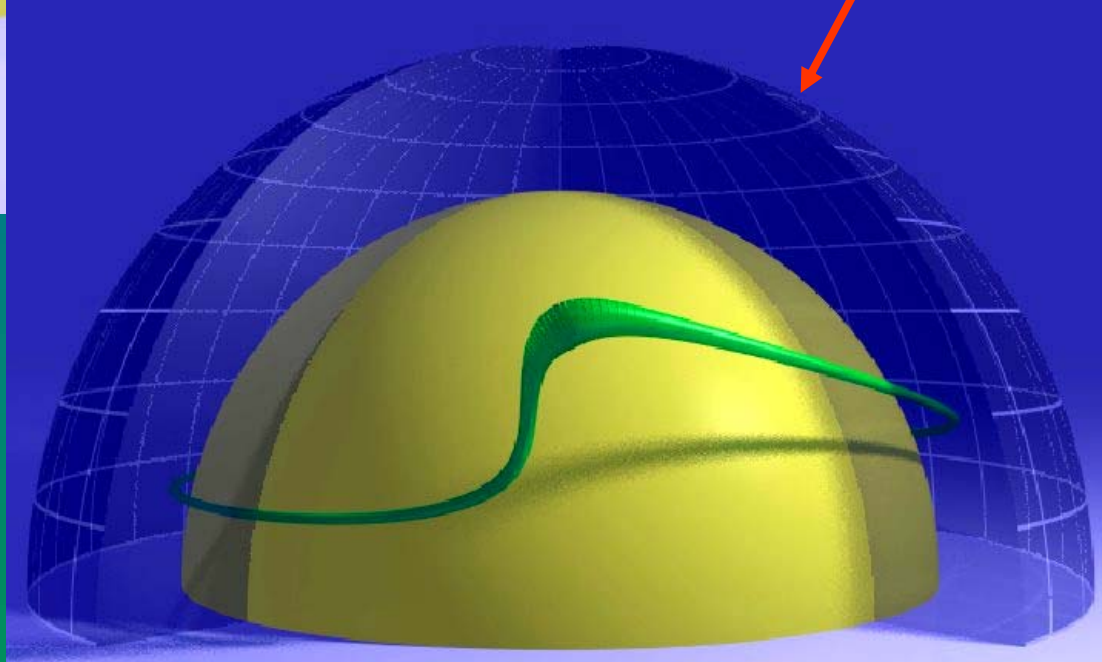




$B = 10$ Tesla



$B = 1$ Tesla



The end...

